

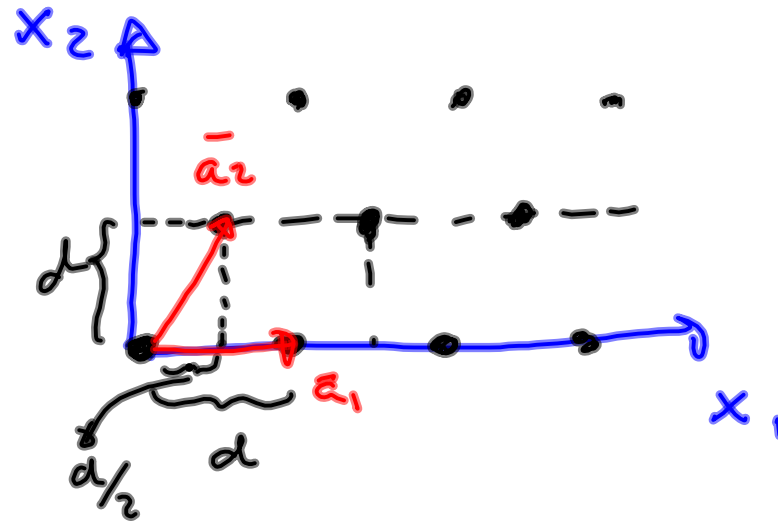
Covariant and contravariant

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basis:

For each system defined in real space by a covariant bases there is a contravariant basis that is orthogonal to the covariant basis. The contravariant basis is defined in the dual or reciprocal space.

Dual space and dual basis in a 2D crystal: Oblique Bravais lattice.



$$\bar{a}_1 = d \hat{e}_1$$

$$\bar{a}_2 = \frac{d}{2} \hat{e}_1 + d \hat{e}_2$$

$$|\bar{a}_1| = d$$

$$|\bar{a}_2| = d \frac{\sqrt{5}}{2}$$

$$\bar{R} = m_1 \bar{a}_1 + m_2 \bar{a}_2 \quad \text{all the points in the lattice.}$$

Let's find $\bar{b}_i \equiv \bar{b}^i$ the basis of the reciprocal lattice.

We know that in the reciprocal lattice a vector \bar{K} is given by:

$$\bar{K} = m_1 \bar{b}_1 + m_2 \bar{b}_2$$

$$\text{we want } \bar{K} \cdot \bar{R} = 2\pi n \quad \textcircled{1}$$

We found that $\textcircled{1}$ means that

$$a_i b^j = 2\pi \delta_{ij}$$

(2π replaced by 1
for a "mathematical"
system as in
HW)

$$\bar{a}_1 \cdot \bar{b}_1 = d b_{1x} = 2\pi \Rightarrow \boxed{b_{1x} = \frac{2\pi}{d}}$$

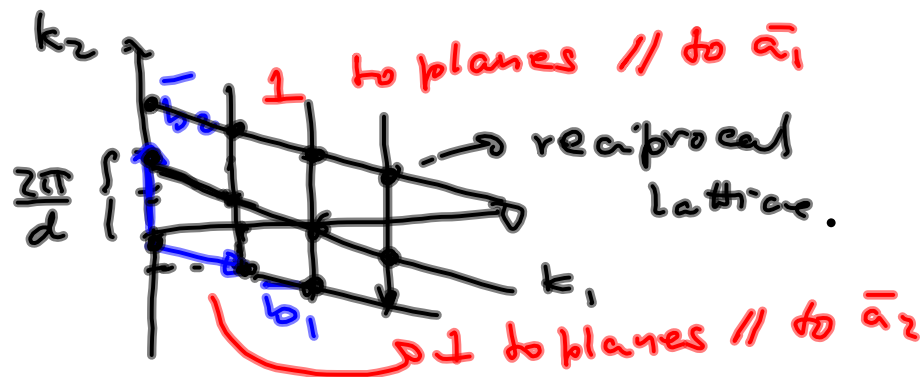
$$\bar{a}_2 \cdot \bar{b}_2 = \frac{d}{2} b_{2x} + d b_{2y} = 2\pi \Rightarrow \boxed{b_{2y} = \frac{2\pi}{d}}$$

