## Announcements:

- Average of HW1 was 48/60 points which is 32/40%, within expected range.
- For HW2, in problem 2.7(b) you need some integrals of the "x sin(x)" type. Just find them in tables or use Mathematica.

Also, use a first page with only your name, and the rest blank. Grader will write provide grade on the second page. Privacy issue.

- Thursday 30 minutes after class, I must give an invited zoom talk. If you send questions Wed late, I may not be able to answer fast. Try to send HW2 related questions not Wed very late, but today afternoon/evening or Wed morning if possible.
- October 4 will be by Zoom.
- Test 2 will be November 1.

(3) Returning to page 28 Ch2 book. How to include time? There are several possible values of E, say  $E_1, E_2, E_3, ...,$  as found in square well example. For each "allowed" energy, there is a solution of time-indep. Sch. Eq. with its "phase factor"

 $\Psi_1(x,t) = \psi_1(x)e^{-iE_1t/\hbar}, \quad \Psi_2(x,t) = \psi_2(x)e^{-iE_2t/\hbar}, \quad \dots$ 

Make a linear 
$$\Psi(x,t) = \sum_{n=1}^{\infty} c_n \psi_n(x) e^{-iE_nt/\hbar}$$
  
combination:

Statement: any wave function  $\Psi(x,t)$  can be written as above. The  $c_n$ 's are the same as before i.e. time INDEPENDENT. But the linear combination above is NOT a stationary state i.e.  $|\Psi(x,t)|^2$  at fixed x, changes with t.

## **Not in book:** Is this a solution of the time-dep. Sch. Eq. with V(x)?

$$\hat{H}\Psi(x,t) = \sum_{n=1}^{\infty} c_n \underbrace{\hat{H}\psi_n(x)e^{-iE_nt/\hbar}}_{E_n\psi_n(x)}$$
(a)

$$i\hbar \frac{\partial}{\partial t} \Psi(x,t) = \sum_{n=1}^{\infty} c_n \psi_n(x) i\hbar \frac{\partial}{\partial t} e^{-iE_n t/\hbar}$$
 (b)  
 $E_n e^{-iE_n t/\hbar}$ 

(a) is equal to (b)  $\rightarrow$  The linear combination is solution of the time-dependent Sch. Eq., although is not a stationary state.

## Final procedure for $\Psi(x,t)$ in square well:

Given an arbitrary  $\Psi(x,0)$  -- satisfying the Bound. Conditions -- you are asked for  $\Psi(x,t)$ .

(1) Find stationary states and energies.

(2) "Somehow" do the integrals for  $c_n = \sqrt{\frac{2}{a}} \int_0^a \sin\left(\frac{n\pi}{a}x\right) \Psi(x, 0) dx$  coefficients:

(3) Done! 
$$\Psi(x,t) = \sum_{n=1}^{\infty} c_n \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi}{a}x\right) e^{-i(n^2\pi^2\hbar/2ma^2)t}$$

ψ**(**x,0)

X

The procedure is general but of course  $\psi_n(x)$  and  $E_n$  are diff for diff potentials. This formula is only for square well. Example 2.1, page 29 book: Assume you are given at t=0 (typical exam problem!):

Real numbers for simplicity. Assume state normalized i.e.  $|c_1|^2 + |c_2|^2 = 1$ 

$$\Psi(x,0) = c_1 \psi_1(x) + c_2 \psi_2(x)$$

Stationary states sin(nπx/a). NOT any arbitrary function like, say, e -|×|

PUNCHLINE: Time dependence is now **trivially** obtained! Eqs 2.16-2.17

$$\Psi(x,t) = c_1 \psi_1(x) e^{-iE_1 t/\hbar} + c_2 \psi_2(x) e^{-iE_2 t/\hbar}$$

 $|\Psi(x,t)|^2 = (c_1\psi_1 e^{iE_1t/\hbar} + c_2\psi_2 e^{iE_2t/\hbar})(c_1\psi_1 e^{-iE_1t/\hbar} + c_2\psi_2 e^{-iE_2t/\hbar})$ 

 $= c_1^2 \psi_1^2 + c_2^2 \psi_2^2 + 2c_1 c_2 \psi_1 \psi_2 \cos[(E_2 - E_1)t/\hbar].$ 

The prob. density is now time dependent even if stationary states are combined ("quantum beat").

