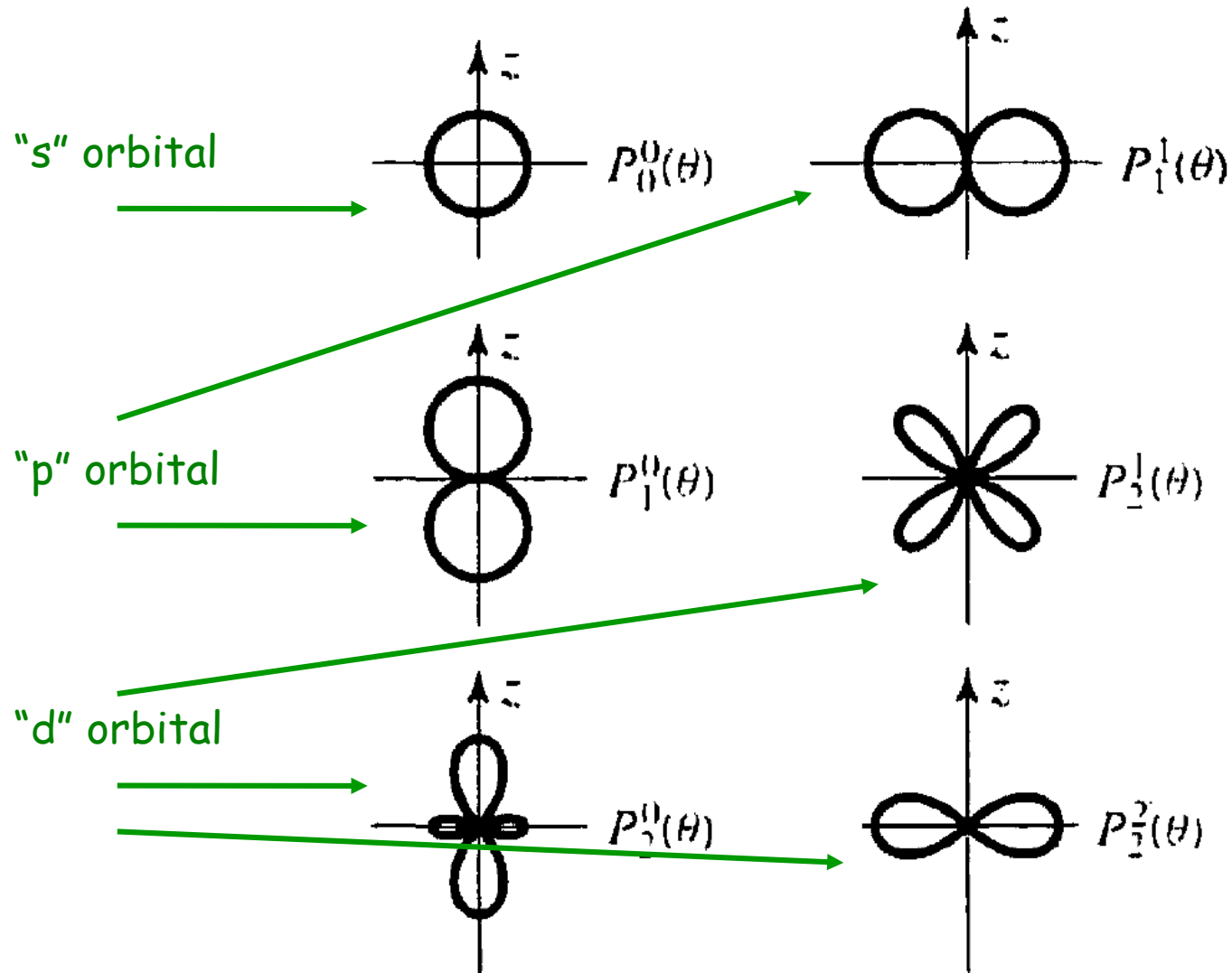


If you graph these functions, the "orbitals" that you have seen many times before since high school start appearing!



