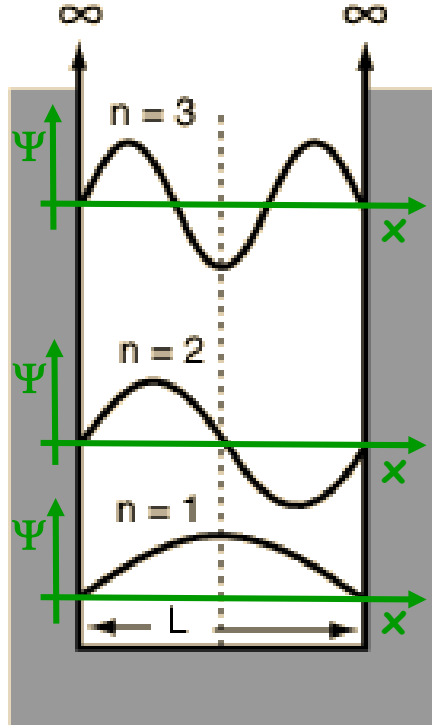


Learn to identify "symmetries"! This one is invariant under reflection

Separating in even-odd is valid only for even $V(x)=V(-x)$ potentials



$x = 0$ at left wall of box.

2 nd excited state	even	2 nodes
1 st excited state	odd	1 node
ground state	even	0 node

Remember there is a $e^{-iE_n t/\hbar}$ multiplying always.

Thus, $\text{Re } \Psi$ and $\text{Im } \Psi$ parts are oscillating with time. But $|\Psi|^2$ is time independent.

Two neat properties of the solutions.

(I) They are **orthonormal**.

$$\int \psi_m(x)^* \psi_n(x) dx = \delta_{mn} \left\{ \begin{array}{l} \text{Kronecker delta} \\ =1 \text{ if } m \text{ equal to } n \\ =0 \text{ if } m \text{ diff from } n \end{array} \right.$$

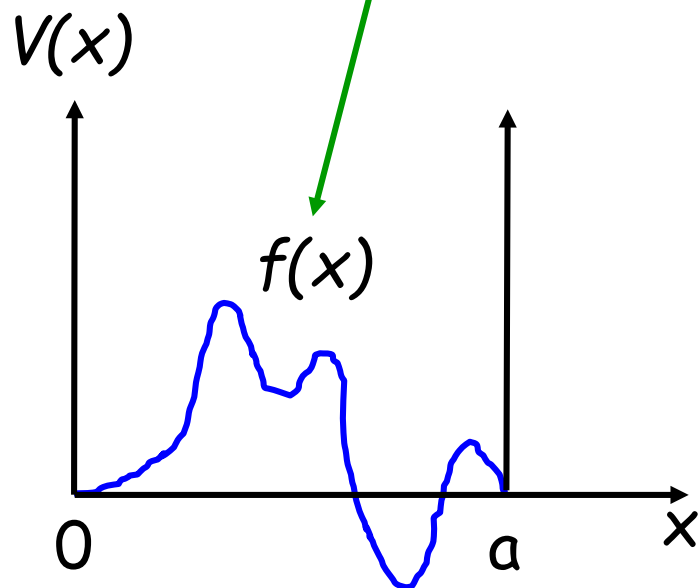
If **$m=n$** this is obvious from normalization done.

If **$m \neq n$** , left as exercise (please check):

$$\begin{aligned} \int \psi_m(x)^* \psi_n(x) dx &= \frac{2}{a} \int_0^a \sin\left(\frac{m\pi}{a}x\right) \sin\left(\frac{n\pi}{a}x\right) dx \\ &= \frac{1}{a} \int_0^a \left[\cos\left(\frac{m-n}{a}\pi x\right) - \cos\left(\frac{m+n}{a}\pi x\right) \right] dx \\ &= \left\{ \frac{1}{(m-n)\pi} \sin\left(\frac{m-n}{a}\pi x\right) - \frac{1}{(m+n)\pi} \sin\left(\frac{m+n}{a}\pi x\right) \right\} \Big|_0^a \\ &= \frac{1}{\pi} \left\{ \frac{\sin[(m-n)\pi]}{(m-n)} - \frac{\sin[(m+n)\pi]}{(m+n)} \right\} = 0. \end{aligned}$$

(II) They are **complete**. This means coefficients c_n can always be found such that **any** wave function inside the square well can be written as

$$f(x) = \sum_{n=1}^{\infty} c_n \psi_n(x) = \underbrace{\sqrt{\frac{2}{a}} \sum_{n=1}^{\infty} c_n \sin\left(\frac{n\pi}{a}x\right)}_{\text{This expression is not surprising. This is just the sine Fourier series of } f(x).}$$