**Example 2.2:** how to apply the recipe (typical exam problem!) Usual A used to normalize.  $\Psi(x, 0)$  $\Psi(x,0) = Ax(a-x), \quad (0 \le x \le a)$  $\frac{Aa^2}{4}$ Given to you as input, not as easy as previous page problem. "Even" with respect to middle reflection. Resembles "cosine" а ground state but it's not. Normalize first:  $1 = \int_0^a |\Psi(x,0)|^2 dx = |A|^2 \int_0^a x^2 (a-x)^2 dx \quad \longrightarrow \quad A = \sqrt{\frac{30}{a^5}}$ 

## Find coefficients $c_n$ : $c_n = \sqrt{\frac{2}{a}} \int_0^a \sin(\frac{n\pi}{a}x) \Psi(x,0) dx$ Only *n* odd is nonzero $c_n = \sqrt{\frac{2}{a}} \int_0^a \sin(\frac{n\pi}{a}x) \sqrt{\frac{30}{a^5}} x(a-x) dx = 8\sqrt{15}/(n\pi)^3$ $\Psi(x,0)$ See integration process in book. Even *n* (i.e. odd functions) gives 0 by symmetry.

You are ready to write the final answer:

$$\Psi(x,t) = \sqrt{\frac{30}{a}} \left(\frac{2}{\pi}\right)^3 \sum_{n=1,3,5,...} \frac{1}{n^3} \sin\left(\frac{n\pi}{a}x\right) e^{-in^2\pi^2\hbar t/2ma^2}$$
  
Often forgotten by  
students! It contains the  
time dependence!

We can learn more from this example:

For instance,  $\sum_{n} |c_{n}|^{2} = 1$ . You can verify that adding, say, the first 10 terms. General reason?

$$1 = \int |\Psi(x,0)|^2 dx = \int \left(\sum_{m=1}^{\infty} c_m \psi_m(x)\right)^* \left(\sum_{n=1}^{\infty} c_n \psi_n(x)\right) dx$$

$$= \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} c_m^* c_n \int \psi_m(x)^* \psi_n(x) \, dx$$
$$= \sum_{n=1}^{\infty} \sum_{m=1}^{\infty} c_m^* c_n \delta_{mn} = \sum_{n=1}^{\infty} |c_n|^2.$$

This happens only if the given  $\Psi(x,0)$  is already normalized to 1.

Thus,  $\sum_{n=1}^{\infty} |c_n|^2 = 1$  is equivalent to normalization 1, i.e. the probability of finding the particle "somewhere" inside the well is 100%.

Moreover,  $|c_1|^2 = 0.99855...$  i.e. other coefficients are very small. Why? Because  $\Psi(x,0)$  closely resembles the ground state! Develop intuition!

If you measure the energy, it will be shown that  $|c_n|^2$  is the probability that you will find  $E_n$  as result. Here, the chance of measuring  $E_1$  is high.

It can also be shown  $\langle \hat{H} \rangle = \sum_{n=1}^{\infty} |c_n|^2 E_n$ (book p37) that:

In general, for  $\langle \hat{O} \rangle$ , with  $\hat{O}$  any operator, this is not true.

Intuition! Read book:  $\langle \hat{H} \rangle$  is only slightly above the ground state energy  $E_1$  compatible with  $|c_1|^2 \sim 1$ 



The information about the provided initial wave function is hidden in the coefficients via

$$c_n = \sqrt{\frac{2}{a}} \int_0^a \sin\left(\frac{n\pi}{a}x\right) \Psi(x,0) \, dx$$