$$\bar{a}_1 \cdot \bar{b}_2 = 0 = d b_{2x} \Rightarrow \boxed{b_{2x} = 0}$$

$$\bar{a}_2 \cdot \bar{b}_1 = 0 = \frac{d}{2} b_{1x} + d b_{1y} \Rightarrow \boxed{b_{1y} = -\frac{b_{1x}}{2} = -\frac{\pi}{d}}$$

$$\bar{b}_1 = \frac{2\pi}{d} \left(1, -\frac{1}{2} \right)$$

$$\bar{b}_2 = \frac{2\pi}{d} (0, 1)$$

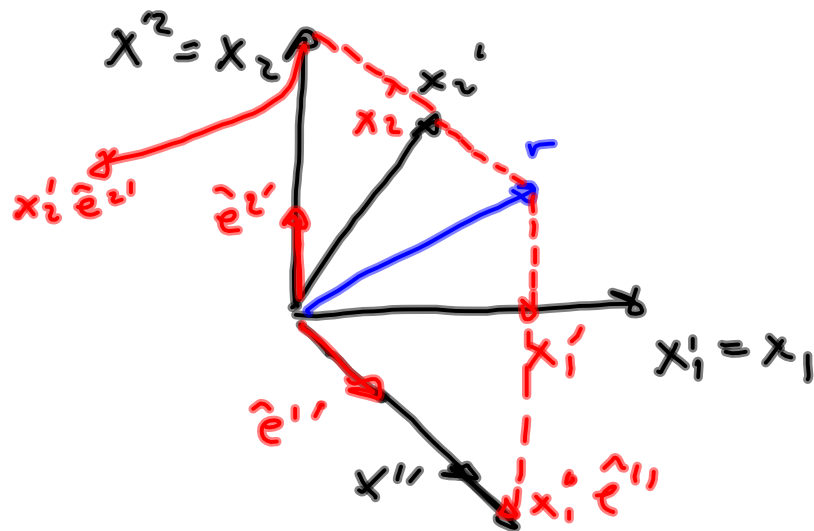
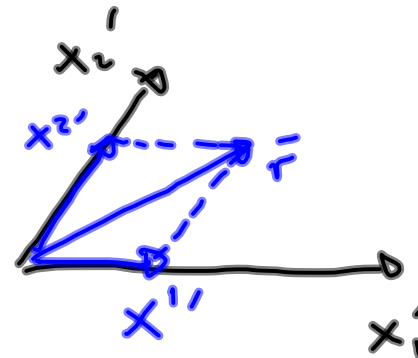


In homework you will show that:

$$\bar{r}' = x'^i \hat{e}'_i \quad \text{in real space.}$$

and

$$\bar{r}' = x'_i \hat{e}^i \quad \text{in dual space.}$$



\hat{e}^i are NOT unit vectors

In dual space and rescaled by \hat{e}^i the + projections (covariant) because the "parallel" projection

Covariant and Contravariant vectors.

We have defined our prototype contravariant vector $\bar{r} = r^i$ that transforms as

