

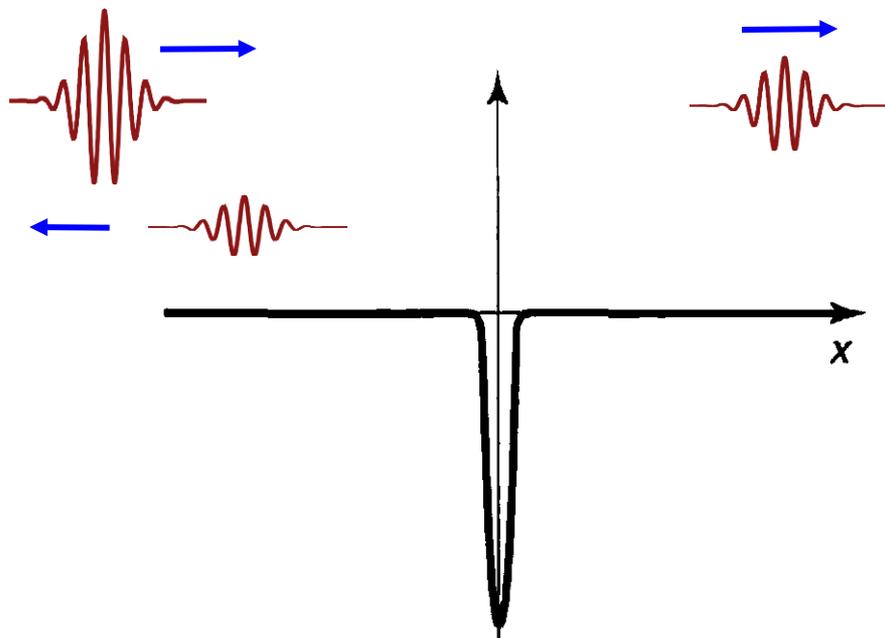
## 4.4: Spin of electrons

This topic is not part of Test 3, but it may be part of Test 1 of QM-412

A classical rigid body, like a planet, can have two kinds of angular momenta: (1)  $L$ , the **orbital** one associated with the center of mass, like Earth around the sun, and (2)  $S$ , the **spin**, like Earth rotating daily about an axis.

In quantum mechanics we already discussed the orbital component  $L$  (related with the electron around the nucleus).

In QM, we also have a spin  $S$  for the electron but ... the electron to the best of our accuracy is a **POINT**, thus cannot rotate.



However, once we measure the location and find the particle at position  $x_0$ , then that "sigma" width is gone. The particle is perfectly at  $x_0$ . **At that moment what radius it has?**

In wave packets as shown, the "finite size" due to "sigma" is the finite size of the wave function related to the probability of finding the particle.

It seems that the radius is smaller than  $10^{-18}$  m according to experiments (see email of Tuesday). Radius of nucleus is  $10^{-15}$  m.

It is a fact of Nature, that elementary particles such as an electron carry an **intrinsic spin angular momentum  $\mathbf{S}$** .

Because the electron is a point, we cannot use the classical formulas  $\mathbf{S} = I\boldsymbol{\omega}$  or  $\mathbf{L} = \mathbf{r} \times \mathbf{p}$

To describe the **intrinsic spin** the math has to be “analogous” to that of  $\mathbf{L}$ . Let us start with the commutators:

$$[L_x, L_y] = i\hbar L_z; \quad [L_y, L_z] = i\hbar L_x; \quad [L_z, L_x] = i\hbar L_y$$

becomes ...

$$[S_x, S_y] = i\hbar S_z, \quad [S_y, S_z] = i\hbar S_x, \quad [S_z, S_x] = i\hbar S_y$$

The eigenfunctions are more "abstract" ...

