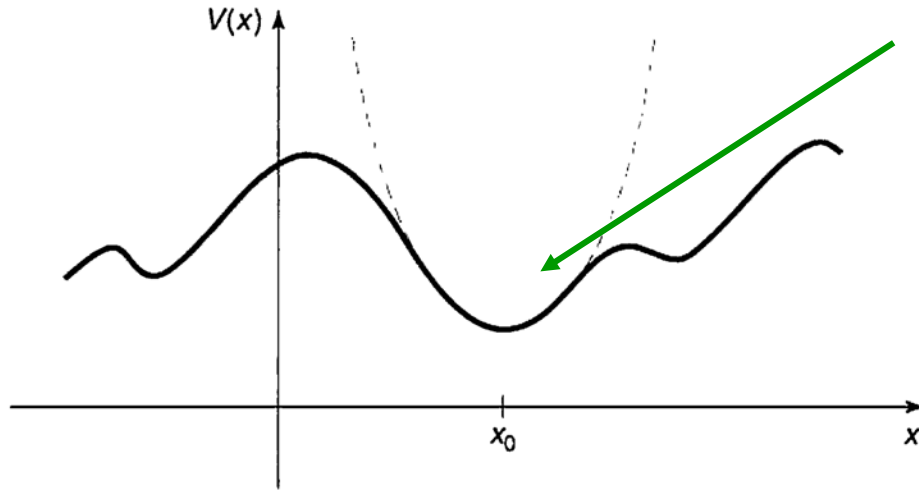


2.3 The harmonic oscillator



Any potential minimum can be approx. by a harmonic oscillator, if oscillations are small. Thus, this is widely used! Even E&M is quantized this way.

We wish to study the harmonic oscillator from the QM perspective:

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + \underbrace{\frac{1}{2} m \omega^2 x^2}_{V(x)} \psi = E \psi$$

$$V(x) = \frac{1}{2} m \omega^2 x^2$$

Two methods will be used to arrive to the same answer:

Algebraic method (uses $\hat{a}+$ and $\hat{a}-$ operators)

Analytic method (uses polynomials)

Warning: the symbol $\hat{}$ will sometimes appear, sometimes not. You have to develop the ability to judge if a quantity is an operator or not.

Preliminary exercise: commutation relations

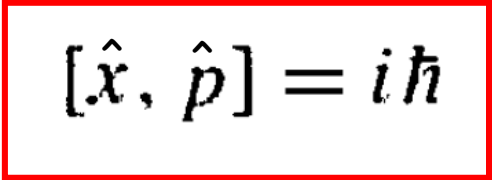
We define the commutator between operators \hat{A} and \hat{B} as:

$$[\hat{A}, \hat{B}] \equiv \hat{A}\hat{B} - \hat{B}\hat{A}$$

Suppose $\hat{A}=\hat{x}$ and $\hat{B}=\hat{p}$. Thus, we want $[\hat{x}, \hat{p}] = (\hat{x}\hat{p} - \hat{p}\hat{x})$. This is **NOT** zero because \hat{p} is the derivative operator.

To know its value you need a general "test function" $f(x)$.

$$[\hat{x}, \hat{p}]f(x) = \left[x \frac{\hbar}{i} \frac{d}{dx}(f) - \frac{\hbar}{i} \frac{d}{dx}(xf) \right] = \frac{\hbar}{i} \left(x \frac{df}{dx} - x \frac{df}{dx} - f \right) = i\hbar f(x)$$


$$[\hat{x}, \hat{p}] = i\hbar$$

In general, operators do not commute!

Finding the commutator of two given operators is a typical test problem.

2.3.1 Algebraic method

Sum of squares

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + \frac{1}{2} m \omega^2 x^2 \psi = E \psi \longrightarrow \frac{1}{2m} [\hat{p}^2 + (m\omega\hat{x})^2] \psi = E \psi$$

$\hat{p} \equiv (\hbar/i)d/dx$

For u, v real numbers $u^2 + v^2 = (iu + v)(-iu + v)$ but not for operators.

This factor is merely for future convenience.