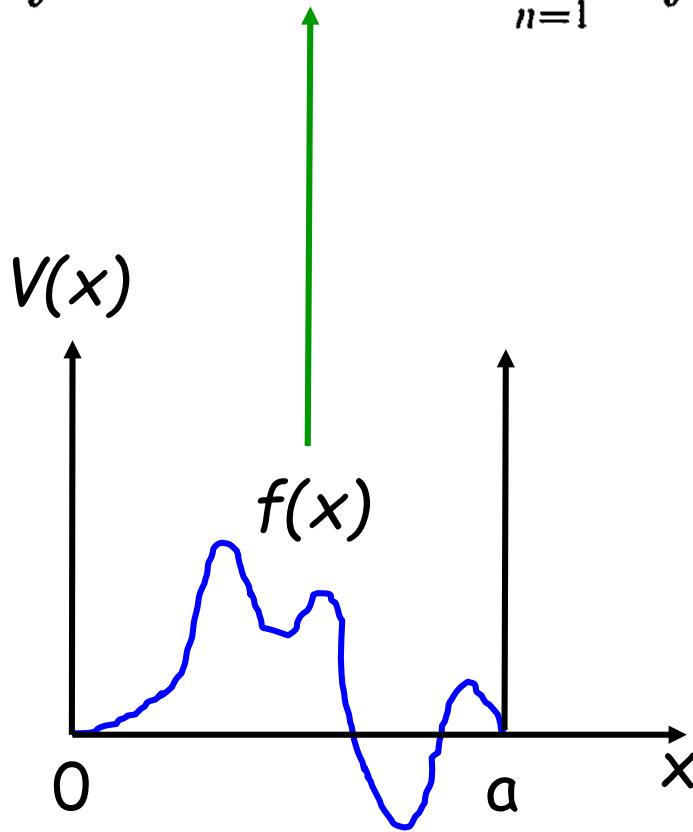


Arbitrary example

- (1) $f(x)$ can be **ANY** function that is 0 outside the well. If not, then it is not acceptable.
- (2) It could be discontinuous inside the well.
- (3) $f(x)$ does **NOT** have a sharp energy but an average energy.

How do we find the coefficients?

$$\int \psi_m(x)^* f(x) dx = \sum_{n=1}^{\infty} c_n \int \psi_m(x)^* \psi_n(x) dx = \sum_{n=1}^{\infty} c_n \delta_{mn} = c_m$$

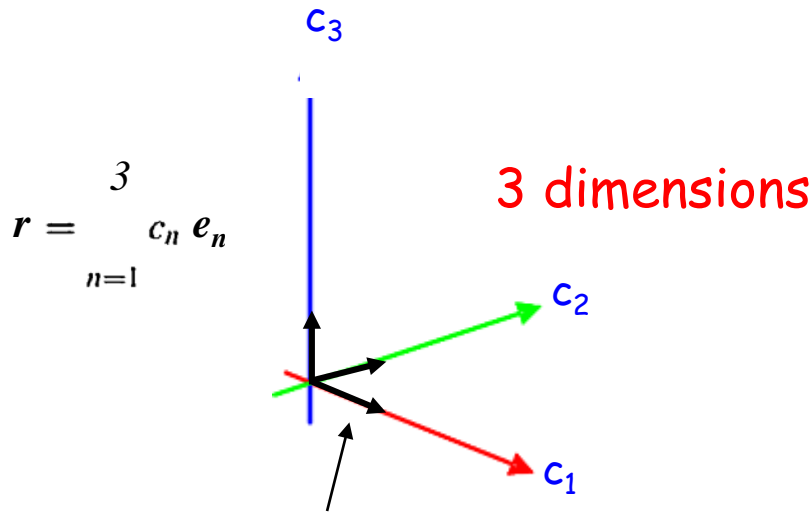


$$c_n = \int \psi_n(x)^* f(x) dx.$$

The integration can be done analytically or numerically.

This page not in book, just FYI.

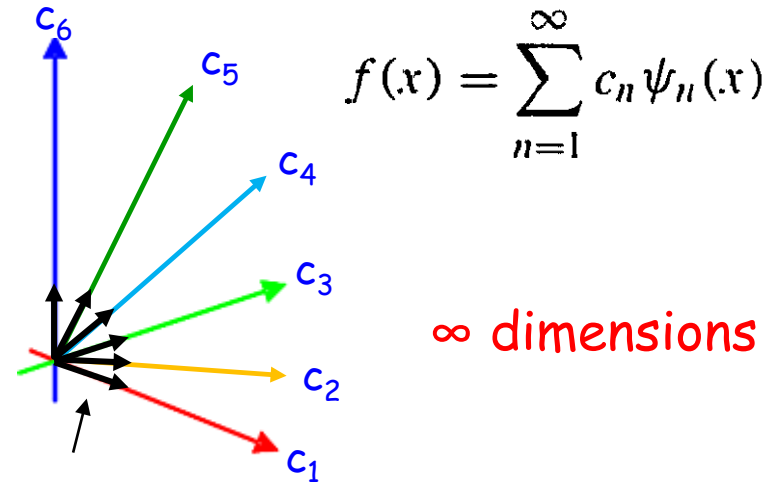
Cartesians axes



Unit vectors are e_1, e_2, e_3 .

