# Sec. 5.1.3. Now all together: the space dependence and the spin (still in 1D for simplicity)

$$\psi_{S}(\mathbf{r}_{1},\mathbf{r}_{2}) = \frac{1}{\sqrt{2}} [\psi_{a}(x_{1})\psi_{b}(x_{2}) + \psi_{b}(x_{1})\psi_{a}(x_{2})]$$

$$\psi_{AS}(\mathbf{r}_{1},\mathbf{r}_{2}) = \frac{1}{\sqrt{2}} [\psi_{a}(x_{1})\psi_{b}(x_{2}) - \psi_{b}(x_{1})\psi_{a}(x_{2})]$$

$$\chi_{S}(\mathbf{S}_{1},\mathbf{S}_{2}) = \frac{1}{\sqrt{2}} (\uparrow \downarrow_{2} + \downarrow \uparrow_{1}) \qquad \chi_{AS}(\mathbf{S}_{1},\mathbf{S}_{2}) = \frac{1}{\sqrt{2}} (\uparrow \downarrow_{2} - \downarrow \uparrow_{1})$$

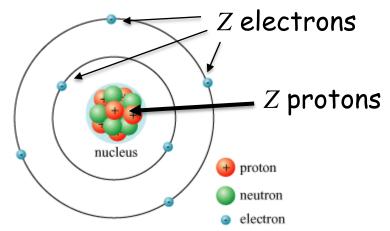
For 2 electrons (i.e. <u>fermions</u>) in two different states:

$$\Psi = \psi_{S}(\mathbf{r}_{1},\mathbf{r}_{2}) \chi_{AS}(\mathbf{S}_{1},\mathbf{S}_{2}) \text{ OK ! } \Psi = \psi_{S}(\mathbf{r}_{1},\mathbf{r}_{2}) \chi_{S}(\mathbf{S}_{1},\mathbf{S}_{2}) \text{ NO}$$

$$\Psi = \psi_{AS}(\mathbf{r}_{1},\mathbf{r}_{2}) \chi_{S}(\mathbf{S}_{1},\mathbf{S}_{2}) \text{ OK ! } \Psi = \psi_{AS}(\mathbf{r}_{1},\mathbf{r}_{2}) \chi_{AS}(\mathbf{S}_{1},\mathbf{S}_{2}) \text{ NO}$$

If a=b, then only  $\Psi = \psi_S(\mathbf{r}_1,\mathbf{r}_2) \chi_{AS}(\mathbf{S}_1,\mathbf{S}_2)$  is possible (Pauli)

#### 5.2 Atoms



#### The Hamiltonian is:

$$H = \sum_{j=1}^{Z} \left\{ -\frac{\hbar^2}{2m} \nabla_j^2 - \left(\frac{1}{4\pi\epsilon_0}\right) \underbrace{\frac{Z \theta^2}{r_j}} \right\} + \frac{1}{2} \left(\frac{1}{4\pi\epsilon_0}\right) \underbrace{\sum_{j\neq k}^{Z} \frac{e^2}{|\mathbf{r}_j - \mathbf{r}_k|}}_{\text{one-body e-p attraction}} \right.$$
 e-e repulsion "the easy part" e-e refulsion "the difficult part"

## Solving exactly is impossible $\otimes$

$$H \ \psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_Z) \chi(\mathbf{s}_1, \mathbf{s}_2, \dots, \mathbf{s}_Z) = \begin{cases} \text{No r dependence} \\ \text{here} \end{cases}$$

$$= E \ \psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_Z) \chi(\mathbf{s}_1, \mathbf{s}_2, \dots, \mathbf{s}_Z)$$

In writing the Sch Eq we assumed that the spins may be coupled among themselves and/or with a uniform magnetic field, but the spin orientation does not depend on position.

Because electrons are fermions, the entire wave function must be antisymmetric.

## 5.2.1 Helium (Z=2)

$$H = \left\{ -\frac{\hbar^2}{2m} \nabla_1^2 - \frac{1}{4\pi \epsilon_0} \frac{2e^2}{r_1} \right\} + \left\{ -\frac{\hbar^2}{2m} \nabla_2^2 - \frac{1}{4\pi \epsilon_0} \frac{2e^2}{r_2} \right\} + \frac{1}{4\pi \epsilon_0} \frac{e^2}{|\mathbf{r}_1 - \mathbf{r}_2|}$$

First, neglect the e-e repulsion (on page 322, Ch 8, we will improve on this)

The space-like portion of the wave function in general will be (before symmetrization):

$$\psi(\mathbf{r}_1, \mathbf{r}_2) = \psi_{nlm}(\mathbf{r}_1)\psi_{n'l'm'}(\mathbf{r}_2)$$

$$E = 4(E_n + E_{n'})$$

For **ground state**, we place both electrons at n=1, l=0, m=0.