The final "touch" requires putting all together

$$Y(\theta, \phi) = \Theta(\theta)\Phi(\phi)$$

where

$$\Theta(\theta) = A P_l^m(\cos \theta)$$

and

$$\Phi(\phi) = e^{im\phi}$$

with the "quantization" condition:

$$l = 0, 1, 2, \dots; \quad m = -l, -l + 1, \dots, -1, 0, 1, \dots, l - 1, l$$

We also need to **normalize** using  $d^3\mathbf{r} = r^2 \sin \theta dr d\theta d\phi$

$$\int |\psi|^2 r^2 \sin \theta dr d\theta d\phi = \underbrace{\int |R|^2 r^2 dr}_{=1} \underbrace{\int |Y|^2 \sin \theta d\theta d\phi}_{=1} = 1$$

For the angular component this means:

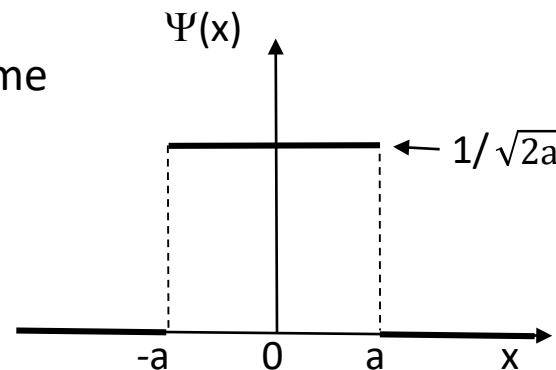
$$\int_0^{2\pi} \int_0^{\pi} |Y|^2 \sin \theta d\theta d\phi = 1$$

By this procedure the famous **spherical harmonics** arise:

one "s"	$Y_0^0 = \left(\frac{1}{4\pi}\right)^{1/2}$	$Y_2^{\pm 2} = \left(\frac{15}{32\pi}\right)^{1/2} \sin^2 \theta e^{\pm 2i\phi}$	
pz	$Y_1^0 = \left(\frac{3}{4\pi}\right)^{1/2} \cos \theta$	$Y_3^0 = \left(\frac{7}{16\pi}\right)^{1/2} (5 \cos^3 \theta - 3 \cos \theta)$	
three "p" px, py	$Y_1^{\pm 1} = \mp \left(\frac{3}{8\pi}\right)^{1/2} \sin \theta e^{\pm i\phi}$	$Y_3^{\pm 1} = \mp \left(\frac{21}{64\pi}\right)^{1/2} \sin \theta (5 \cos^2 \theta - 1) e^{\pm i\phi}$	
five "d"	$Y_2^0 = \left(\frac{5}{16\pi}\right)^{1/2} (3 \cos^2 \theta - 1)$	$Y_3^{\pm 2} = \left(\frac{105}{32\pi}\right)^{1/2} \sin^2 \theta \cos \theta e^{\pm 2i\phi}$	seven "f"
	$Y_2^{\pm 1} = \mp \left(\frac{15}{8\pi}\right)^{1/2} \sin \theta \cos \theta e^{\pm i\phi}$	$Y_3^{\pm 3} = \mp \left(\frac{35}{64\pi}\right)^{1/2} \sin^3 \theta e^{\pm 3i\phi}$	

Question

Consider the normalized wave function in the figure. Assume it is the wave function  $\Psi(x,0)$  at time  $t=0$  of a wave packet problem. (a) Find  $\phi(k)$ . Do the simple integral. (b) Now find  $\Psi(x,t)$ . Leave this result just expressed as an integral in  $k$ .

