

Placing all together: the space and the spin (1D for simplicity)

$$\psi_S(\mathbf{r}_1, \mathbf{r}_2) = \frac{1}{\sqrt{2}} [\psi_a(x_1)\psi_b(x_2) + \psi_b(x_1)\psi_a(x_2)]$$

$$\psi_{AS}(\mathbf{r}_1, \mathbf{r}_2) = \frac{1}{\sqrt{2}} [\psi_a(x_1)\psi_b(x_2) - \psi_b(x_1)\psi_a(x_2)]$$

$$\chi_S(\mathbf{s}_1, \mathbf{s}_2) = \frac{1}{\sqrt{2}} \left(\begin{array}{c} \uparrow \\ 1 \end{array} \begin{array}{c} \downarrow \\ 2 \end{array} + \begin{array}{c} \downarrow \\ 1 \end{array} \begin{array}{c} \uparrow \\ 2 \end{array} \right)$$

$$\chi_{AS}(\mathbf{s}_1, \mathbf{s}_2) = \frac{1}{\sqrt{2}} \left(\begin{array}{c} \uparrow \\ 1 \end{array} \begin{array}{c} \downarrow \\ 2 \end{array} - \begin{array}{c} \downarrow \\ 1 \end{array} \begin{array}{c} \uparrow \\ 2 \end{array} \right)$$

For 2 electrons (i.e. fermions):

$$\Psi = \psi_S(\mathbf{r}_1, \mathbf{r}_2) \chi_{AS}(\mathbf{s}_1, \mathbf{s}_2) \text{ OK!}$$

$$\Psi = \psi_{AS}(\mathbf{r}_1, \mathbf{r}_2) \chi_S(\mathbf{s}_1, \mathbf{s}_2) \text{ OK!}$$

$$\Psi = \psi_S(\mathbf{r}_1, \mathbf{r}_2) \chi_S(\mathbf{s}_1, \mathbf{s}_2) \text{ NO}$$

$$\Psi = \psi_{AS}(\mathbf{r}_1, \mathbf{r}_2) \chi_{AS}(\mathbf{s}_1, \mathbf{s}_2) \text{ NO}$$

Some consequences of AS vs S:

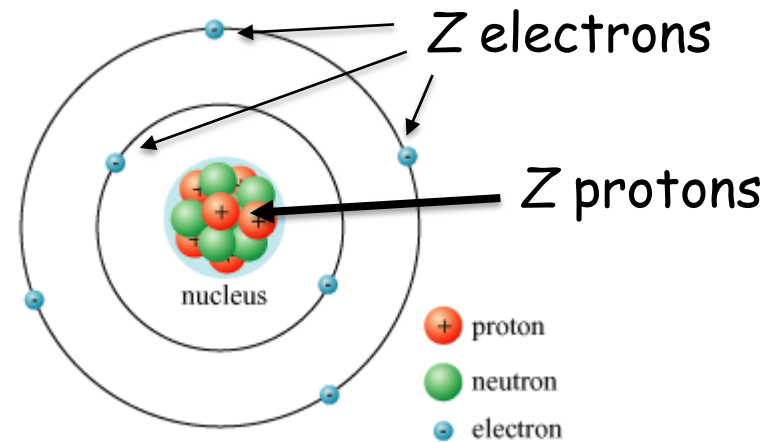
Because the full wave function has a "space portion" and a "spin portion", as shown next for 2e in He, the two possibilities are

$$\Psi_{2e} = \psi_S(\mathbf{r}_1, \mathbf{r}_2) \chi_{AS}(\mathbf{s}_1, \mathbf{s}_2)$$

$$\Psi_{2e} = \psi_{AS}(\mathbf{r}_1, \mathbf{r}_2) \chi_S(\mathbf{s}_1, \mathbf{s}_2)$$

then, all other things equal, the e-e repulsion, that has nothing to do with spin, prefers the AS space portion because electrons are further apart than in S space portion.

5.2 Atoms



The Hamiltonian is:

$$H = \sum_{j=1}^Z \left\{ -\frac{\hbar^2}{2m} \nabla_j^2 - \left(\frac{1}{4\pi\epsilon_0} \right) \frac{Ze^2}{r_j} \right\} + \frac{1}{2} \left(\frac{1}{4\pi\epsilon_0} \right) \sum_{j \neq k}^Z \frac{e^2}{|\mathbf{r}_j - \mathbf{r}_k|}$$

one-body e-p attraction
"the easy part"

e-e repulsion
"the difficult part"

Big challenge:

$$H \psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_Z) \chi(\mathbf{s}_1, \mathbf{s}_2, \dots, \mathbf{s}_Z) =$$

$$= E \psi(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_Z) \chi(\mathbf{s}_1, \mathbf{s}_2, \dots, \mathbf{s}_Z)$$

No \mathbf{r} dependence
here

In writing the Sch Eq we **assumed** that the spins may be coupled among themselves and/or with a uniform magnetic field, but **the spins do not depend on position**.

