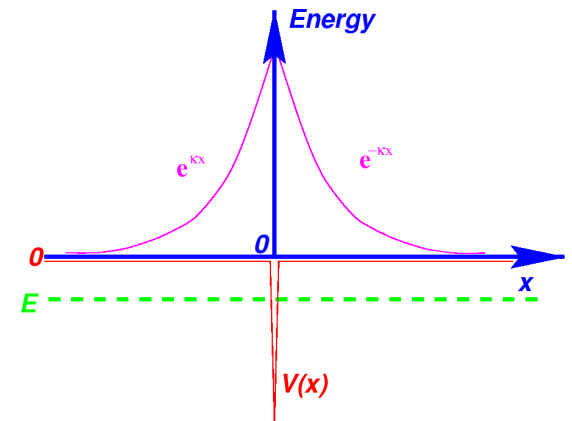


## Example 3.4 book:

Consider a particle located in the (only) bound state of the  $\delta$ -function potential. The wave function is:

$$\Psi(x, t) = \frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2} e^{-iEt/\hbar}$$



**Typical question:** what is the probability of measuring a momentum greater than  $p_0 = m\alpha/\hbar^*$ ? We need to calculate  $|c(p)|^2$  (see next page).

\* Can you confirm that the units are those of momentum?

Reminder of previous lecture:  
eigenfunction of momentum operator  $\hat{p}$ .  
We need an eigenfunction  $f_p(x)$  such that

The diagram shows the momentum eigenvalue equation:  $\hat{p} \text{ operator} \rightarrow \left( \frac{\hbar}{i} \frac{d}{dx} \right) f_p(x) = p f_p(x)$ . Annotations include: a blue arrow pointing from the text " $\hat{p}$  operator" to the differential operator  $\frac{\hbar}{i} \frac{d}{dx}$ ; a blue circle around the differential operator; a blue circle around the function  $f_p(x)$  on the left side of the equation, with a blue arrow pointing from the text "eigenfunction" to it; a blue circle around the eigenvalue  $p$ , with a blue arrow pointing from the text "eigenvalue" to it.

Solution is very easy (but  
normalization is complicated):

$$f_p(x) = \frac{1}{\sqrt{2\pi\hbar}} e^{ipx/\hbar}$$

$$\Psi(x, t) = \frac{\sqrt{m\alpha}}{\hbar} e^{-m\alpha|x|/\hbar^2} e^{-iEt/\hbar}$$

**Repeating:** What is the probability of measuring a momentum greater than  $p_0 = m\alpha/\hbar$  ?

We need to calculate  $|c(p)|^2$ .

$$\begin{aligned} c(p) &= \langle f_p | \Psi \rangle = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} e^{-ipx/\hbar} \Psi(x, t) dx = \\ &= \frac{1}{\sqrt{2\pi\hbar}} \frac{\sqrt{m\alpha}}{\hbar} e^{-iEt/\hbar} \int_{-\infty}^{\infty} e^{-ipx/\hbar} e^{-m\alpha|x|/\hbar^2} dx = \boxed{\sqrt{\frac{2}{\pi}} \frac{p_0^{3/2} e^{-iEt/\hbar}}{p^2 + p_0^2}} \end{aligned}$$

↑  
integral given

The final answer is:

Integral given.  
Answer must be  
a number [0,1]

$$\int_{p_0}^{\infty} |c(p)|^2 dp = \frac{2}{\pi} p_0^3 \int_{p_0}^{\infty} \frac{1}{(p^2 + p_0^2)^2} dp = 0.0908$$

**Just a name:**  $c(p)$ , which can be function of  $t$  but **not**  $x$ , is often called  $\Phi(p,t)$ , the **momentum space wave function**.

A similar problem could be formulated for the **harmonic oscillator** involving Gaussians (problem 3.11, HW8) or the **infinite square well** involving sines, etc., etc.

# Uncertainty Principle (4 pages, prepare for impact)

The standard deviation for any **operator** is

$$\begin{aligned}\sigma_A^2 &= \langle \hat{A}^2 \rangle - \langle \hat{A} \rangle^2 = \langle \Psi | \hat{A}^2 | \Psi \rangle - \langle \Psi | \hat{A} \rangle \langle \hat{A} | \Psi \rangle = \\ &= \langle \Psi | \hat{A}^2 - 2\langle \hat{A} \rangle \hat{A} + \langle \hat{A} \rangle \hat{A} | \Psi \rangle = \langle \Psi | (\hat{A} - \langle \hat{A} \rangle)^2 | \Psi \rangle\end{aligned}$$

If  $\Psi$  normalized to 1  
i.e.  $\langle \Psi | \Psi \rangle = 1$ .

