How about the ground state wave function? Remember:

$$\psi_{nlm}(r,\theta,\phi) = R_{nl}(r) Y_l^m(\theta,\phi)$$

$$R_{nl}(r) = \frac{1}{r} \rho^{l+1} e^{-\rho} v(\rho) \qquad u(\rho) = \rho^{l+1} e^{-\rho} v(\rho) \qquad \rho \equiv \kappa r$$

$$R(r) = u(r)/r$$

n=1 is the lowest energy (the most negative). Since $n = j_{max} + l + 1$, then *l*=0, $j_{max} = 0$. Then

$$\begin{split} \psi_{100}(r,\theta,\phi) &= R_{10}(r)Y_0^0(\theta,\phi) \\ \text{with} \quad R_{10}(r) &= \frac{c_0}{a}e^{-r/a} \text{ because} \\ \text{the polynomial has only a} \qquad Y_0^0(\theta,\phi) &= 1/\sqrt{4\pi} \\ \text{constant } c_0. \end{split}$$

We then normalize, which means we fix the value of c_0 :

$$\int_0^\infty |R_{10}|^2 r^2 \, dr = \frac{|c_0|^2}{a^2} \int_0^\infty e^{-2r/a} \, r^2 \, dr = |c_0|^2 \frac{a}{4} = 1$$

The very final result (celestial music ...) is:

$$\psi_{100}(r,\theta,\phi) = \frac{1}{\sqrt{\pi a^3}} e^{-r/a}$$

which, amazingly, is remarkably simple.

Consider now the excited states, starting with n=2:

$$E_2 = \frac{-13.6 \text{ eV}}{4}$$

= -3.4 eV

Because $n = j_{max} + l + 1$, then n=2 allows for l=0, $j_{max}=1$ or l=1 (m=-1,0,+1), $j_{max}=0$. This will be the 2s and three 2p's orbitals. Degeneracy 4 (in general, degeneracy is n²).

(1) If *I=0*, $j_{max}=1$, then $c_{jmax+1}=c_2=0$. Use $c_1 = [2(j + l + 1 - n) / (j+1)(j+2l+2)] c_0$ with n=2, j=0 (because c_0 means j=0) and *I=0*, and you get $c_1 = -c_0$ so the polynomial becomes $v(\rho)=c_0(1-\rho)$ with c_0 again used to normalize.

$$R_{20}(r) = \frac{c_0}{2a} \left(1 - \frac{r}{2a} \right) e^{-r/2a}$$

$$Y_0^0(\theta,\phi) = 1/\sqrt{4\pi}$$

The factors 2 arise from $\rho = (1/an)r$ with n=2.

(2) If *I=1*, j_{max} = 0, you get c_{jmax+1} = c_1 = 0 so the polynomial is only c_0 as it happens in the ground state.

Use the general formula $R_{nl}(r) = \frac{1}{r}\rho^{l+1}e^{-\rho}v(\rho)$ with *l=1* and *n=2*.

$$R_{21}(r) = \frac{c_0}{4a^2} r e^{-r/2a} \qquad \begin{array}{c} \rho = (1/an)r \\ \text{with n=2} \end{array}$$

The front factor comes from $\rho^2/r = (r/2a)^2/r$.

The spherical harmonics are:

$$Y_1^0 = \left(\frac{3}{4\pi}\right)^{1/2} \cos \theta$$
$$Y_1^{\pm 1} = \mp \left(\frac{3}{8\pi}\right)^{1/2} \sin \theta e^{\pm i\phi}$$

In general, the polynomials $v(\rho)$ are called **associated Laguerre polynomials**

$$v(\rho) = L_{n-l-1}^{2l+1}(2\rho)$$
 $\rho = (1/an) r$

and there are tables with these polynomials.

Putting all together, for arbitrary (*n,l,m*) the normalized wave functions are:

$$\psi_{nlm} = \sqrt{\left(\frac{2}{na}\right)^3 \frac{(n-l-1)!}{2n[(n+l)!]^3}} e^{-r/na} \left(\frac{2r}{na}\right)^l \left[L_{n-l-1}^{2l+1}\left(\frac{2r}{na}\right)\right] Y_l^m(\theta,\phi)$$

$$R_{nl}(r)$$

$$\begin{array}{l} R_{10} = 2a^{-3/2} \exp(-r/a) \\ R_{20} = \frac{1}{\sqrt{2}} a^{-3/2} \left(1 - \frac{1}{2} \frac{r}{a}\right) \exp(-r/2a) \\ R_{21} = \frac{1}{\sqrt{24}} a^{-3/2} \frac{r}{a} \exp(-r/2a) \\ R_{30} = \frac{2}{\sqrt{27}} a^{-3/2} \left(1 - \frac{2}{3} \frac{r}{a} + \frac{2}{27} \left(\frac{r}{a}\right)^2\right) \exp(-r/3a) \\ R_{31} = \frac{8}{27\sqrt{6}} a^{-3/2} \left(1 - \frac{1}{6} \frac{r}{a}\right) \left(\frac{r}{a}\right) \exp(-r/3a) \\ R_{32} = \frac{4}{81\sqrt{30}} a^{-3/2} \left(1 - \frac{1}{6} \frac{r}{a}\right) \left(\frac{r}{a}\right) \exp(-r/3a) \\ R_{40} = \frac{1}{4} a^{-3/2} \left(1 - \frac{3}{4} \frac{r}{a} + \frac{1}{8} \left(\frac{r}{a}\right)^2 - \frac{1}{192} \left(\frac{r}{a}\right)^3\right) \exp(-r/4a) \\ R_{41} = \frac{\sqrt{5}}{16\sqrt{3}} a^{-3/2} \left(1 - \frac{1}{4} \frac{r}{a} + \frac{1}{80} \left(\frac{r}{a}\right)^2\right) \frac{r}{a} \exp(-r/4a) \\ R_{42} = \frac{1}{64\sqrt{5}} a^{-3/2} \left(1 - \frac{1}{12} \frac{r}{a}\right) \left(\frac{r}{a}\right)^2 \exp(-r/4a) \\ R_{43} = \frac{1}{768\sqrt{35}} a^{-3/2} \left(\frac{r}{a}\right)^3 \exp(-r/4a) \\ \end{array}$$