$$E_n = -\left[\frac{m}{2\hbar^2} \left(\frac{2e^2}{4\pi\epsilon_0}\right)^2\right] \frac{1}{n^2}$$

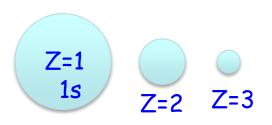
$$\psi_0(\mathbf{r}_1, \mathbf{r}_2) = \psi_{100}(\mathbf{r}_1)\psi_{100}(\mathbf{r}_2) = \frac{8}{\pi a^3}e^{-2(r_1+r_2)/a}$$

$$E_0 = 8(-13.6 \text{ eV}) = -109 \text{ eV}$$
  
 $4+4 = 2^2+2^2$ 

Rapidly with increasing Z, big energies are induced!

-109 eV with Z=2 vs -13.6 eV with Z=1

Bohr radius reduced by factor 2; in general by a factor Z.



### Some consequences of AS vs S:

Because the full wave function has a "space portion" and a "spin portion", the first excited states of He have two possibilities

$$\Psi_{2e} = \psi_{S}(\mathbf{r}_{1},\mathbf{r}_{2}) \chi_{AS}(\mathbf{S}_{1},\mathbf{S}_{2})$$

$$\Psi_{2e} = \psi_{AS}(\mathbf{r}_{1},\mathbf{r}_{2}) \chi_{S}(\mathbf{S}_{1},\mathbf{S}_{2}) \leftarrow \text{First excited state}$$
is triplet S=1

then, with all other things equal, the thus-far ignored e-e repulsion, that has nothing to do with spin, prefers the AS space portion because electrons are further apart than in the S space portion (see page 211). Confirmed experimentally that the first excited state has spin 1.

Now repeating using the Chemistry class cartoons: The space portion is symmetric, thus the spin portion must be antisymmetric.

$$\psi = \psi(\mathbf{r}_1, \mathbf{r}_2) \chi(\mathbf{s}_1, \mathbf{s}_2) = \frac{8}{\pi a^3} e^{-2(r_1+r_2)/a} \frac{1}{\sqrt{2}} \left( \uparrow_1 \downarrow_2 - \downarrow_1 \uparrow_2 \right)$$

The ground state cartoon n=3  $\frac{1}{n=2}$   $\frac{1}{n=1}$   $\frac{1}{n=1}$   $\frac{1}{n=1}$   $\frac{1}{n=1}$ 

#### This portion is repeated: Excited states? Two options

•••

$$\Psi_{\text{singlet}} = \psi_{\text{S}}(\mathbf{r}_{1}, \mathbf{r}_{2}) \chi_{\text{AS}}(\mathbf{S}_{1}, \mathbf{S}_{2}) = \frac{1}{\sqrt{2}} \left[ \psi_{100}(\mathbf{r}_{1}) \psi_{200}(\mathbf{r}_{2}) + \psi_{200}(\mathbf{r}_{1}) \psi_{100}(\mathbf{r}_{2}) \right] \frac{1}{\sqrt{2}} \left( \uparrow \downarrow_{1} - \downarrow_{1} \uparrow_{2} \right)$$

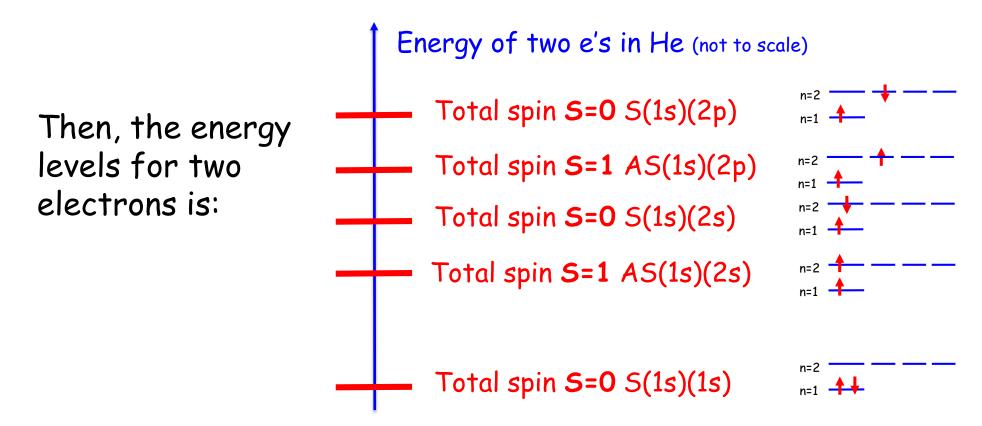
$$\Psi_{\text{triplet}} = \psi_{AS}(\mathbf{r}_{1},\mathbf{r}_{2}) \chi_{S}(\mathbf{S}_{1},\mathbf{S}_{2}) =$$