First, let us switch to the **Ch. 3 notation** using an abstract Hilbert space notation:

$$L^2 f_l^m = \hbar^2 l(l+1) f_l^m; \quad L_z f_l^m = \hbar m f_l^m \quad \text{to}$$

$$L^2 |l m_l\rangle = \hbar^2 l(l+1) |l m_l\rangle; \quad L_z |l m_l\rangle = \hbar m_l |l m_l\rangle$$

For  $L^2$  and  $L_z$  using  $Y_l^m(\theta, \phi)$  or  $|l m_l\rangle$  is the SAME.

But for the intrinsic spin the Ch. 3 notation is the **ONLY** way because there are **no angles** to use.

Because in lecture Nov. 20 we arrived all the way to the eigenvalues by only using the commutators, then we simply **repeat** the operation **line by line** and find:

$$S^2 |s m_s\rangle = \hbar^2 s(s+1) |s m_s\rangle; \quad S_z |s m_s\rangle = \hbar m_s |s m_s\rangle$$

$$s = 0, \frac{1}{2}, 1, \frac{3}{2}, \dots; \quad m = -s, -s + 1, \dots, s - 1, s$$

The **spin of each type of particle is FIXED**, not like the orbital angular momentum that you can change by emission or absorption of energy.

### 4.4.1: Spin $\frac{1}{2}$ (electrons, quarks)

Use  $\mathbf{S}^2 |s m_s\rangle = \hbar^2 s(s+1) |s m_s\rangle$ ;  $S_z |s m_s\rangle = \hbar m_s |s m_s\rangle$

Specialize for  $s=1/2$ . Then, there are only two states, which in abstract form (no angles  $\theta$  and  $\phi$ !) are:

$$|\frac{1}{2} \frac{1}{2}\rangle \text{ and } |\frac{1}{2} -\frac{1}{2}\rangle$$

We call them spin "up" or  $\uparrow$  and spin "down" or  $\downarrow$ .

There is another, still abstract, way to represent spins up and down. It is using so-called "spinors"

$$\chi_+ = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \chi_- = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

We can combine the "up" and "down" linearly at will.  
So the spin could point "sideways" for instance.

$$\chi = \begin{pmatrix} a \\ b \end{pmatrix} = a\chi_+ + b\chi_-.$$

If we use spinors for the states, then what do we use for the operators such as  $L^2$ ? Certainly we cannot use derivatives of angles. There are no angles!

From the two equations ...

$$\mathbf{S}^2 \left| \frac{1}{2} \frac{1}{2} \right\rangle = \hbar^2 \frac{1}{2} \left( \frac{1}{2} + 1 \right) \left| \frac{1}{2} \frac{1}{2} \right\rangle$$

$$\mathbf{S}^2 \left| \frac{1}{2} -\frac{1}{2} \right\rangle = \hbar^2 \frac{1}{2} \left( \frac{1}{2} + 1 \right) \left| \frac{1}{2} -\frac{1}{2} \right\rangle$$

... it can be deduced (see book, easy) that:

$$\mathbf{S}^2 = \frac{3}{4} \hbar^2 \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

From the other two equations ...

$$S_z \left| \frac{1}{2} \frac{1}{2} \right\rangle = \frac{1}{2} \hbar \left| \frac{1}{2} \frac{1}{2} \right\rangle$$

$$S_z \left| \frac{1}{2} -\frac{1}{2} \right\rangle = -\frac{1}{2} \hbar \left| \frac{1}{2} -\frac{1}{2} \right\rangle$$

... it can be deduced (see book, easy) that:

$$\mathbf{S}_z = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Using the  $S_+$  and  $S_-$  operators it can be shown (book):

$$\mathbf{S}_x = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \mathbf{S}_y = \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

Dropping the  $\hbar/2$  factor defines  
the famous **Pauli matrices**:

$$\sigma_x \equiv \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y \equiv \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_z \equiv \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Returning to the general combination:

