## Problem 1:

Since the potential on the surfaces is given we need to use the Green function for Dirichlet boundary conditions that was obtained in class:

$$G_D(x, y, x', y') = 4 \sum_{n=1}^{\infty} \frac{e^{-\frac{\pi n}{a}(x_> - x_<)}}{n} \sin \frac{n\pi y}{a} \sin \frac{n\pi y'}{a},$$
 (1)

where  $x_{>}(x_{<})$  is the larger (smaller) between x and x'.

In the problem we are considering the density of charge is

$$\rho(x', y') = q\delta(x')\delta(y' - d). \tag{2}$$

The potential inside the volume defined by  $-\infty \le x \le \infty$  and  $0 \le y \le a$  is given by:

$$\Phi(x,y) = \frac{1}{4\pi\epsilon_0} \int_V G(x,x',y,y') \rho(x',y') dx' dy' - \frac{1}{4\pi} \oint_S \Phi_s \frac{\partial G_D}{\partial n'} dS'. \tag{3}$$

Since q is at x' = 0, I expect to have two different expressions for the potential, one for  $x \le 0$  and another for x > 0. Let us first calculate the volume integral which would give us the potential for the charge if all the surfaces were grounded. Because of Eq.(2) we obtain for  $x \leq 0$ :

$$\Phi_V^I(x,y) = \frac{q}{\pi\epsilon_0} \sum_{n=1}^{\infty} \frac{e^{\frac{\pi nx}{a}}}{n} \sin \frac{n\pi y}{a} \sin \frac{n\pi d}{a},\tag{4}$$

and for  $x \geq 0$ :

$$\Phi_V^{II}(x,y) = \frac{q}{\pi\epsilon_0} \sum_{n=1}^{\infty} \frac{e^{-\frac{\pi nx}{a}}}{n} \sin\frac{n\pi y}{a} \sin\frac{n\pi d}{a}.$$
 (5)

The two expressions can be combined as

$$\Phi_V(x,y) = \frac{q}{\pi\epsilon_0} \sum_{n=1}^{\infty} \frac{e^{-\frac{\pi n}{a}(x_> - x_<)}}{n} \sin \frac{n\pi y}{a} \sin \frac{n\pi d}{a},\tag{6}$$

where  $x_{>}(x_{<})$  is the larger (smaller) between x and 0.

Now let's calculate the surface integral that would give the potential due to the surface potentials in the absence of charges. Notice that this potential only depends on y and it is trivially given by  $\Phi(x,y) = \frac{Vy}{a}$ . Using the Green function we should obtain the same result, but it will be expanded in terms of  $\sin \frac{n\pi y}{a}$ .

The surface integral is given by:

$$-\frac{1}{4\pi} \oint_{S} \Phi_{s} \frac{\partial G_{D}}{\partial n'} dS'. \tag{7}$$

In this case the "surface" is the line parallel to the x-axis at y=a. The normal is  $\hat{n}'=\hat{y}'$ . Thus, we have to take the derivative with respect to y' and evaluate it at y' = a. Then,

$$\frac{\partial G_D}{\partial n'}|_S = \frac{\partial G_D}{\partial y'}|_{y'=a} = 4\pi \sum_{n=1}^{\infty} \frac{e^{-\frac{\pi n}{a}(x_> - x_<)}}{a} \sin \frac{n\pi y}{a} (-1)^n, \tag{8}$$

where we have used that  $\cos(n\pi y'/a)|_{y'=a} = (-1)^n$ .

Then the surface integral becomes:

$$\Phi_s(x,y) = -\frac{V}{a} \sum_{n=1}^{\infty} \frac{(-1)^n}{n} \sin \frac{n\pi y}{a} \int_{-\infty}^{\infty} e^{-\frac{\pi n}{a}(x_> - x_<)} dx'.$$
 (9)

Then we need to split the integral in two pieces: for  $-\infty \le x' \le x$  we have that  $x_< = x'$  and  $x_> = x$  while for  $x \le x' \le \infty$  we have that  $x_< = x$  and  $x_> = x'$ . Then we have that

$$\int_{-\infty}^{\infty} e^{-\frac{\pi n}{a}(x > -x_{<})} dx' = \int_{-\infty}^{x} e^{-\frac{\pi n}{a}(x - x')} dx' + \int_{x}^{\infty} e^{-\frac{\pi n}{a}(x' - x)} dx' = \frac{2a}{n\pi}.$$
 (10)

Putting this result in Eq.(9) we obtain:

$$\Phi_s(x,y) = -\frac{2V}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^n}{n} \sin \frac{n\pi y}{a}.$$
 (11)

Then the total potential is given by the sum of Eq.(6) and Eq.(11):

$$\Phi(x,y) = \frac{q}{\pi\epsilon_0} \sum_{n=1}^{\infty} \frac{e^{-\frac{\pi n}{a}(x_> - x_<)}}{n} \sin\frac{n\pi y}{a} \sin\frac{n\pi d}{a} - \frac{2V}{\pi} \sum_{n=1}^{\infty} \frac{(-1)^n}{n} \sin\frac{n\pi y}{a}.$$
 (12)

Important detail: notice that this is the same result that you obtained in Problem 6 of Hw#8. The surface term, i.e., Eq.(11), provides the contribution to the potential coming from the plane at potential V which was Vy/a. In Eq.(11) y is expanded in terms of the complete orthogonal basis formed by the functions  $\sin \frac{n\pi y}{a}$  in the interval [0,a].