Because electrons are fermions, the entire wave function must be antisymmetric.

5.2.1 Helium ($Z=2$)

$$H = \left\{ -\frac{\hbar^2}{2m} \nabla_1^2 - \frac{1}{4\pi\epsilon_0} \frac{2e^2}{r_1} \right\} + \left\{ -\frac{\hbar^2}{2m} \nabla_2^2 - \frac{1}{4\pi\epsilon_0} \frac{2e^2}{r_2} \right\} + \frac{1}{4\pi\epsilon_0} \frac{e^2}{|\mathbf{r}_1 - \mathbf{r}_2|}$$

First neglect the e-e repulsion
(on page 299, Ch 7, we will improve on this)

The space-like portion of the wave function in general will be (before symmetrization):

$$\psi(\mathbf{r}_1, \mathbf{r}_2) = \psi_{nlm}(\mathbf{r}_1) \psi_{n'l'm'}(\mathbf{r}_2)$$

$$E = 4(E_n + E_{n'})$$

For ground state, we place both electrons at $n=1, l=0, m=0$.

$$E_n = - \left[\frac{m}{2\hbar^2} \left(\frac{2e^2}{4\pi\epsilon_0} \right)^2 \right] \frac{1}{n^2}$$

$$\psi_0(\mathbf{r}_1, \mathbf{r}_2) = \psi_{100}(\mathbf{r}_1) \psi_{100}(\mathbf{r}_2) = \frac{8}{\pi a^3} e^{-2(r_1+r_2)/a}$$

$$E_0 = 8(-13.6 \text{ eV}) = -109 \text{ eV}$$

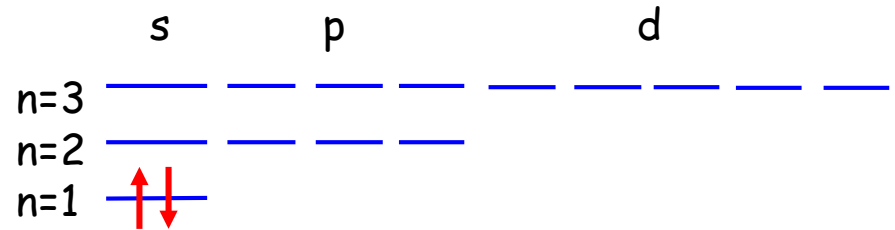
$$\hookrightarrow 4+4 = 2^2+2^2$$

Rapidly inducing big energies!

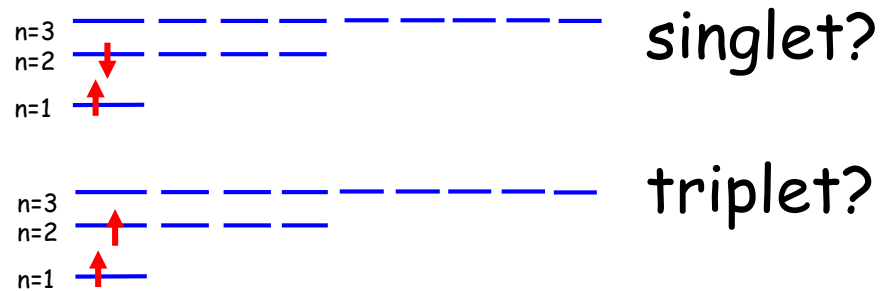
The **space** portion is **symmetric**, thus the **spin** portion must be **antisymmetric**.

$$\psi = \psi(\mathbf{r}_1, \mathbf{r}_2) \chi(\mathbf{s}_1, \mathbf{s}_2) = \frac{8}{\pi a^3} e^{-2(r_1+r_2)/a} \frac{1}{\sqrt{2}} (\uparrow_1 \downarrow_2 - \downarrow_1 \uparrow_2)$$

The cartoon version is:



Excited states?



Excited states? Two options ...