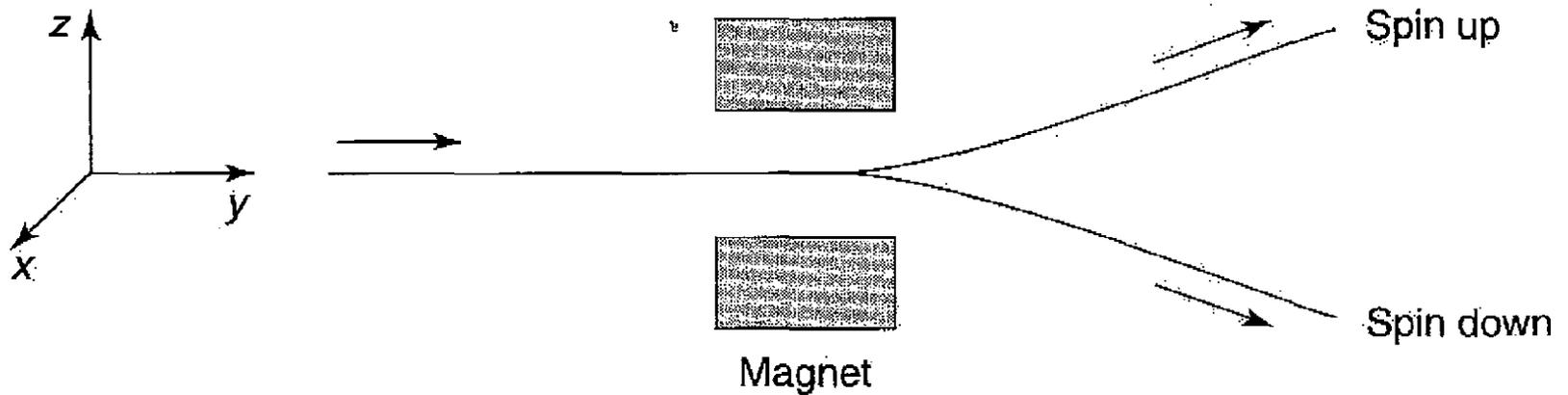
$$\chi = \begin{pmatrix} a \\ b \end{pmatrix} = a\chi_+ + b\chi_-.$$

You have to normalize i.e.  $|a|^2 + |b|^2 = 1$

$|a|^2$  is the probability of measuring spin up.

$|b|^2$  is the probability of measuring spin down.

How do you measure a spin? The same as any magnetic moment, like the orbital "l". You introduce the particle in a magnetic field. Also there is something called the **Stern-Gerlach experiment**:



Returning, again!, to the general combination:

$$\chi = \begin{pmatrix} a \\ b \end{pmatrix} = a\chi_+ + b\chi_-.$$

What is the probability that the spin points say along the *positive x axis*?

To answer this question, first you have to diagonalize the 2x2 Pauli matrix "x". Turns out the "eigenspinors" are:

$$\chi_+^{(x)} = \begin{pmatrix} \frac{1}{\sqrt{2}} \\ 1 \\ \frac{1}{\sqrt{2}} \end{pmatrix}, \left( \text{eigenvalue} + \frac{\hbar}{2} \right); \quad \chi_-^{(x)} = \begin{pmatrix} \frac{1}{\sqrt{2}} \\ -1 \\ \frac{1}{\sqrt{2}} \end{pmatrix}, \left( \text{eigenvalue} - \frac{\hbar}{2} \right)$$

The general combination ...

$$\chi = \begin{pmatrix} a \\ b \end{pmatrix} = a\chi_+ + b\chi_-.$$

... can now be written as:

$$\chi = \left( \frac{a+b}{\sqrt{2}} \right) \chi_+^{(x)} + \left( \frac{a-b}{\sqrt{2}} \right) \chi_-^{(x)}$$

$$(1/2)|a+b|^2$$

is the probability of measuring spin up along x.

$$(1/2)|a-b|^2$$

is the probability of measuring spin down along x.