$$r'^i = \frac{\partial x'^i}{\partial x^j} r^j$$

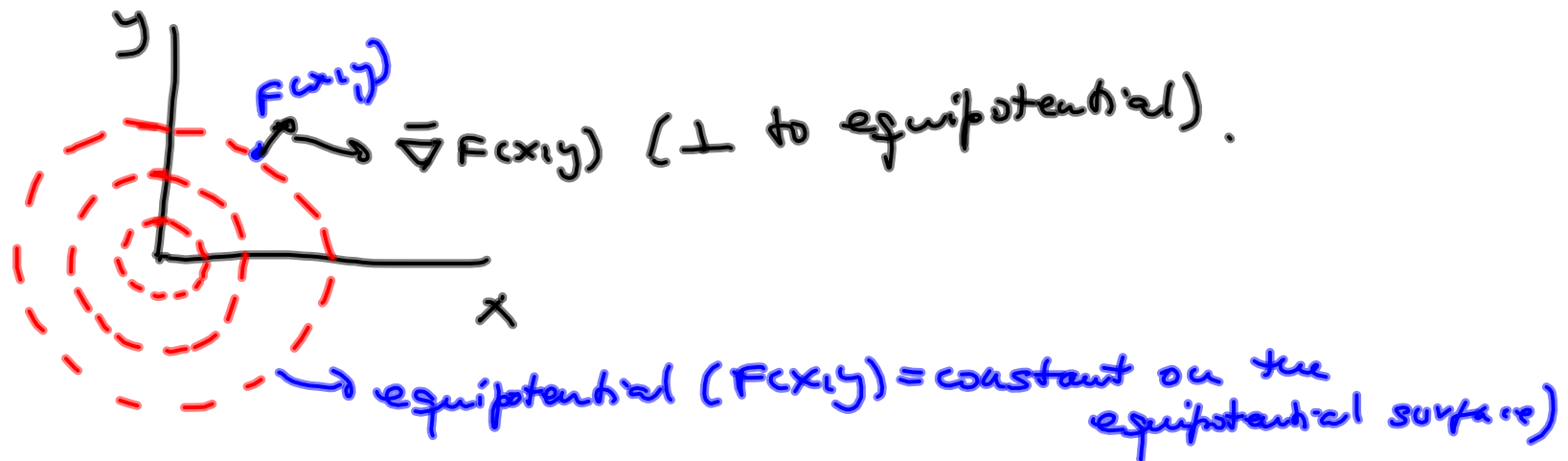
going from a system K to a system K' .

Now let's see how the gradient of an scalar function transforms going from K to K' .

Gradient

The gradient of a scalar field $F(x,y)$ is a vector.

$$\vec{\nabla} F(x,y) = \frac{\partial F}{\partial x} \hat{x} + \frac{\partial F}{\partial y} \hat{y} = \left(\frac{\partial F}{\partial x}, \frac{\partial F}{\partial y} \right)$$



$$\nabla F \equiv \frac{\partial F}{\partial x^i} \equiv \bar{B} \equiv B_i \rightarrow \text{we'll see why}$$

$$\bar{r} \equiv r^i \hat{e}_i$$

rotation
real meaning

Remember that $r'^i = \frac{\partial x'^i}{\partial x^j} r^j$

Now let's find out how \bar{B} transforms:
 ① $F(x', y') = F(x, y)$ because it is a scalar.

$$\nabla' F = \bar{B}' = (B'_x, B'_y) = \left(\frac{\partial F(x', y')}{\partial x'}, \frac{\partial F(x', y')}{\partial y'} \right) =$$

$$\stackrel{\text{①}}{=} \left(\frac{\partial F(x, y)}{\partial x'}, \frac{\partial F(x, y)}{\partial y'} \right) = \left(\underbrace{\frac{\partial F(x, y)}{\partial x}}_{B_x} \frac{\partial x}{\partial x'} + \underbrace{\frac{\partial F(x, y)}{\partial y}}_{B_y} \frac{\partial y}{\partial x'}, \right.$$

$$\left. \underbrace{\frac{\partial F(x, y)}{\partial x}}_{B_x} \frac{\partial x}{\partial y'} + \underbrace{\frac{\partial F(x, y)}{\partial y}}_{B_y} \frac{\partial y}{\partial y'} \right) = \left(B_x \frac{\partial x}{\partial x'} + B_y \frac{\partial y}{\partial x'}, B_x \frac{\partial x}{\partial y'} + B_y \frac{\partial y}{\partial y'} \right)$$

In compact way:

$$\textcircled{1} \quad B'_i = \frac{\partial x_j}{\partial x'^i} B_j$$

for r^i we had

$$\textcircled{2} \quad r'^i = \frac{\partial x'^i}{\partial x^j} r^j$$

Check:

$$B'_1 = \frac{\partial x_1}{\partial x'^1} B_1 + \frac{\partial x_2}{\partial x'^1} B_2$$