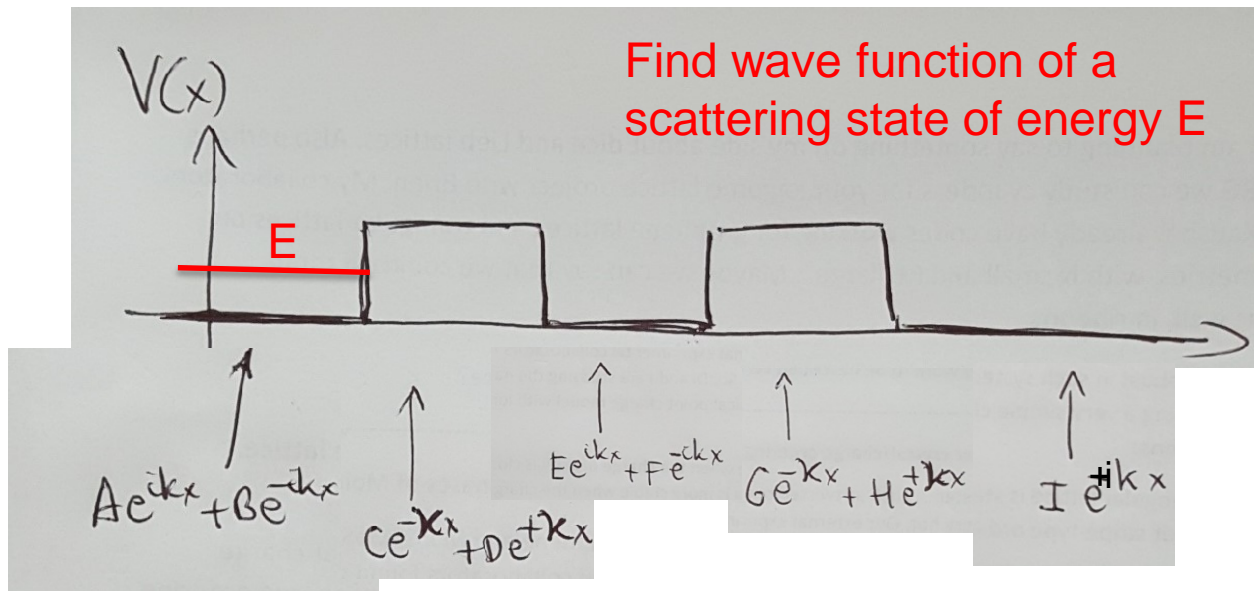


$$(a) \quad \phi(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} \Psi(x,0) e^{-ikx} dx = \boxed{\frac{\sin(ka)}{\sqrt{\pi a} k}}$$

$$(b) \quad \Psi(x,t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} \phi(k) e^{i(kx - \frac{\hbar k^2}{2m} t)} dk =$$

$$= \boxed{\frac{1}{\pi \sqrt{2a}} \int_{-\infty}^{+\infty} \frac{\sin(ka)}{k} e^{i(kx - \frac{\hbar k^2}{2m} t)} dk}$$

Question



How many unknowns? 9  $A, B, C, D, E, F, G, H, I$   
 $E$  being scattering state is arbitrary

How many boundary condition equations?  
4 frontiers, i.e. 8 (4 and  $\frac{d\psi}{dx}$  continuous)

Best dividing by  $A$  we get 8 unknowns.

Question

Consider the infinite square well of width "a".

Is the wave function  $\Psi(x) = \cos(\pi x/a)$  an eigenstate of the momentum operator?

Apply p operator over  $\Psi(x)$  and see if you obtain a number times  $\Psi(x)$

$$(-i \hbar d/dx)\Psi(x) = (-i \hbar) d/dx \cos(\pi x/a) = (i\pi\hbar/a) \sin(\pi x/a)$$