By analogy try $\hat{a}_{\pm} \equiv \frac{1}{\sqrt{2\hbar m\omega}} (\mp i\hat{p} + m\omega\hat{x})$ and cross fingers.

$$\hat{a}_- \hat{a}_+ = \frac{1}{2\hbar m\omega} (i\hat{p} + m\omega\hat{x})(-i\hat{p} + m\omega\hat{x}) = \frac{1}{2\hbar m\omega} [\hat{p}^2 + (m\omega\hat{x})^2 - im\omega(\hat{x}\hat{p} - \hat{p}\hat{x})]$$

$$\hat{a}_- \hat{a}_+ = \frac{1}{2\hbar m\omega} [\hat{p}^2 + (m\omega\hat{x})^2] - \frac{i}{2\hbar} [\hat{x}, \hat{p}] = \frac{1}{\hbar\omega} \hat{H} + \frac{1}{2} \longrightarrow \hat{H} = \hbar\omega \left(\hat{a}_- \hat{a}_+ - \frac{1}{2} \right)$$

$[\hat{x}, \hat{p}] = i\hbar$

$$\hat{H} = \hbar\omega \left(\hat{a}_- \hat{a}_+ - \frac{1}{2} \right)$$

Or, using the other order,
we get (left as exercise):

$$\hat{H} = \hbar\omega \left(\hat{a}_+ \hat{a}_- + \frac{1}{2} \right)$$

Theorem: if ψ satisfies $\hat{H}\psi = E\psi$,
then $\hat{H}(\hat{a}_+\psi) = (E + \hbar\omega)(\hat{a}_+\psi)$

$$\begin{aligned} \hat{H}(\hat{a}_+\psi) &= \hbar\omega \left(\hat{a}_+ \hat{a}_- + \frac{1}{2} \right) (\hat{a}_+\psi) = \hbar\omega \left(\hat{a}_+ \hat{a}_- \hat{a}_+ + \frac{1}{2} \hat{a}_+ \right) \psi \\ &= \hbar\omega \hat{a}_+ \left(\hat{a}_- \hat{a}_+ + \frac{1}{2} \right) \psi = \hat{a}_+ \left[\hbar\omega \left(\hat{a}_- \hat{a}_+ + 1 - \frac{1}{2} \right) \psi \right] \\ &= \hat{a}_+ (\hat{H} + \hbar\omega) \psi = \hat{a}_+ (E + \hbar\omega) \psi = (E + \hbar\omega) (\hat{a}_+\psi). \end{aligned}$$

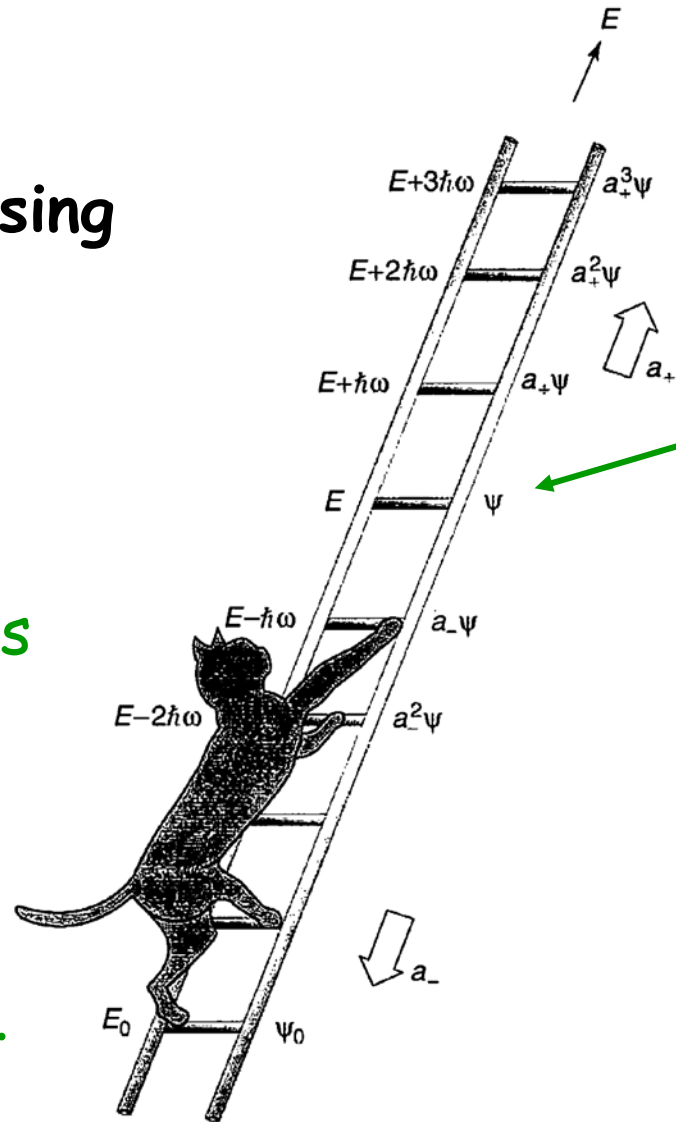
If I know one solution, I know another solution, etc. ...