For **Hermitian operator** this can be rewritten as:

$$\begin{aligned}\text{For } \hat{A}: \quad \sigma_A^2 &= \langle (\hat{A} - \langle \hat{A} \rangle) \Psi | (\hat{A} - \langle \hat{A} \rangle) \Psi \rangle = \langle f | f \rangle \quad \text{where} \\ f &\equiv (\hat{A} - \langle \hat{A} \rangle) \Psi\end{aligned}$$

$$\text{For } \hat{B}: \quad \sigma_B^2 = \langle g | g \rangle \quad g \equiv (\hat{B} - \langle \hat{B} \rangle) \Psi$$

Consider the Schwartz inequality:

$$\sigma_A^2 \sigma_B^2 = \langle f|f \rangle \langle g|g \rangle \geq |\langle f|g \rangle|^2$$



$z$

We use now the following property of complex numbers:

$$|z|^2 = [\text{Re}(z)]^2 + [\text{Im}(z)]^2 \geq [\text{Im}(z)]^2 = \left[ \frac{1}{2i} (z - z^*) \right]^2$$

Consider  $z = \langle f|g \rangle$

Then ....

$$\sigma_A^2 \sigma_B^2 \geq \left( \frac{1}{2i} [\langle f|g \rangle - \langle g|f \rangle] \right)^2$$

Then (repeated) ....  $\sigma_A^2 \sigma_B^2 \geq \left( \frac{1}{2i} [\langle f|g \rangle - \langle g|f \rangle] \right)^2$

$\hat{A}$  is Hermitian

$$\begin{aligned} \langle f|g \rangle &= \langle (\hat{A} - \langle \hat{A} \rangle) \Psi | (\hat{B} - \langle \hat{B} \rangle) \Psi \rangle = \langle \Psi | (\hat{A} - \langle \hat{A} \rangle) (\hat{B} - \langle \hat{B} \rangle) \Psi \rangle \\ &= \langle \Psi | (\hat{A} \hat{B} - \hat{A} \langle \hat{B} \rangle - \hat{B} \langle \hat{A} \rangle + \langle \hat{A} \rangle \langle \hat{B} \rangle) \Psi \rangle \\ &= \langle \Psi | \hat{A} \hat{B} \Psi \rangle - \langle \hat{B} \rangle \langle \Psi | \hat{A} \Psi \rangle - \langle \hat{A} \rangle \langle \Psi | \hat{B} \Psi \rangle + \langle \hat{A} \rangle \langle \hat{B} \rangle \langle \Psi | \Psi \rangle \\ &= \langle \hat{A} \hat{B} \rangle - \langle \hat{B} \rangle \langle \hat{A} \rangle - \langle \hat{A} \rangle \langle \hat{B} \rangle + \langle \hat{A} \rangle \langle \hat{B} \rangle \\ &= \langle \hat{A} \hat{B} \rangle - \langle \hat{A} \rangle \langle \hat{B} \rangle. \end{aligned}$$

Make sure you understand every step

$$\langle g|f \rangle = \langle \hat{B} \hat{A} \rangle - \langle \hat{A} \rangle \langle \hat{B} \rangle \quad \text{Left as exercise}$$

In summary:  $\langle f|g \rangle - \langle g|f \rangle = \langle \hat{A} \hat{B} \rangle - \langle \hat{B} \hat{A} \rangle = \langle [\hat{A}, \hat{B}] \rangle$

$$[\hat{A}, \hat{B}] \equiv \hat{A} \hat{B} - \hat{B} \hat{A}$$

## Generalized uncertainty principle:

$$\sigma_A^2 \sigma_B^2 \geq \left( \frac{1}{2i} \langle [\hat{A}, \hat{B}] \rangle \right)^2$$

Assumes  $\langle \dots \rangle$  is in a normalized to 1 state, and both operators Hermitian.

As special case, if  $\hat{A} = \hat{x}$  and  $\hat{B} = \hat{p}$ , then  $[\hat{x}, \hat{p}] = i\hbar$

$$\sigma_x^2 \sigma_p^2 \geq \left( \frac{1}{2i} i\hbar \right)^2 = \left( \frac{\hbar}{2} \right)^2 \quad \sigma_x \sigma_p \geq \frac{\hbar}{2}$$

This was the last item of Ch. 3 for us