or

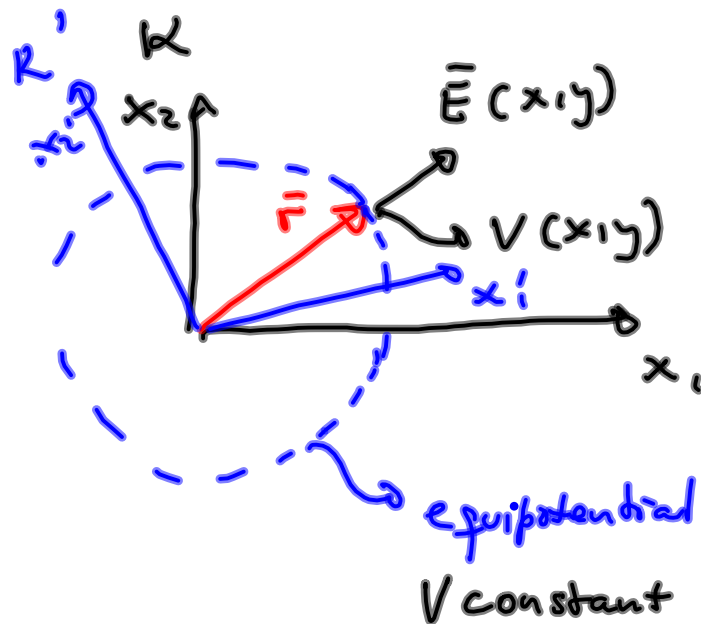
$$B'_x = \frac{\partial x}{\partial x'} B_x + \frac{\partial y}{\partial x'} B_y$$

So $\textcircled{1}$ transforms as
the inverse of $\textcircled{2}$

B_i is our prototype covariant vector.

Example of covariant vector: Electric Field.

Consider a point charge in 2D:



$$\text{At } \vec{r} \quad V(x_1, y) = V(x'_1, y')$$

In K

$$\begin{aligned} V(x_1, y) &= V(r) = -\frac{A}{r} = \\ &= -\frac{A}{(x^2 + y^2)^{1/2}} \end{aligned}$$

In K'

$$\begin{aligned} V(x'_1, y'_1) &= V(r') = -\frac{A}{r'} = \\ &= -\frac{A}{(x'^2 + y'^2)^{1/2}} \end{aligned}$$

Let's obtain \bar{E} and \bar{E}' :

$$\begin{aligned} \text{In } \mathcal{K}: \quad \bar{E}(x, y) &= -\bar{\nabla} V(x, y) = \left(-\frac{\partial V}{\partial x}, -\frac{\partial V}{\partial y}\right) = \\ &= \frac{A}{r^3}(x, y) = \frac{A}{r^3} \bar{r} \end{aligned}$$

$$\begin{aligned} \text{In } \mathcal{K}': \quad \bar{E}'(x', y') &= -\bar{\nabla}' V(x', y') = \left(-\frac{\partial V}{\partial x'}, -\frac{\partial V}{\partial y'}\right) = \\ &= \frac{A}{r'^3}(x', y') = \frac{A}{r'^3} \bar{r}' \end{aligned}$$

$$E'_i(x') = -\frac{\partial V'}{\partial x'^i} = -\frac{\partial V}{\partial x^j} \frac{\partial x^j}{\partial x'^i} = E_j \frac{\partial x^j}{\partial x'^i} \quad \text{covariant vector}$$

Now you see that the work done by the electric force can be written as:

$$dW = \vec{F} \cdot d\vec{S} = q \vec{E} \cdot d\vec{S} \equiv q \begin{matrix} E_i & dS^i \\ \downarrow & \downarrow \\ \text{covariant} & \text{contravariant} \end{matrix}$$

$$= q (E_x, E_y, E_z) \begin{pmatrix} dx \\ dy \\ dz \end{pmatrix}$$

Notation:

$$\frac{\partial}{\partial x^i} \equiv \partial_i \quad \text{covector}$$

$$\bar{\nabla} F = \bar{B} \equiv B_i = \frac{\partial F}{\partial x^i} \equiv \partial_i F$$

Total differential:

$$dF = \bar{\nabla} F \cdot d\bar{r} = \partial_i F dx^i \quad \text{scalar}$$

Divergence:

$$\bar{\nabla} \cdot \bar{G} = \frac{\partial G_x}{\partial x} + \frac{\partial G_y}{\partial y} + \frac{\partial G_z}{\partial z} \equiv \partial_i G^i \quad \text{scalar}$$