$$= \frac{1}{\sqrt{2}} \left[ \psi_{100}(\mathbf{r}_{1}) \psi_{200}(\mathbf{r}_{2}) - \psi_{200}(\mathbf{r}_{1}) \psi_{100}(\mathbf{r}_{2}) \right] \frac{1}{\sqrt{2}} \left( \uparrow_{1} \downarrow_{2} + \downarrow_{1} \uparrow_{2} \right)$$

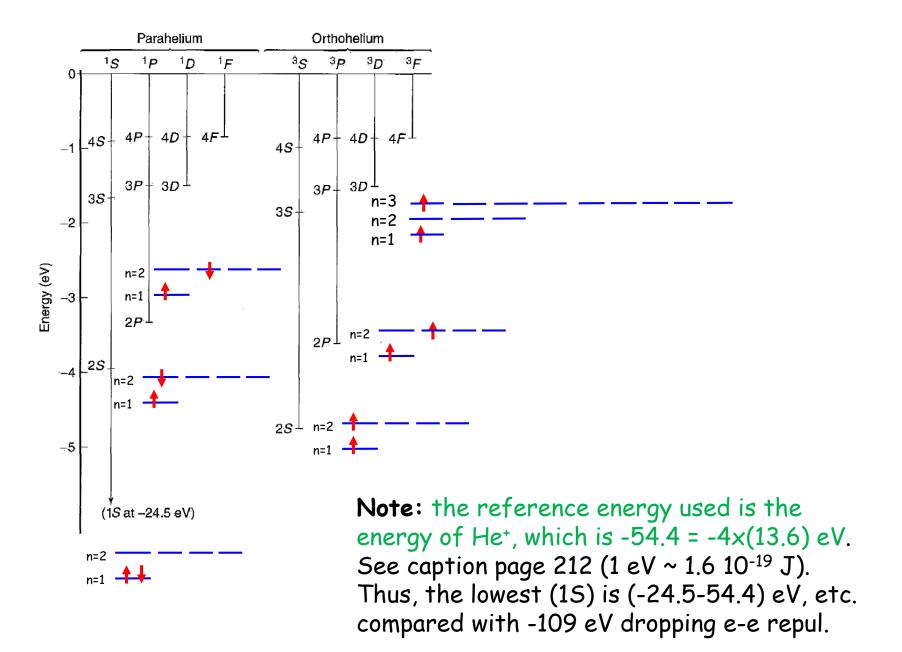
If e-e neglected, then singlet and triplet are degenerate

If e-e is brought back, at least qualitatively, then the degeneracy must be broken

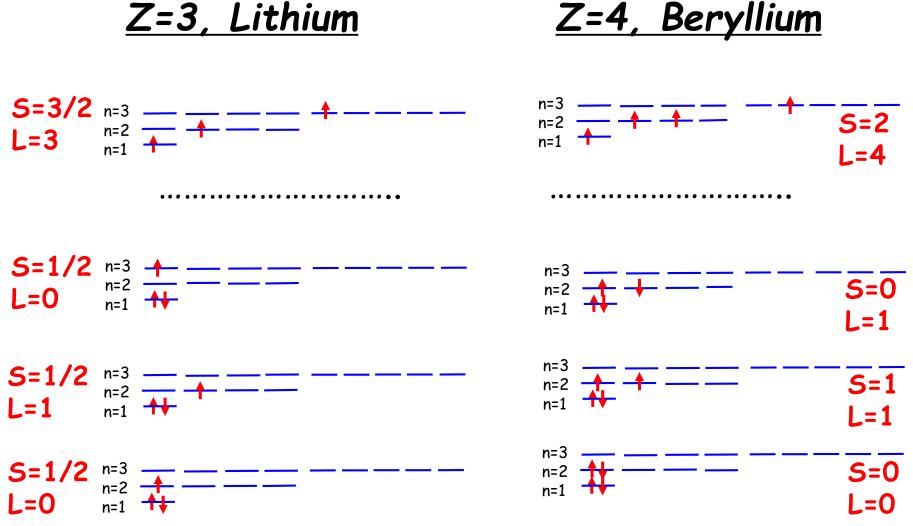
Due to the effective "exchange forces", the AS space-like combination keeps the two electrons a bit further apart ... (for AS sector the exchange force is "repulsive"; for S sector is "attractive")



#### Real numbers from book



#### Not in book, excited states



The excited states become complicated very fast!