Any **vector** can be expanded in the **orthonormal** basis e_1, e_2, e_3 .
E.g. (2,0,-3)

Square-well solutions



"Unit vectors" are $\psi_1, \psi_2, \psi_3, \psi_4, \psi_5, \dots$

Any **wave function** can be expanded in the **orthonormal** basis ψ_n
E.g. (1,2,0,-4,10,0.1,...)

All these properties are not pathological of the square well but very generic.

(3) Returning to page 28 Ch2 book. There are several possible values of E , say E_1, E_2, E_3, \dots , as found in square well example. For each "allowed" energy, there is a solution of time-indep. Sch. Eq. with its "phase factor"

$$\Psi_1(x, t) = \psi_1(x)e^{-iE_1t/\hbar}, \quad \Psi_2(x, t) = \psi_2(x)e^{-iE_2t/\hbar}, \quad \dots$$

Make a linear combination:

$$\Psi(x, t) = \sum_{n=1}^{\infty} c_n \psi_n(x) e^{-iE_n t/\hbar}$$

Statement: any wave function $\Psi(x,t)$ can be written as above. The c_n 's are the same as before i.e. time INDEPENDENT. But the linear combination above is NOT a stationary state i.e. $|\Psi(x,t)|^2$ at fixed x , changes with t .

Not in book:

Is this a solution of the time-dep. Sch. Eq. with $V(x)$?

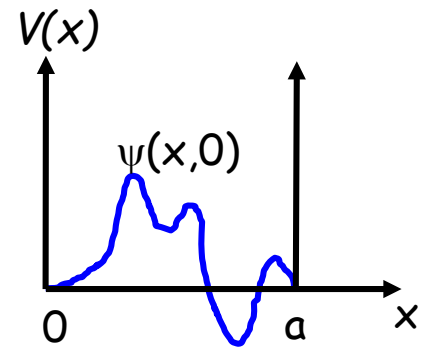
$$\hat{H} \Psi(x, t) = \sum_{n=1}^{\infty} c_n \underbrace{\hat{H} \psi_n(x)}_{E_n \psi_n(x)} e^{-i E_n t / \hbar} \quad (\text{a})$$

$$i \hbar \frac{\partial}{\partial t} \Psi(x, t) = \sum_{n=1}^{\infty} c_n \psi_n(x) \underbrace{i \hbar \frac{\partial}{\partial t} e^{-i E_n t / \hbar}}_{E_n e^{-i E_n t / \hbar}} \quad (\text{b})$$

(a) is equal to (b) \rightarrow The linear combination is solution of the time-dependent Sch. Eq.

Final recipe for $\Psi(x,t)$ in square well:

Given an arbitrary $\Psi(x,0)$ -- that satisfies the BCond -- you want $\Psi(x,t)$.



(1) Find stationary states and energies.

(2) "Somehow" do the integrals for coefficients:

$$c_n = \sqrt{\frac{2}{a}} \int_0^a \sin\left(\frac{n\pi}{a}x\right) \Psi(x,0) dx$$

Done!

$$\Psi(x,t) = \sum_{n=1}^{\infty} c_n \underbrace{\sqrt{\frac{2}{a}} \sin\left(\frac{n\pi}{a}x\right)}_{\Psi_n(x)} \underbrace{e^{-iE_n t/\hbar}}_{e^{-i(n^2\pi^2\hbar/2ma^2)t}}$$

The procedure is general but of course $\Psi_n(x)$ and E_n are diff for diff potentials. Here we use the square well.