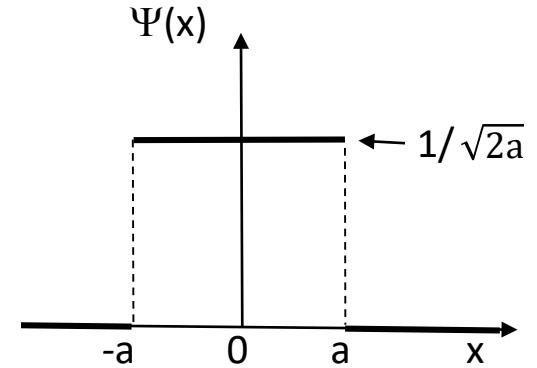
**Answer: No, it is not eigenstate.**

Note: you can also use Euler formulas,  $(-i \hbar d/dx)\Psi(x) = (-i \hbar/2) d/dx (e^{i\pi x/a} + e^{-i\pi x/a}) = (-i \hbar/2) (i\pi/a) [e^{i\pi x/a} - e^{-i\pi x/a}] = (i\pi\hbar/a) \sin(\pi x/a)$

Consider the normalized wave function in the figure:

(a) Find the coefficient  $c(p) = \langle f_p | \Psi \rangle$ , where  $f_p$  is the eigenfunction of the linear momentum with eigenvalue  $p$ .

(b) Find the probability that in a measurement, the momentum  $p$  is between 0 and  $\infty$ . Use  $\int_0^{\infty} \frac{\sin^2(x)}{x^2} dx = \frac{\pi}{2}$



(a) Find  $c(p)$ :

$$c(p) = \langle f_p | \Psi \rangle = \frac{1}{\sqrt{2\pi\hbar}} \int_{-a}^a e^{-i\frac{p}{\hbar}x} \frac{1}{\sqrt{2a}} dx = \boxed{\sqrt{\frac{\hbar}{\pi a}} \frac{1}{p} \sin\left(\frac{pa}{\hbar}\right)}$$

(b)

$$P_{\text{prob.}} = \int_{[0, \infty]} |c(p)|^2 dp = \boxed{\frac{\hbar}{\pi a} \int_0^{\infty} \frac{1}{p^2} \sin^2\left(\frac{pa}{\hbar}\right) dp} = \boxed{\frac{1}{2}}$$

Question

$\hat{A}$  and  $\hat{B}$  are Hermitian.

Thus,  $\hat{A} = \hat{A}^\dagger$  and  $\hat{B} = \hat{B}^\dagger$ .

Find the Hermitian of the product operator  $\hat{O} = \hat{A}\hat{B}$ .

We are asked to find the operator  $\hat{O}$  that

satisfies  $\langle \hat{O}^\dagger \psi | \psi \rangle = \langle \psi | \hat{O} \psi \rangle$

Let us start. We need to move  $\hat{A}$  and  $\hat{B}$  to the other side.

$$\langle \psi | \underbrace{\hat{A}\hat{B}}_{\hat{O}} \psi \rangle = \langle \hat{A} \psi | \hat{B} \psi \rangle = \langle \underbrace{\hat{B}\hat{A}}_{\hat{O}^\dagger} \psi | \psi \rangle$$

use  $\hat{A}^\dagger = \hat{A}$

then,  $\hat{O}^\dagger = \hat{B}\hat{A} \neq \hat{O} = \hat{A}\hat{B}$



Question

$$\text{Find } \left(\frac{d}{dx}\right)^\dagger$$

$$\langle f | \frac{d}{dx} g \rangle = \int_{-\infty}^{+\infty} f^*(x) \frac{d}{dx} g(x) dx =$$

by parts  $\uparrow$   $=$

$$\left. f^*(x) g(x) \right|_{-\infty}^{+\infty} - \int_{-\infty}^{+\infty} \left[ \frac{d}{dx} f^*(x) \right] g(x) dx = \int_{-\infty}^{+\infty} \left[ -\frac{d}{dx} f^*(x) \right] g(x) dx$$

$\approx 0$  because  
we use square  
integrable functions

$$f(x) \rightarrow 0$$

$x \rightarrow \pm\infty$

$$g(x) \rightarrow 0$$

$x \rightarrow \pm\infty$

$$= \langle -\frac{d}{dx} f | g \rangle$$

Then  $\left(\frac{d}{dx}\right)^\dagger = -\frac{d}{dx}$