Another theorem: if ψ satisfies $\hat{H}\psi = E\psi$, then $\hat{H}(\hat{a}_-\psi) = (E - \hbar\omega)(\hat{a}_-\psi)$ (left as exercise)

\hat{a}_+ is the raising operator

Equally spaced energies! This is why it works.

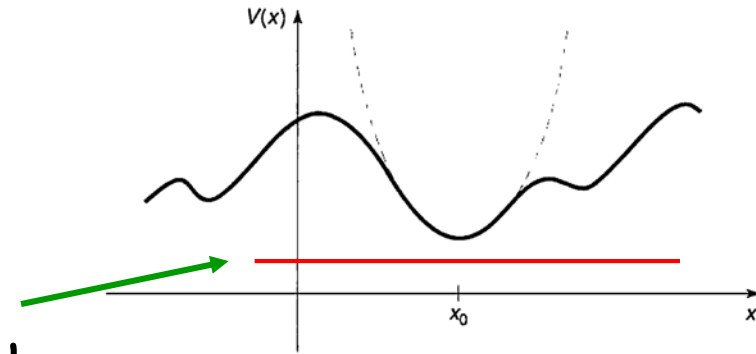
Not general: valid for oscillators only.



If I am given one solution, I get infinite more.

\hat{a}_- is the lowering operator

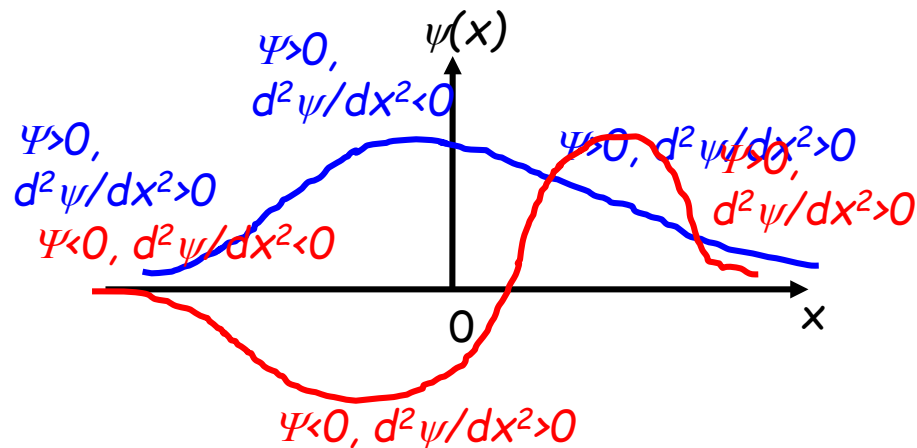
However, there is a problem:
the energy cannot continue
going down!



Theorem: E less than $V(x)$ cannot happen

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + V \psi = E \psi \quad \longrightarrow \quad \frac{d^2 \psi}{dx^2} = \frac{2m}{\hbar^2} \underbrace{[V(x) - E]}_{>0} \psi$$

If this is >0 for all x , as in the **incorrect** sketch above, then ψ and $d^2 \psi/dx^2$ must have same sign. Integrable functions do not satisfy this constraint.



Because the energy cannot
continue going down forever,
the chain down must stop...