Example 2.1, page 29 book. Assume you are given at $t=0$:

Real numbers for simplicity.

$$\Psi(x, 0) = c_1 \psi_1(x) + c_2 \psi_2(x)$$

Stationary states $\sin(n\pi x/a)$. NOT any arbitrary function like, say, $e^{-|x|}$

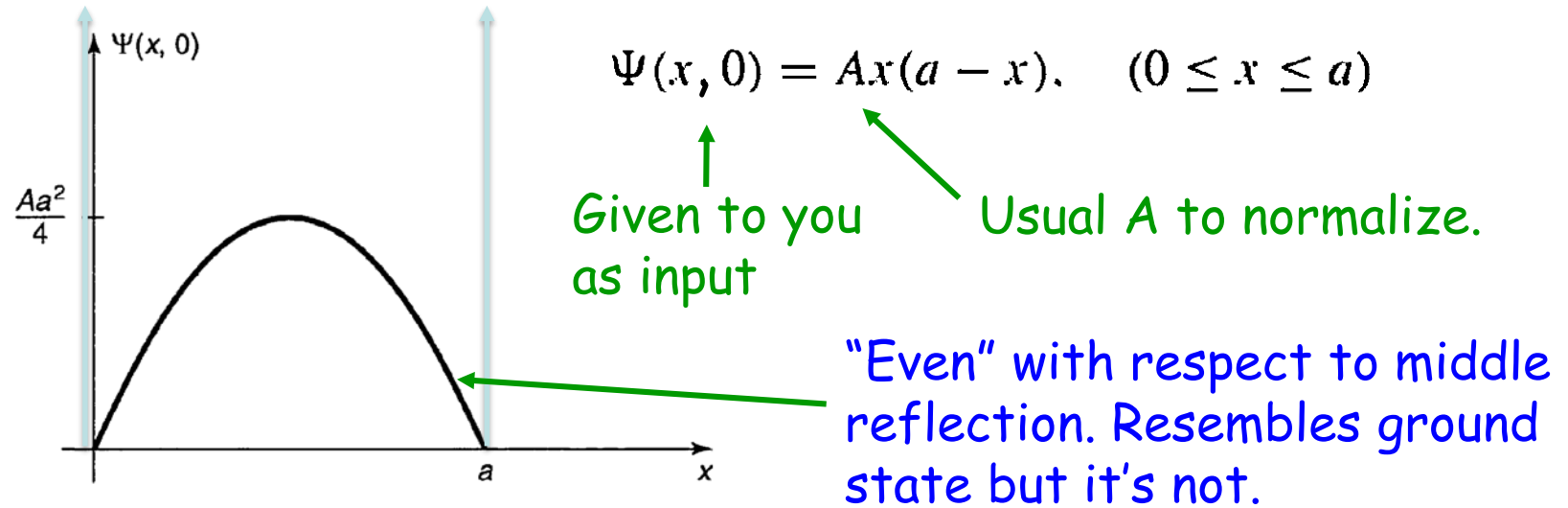
IMPORTANT: Time dependence is now trivially obtained! Eqs 2.16-2.17


$$\Psi(x, t) = c_1 \psi_1(x) e^{-iE_1 t/\hbar} + c_2 \psi_2(x) e^{-iE_2 t/\hbar}$$

$$\begin{aligned} |\Psi(x, t)|^2 &= (c_1 \psi_1 e^{iE_1 t/\hbar} + c_2 \psi_2 e^{iE_2 t/\hbar})(c_1 \psi_1 e^{-iE_1 t/\hbar} + c_2 \psi_2 e^{-iE_2 t/\hbar}) \\ &= c_1^2 \psi_1^2 + c_2^2 \psi_2^2 + 2c_1 c_2 \psi_1 \psi_2 \cos[(E_2 - E_1)t/\hbar]. \end{aligned}$$

The prob. density is now time **dependent** even if stationary states are combined ("quantum beat").

Example 2.2: how to apply the recipe (typical exam problem)



Normalize first:

$$1 = \int_0^a |\Psi(x, 0)|^2 dx = |A|^2 \int_0^a x^2 (a - x)^2 dx \longrightarrow A = \sqrt{\frac{30}{a^5}}$$

Find coefficients c_n :

$$c_n = \sqrt{\frac{2}{a}} \int_0^a \sin\left(\frac{n\pi}{a}x\right) \underbrace{\sqrt{\frac{30}{a^5}} x(a-x) dx}_{\Psi(x,0)} = 8\sqrt{15}/(n\pi)^3$$

n odd only nonzero

$\Psi_n(x)$ $\Psi(x,0)$

See integration process in book. Even n (i.e. odd functions) gives 0 by symmetry.

Then, you can write the final answer:

$$\Psi(x, t) = \sqrt{\frac{30}{a}} \left(\frac{2}{\pi}\right)^3 \sum_{n=1,3,5\dots} \frac{1}{n^3} \sin\left(\frac{n\pi}{a}x\right) e^{-in^2\pi^2\hbar t/2ma^2}$$

Often forgotten by students! It contains the time dependence!

$\underbrace{\hspace{10em}}_{\rightarrow e^{-iE_n t/\hbar}}$