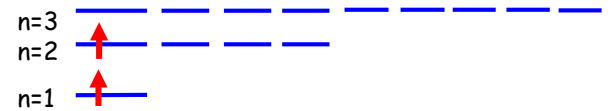
$$\Psi_{\text{singlet}} = \psi_S(\mathbf{r}_1, \mathbf{r}_2) \chi_{AS}(\mathbf{s}_1, \mathbf{s}_2) =$$

$$= \frac{1}{\sqrt{2}} [\psi_{100}(\mathbf{r}_1) \psi_{200}(\mathbf{r}_2) + \psi_{200}(\mathbf{r}_1) \psi_{100}(\mathbf{r}_2)] \quad \frac{1}{\sqrt{2}} (\uparrow_1 \downarrow_2 - \downarrow_1 \uparrow_2)$$



$$\Psi_{\text{triplet}} = \psi_{AS}(\mathbf{r}_1, \mathbf{r}_2) \chi_S(\mathbf{s}_1, \mathbf{s}_2) =$$

$$= \frac{1}{\sqrt{2}} [\psi_{100}(\mathbf{r}_1) \psi_{200}(\mathbf{r}_2) - \psi_{200}(\mathbf{r}_1) \psi_{100}(\mathbf{r}_2)] \quad \frac{1}{\sqrt{2}} (\uparrow_1 \downarrow_2 + \downarrow_1 \uparrow_2)$$

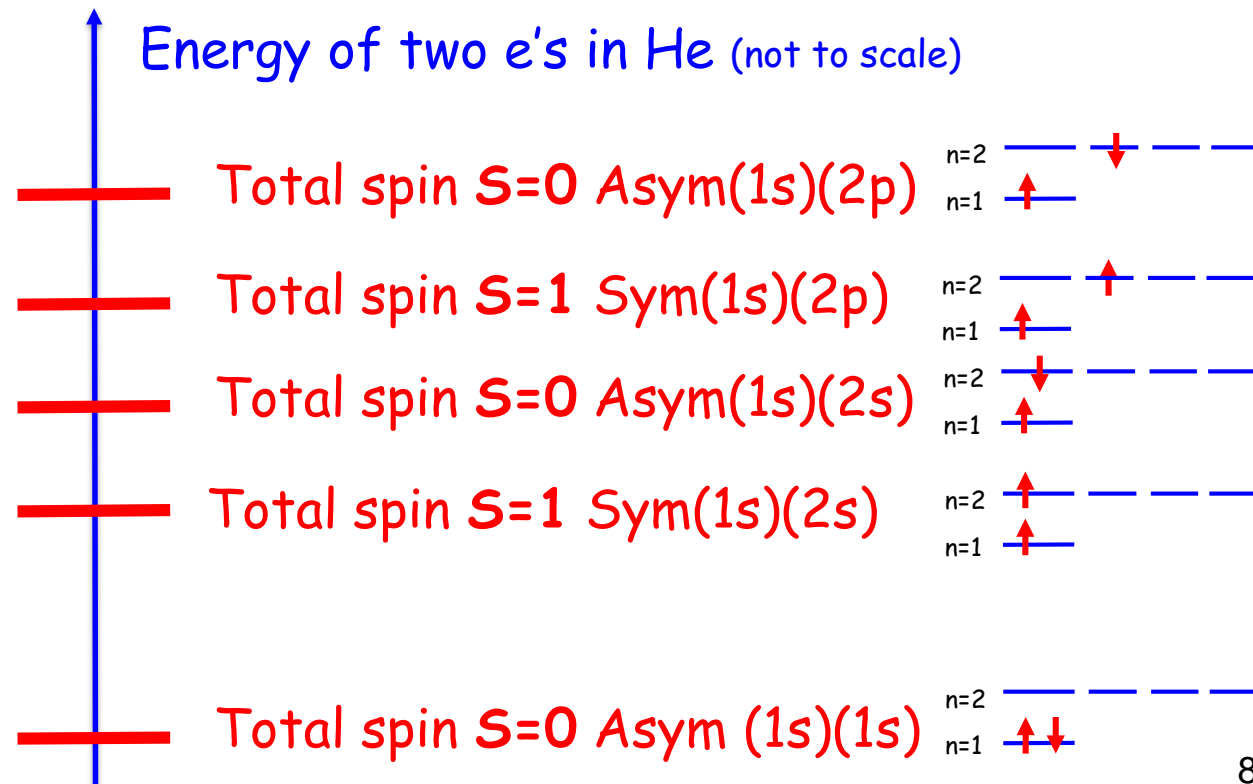


If e-e neglected, then singlet
and triplet are degenerate

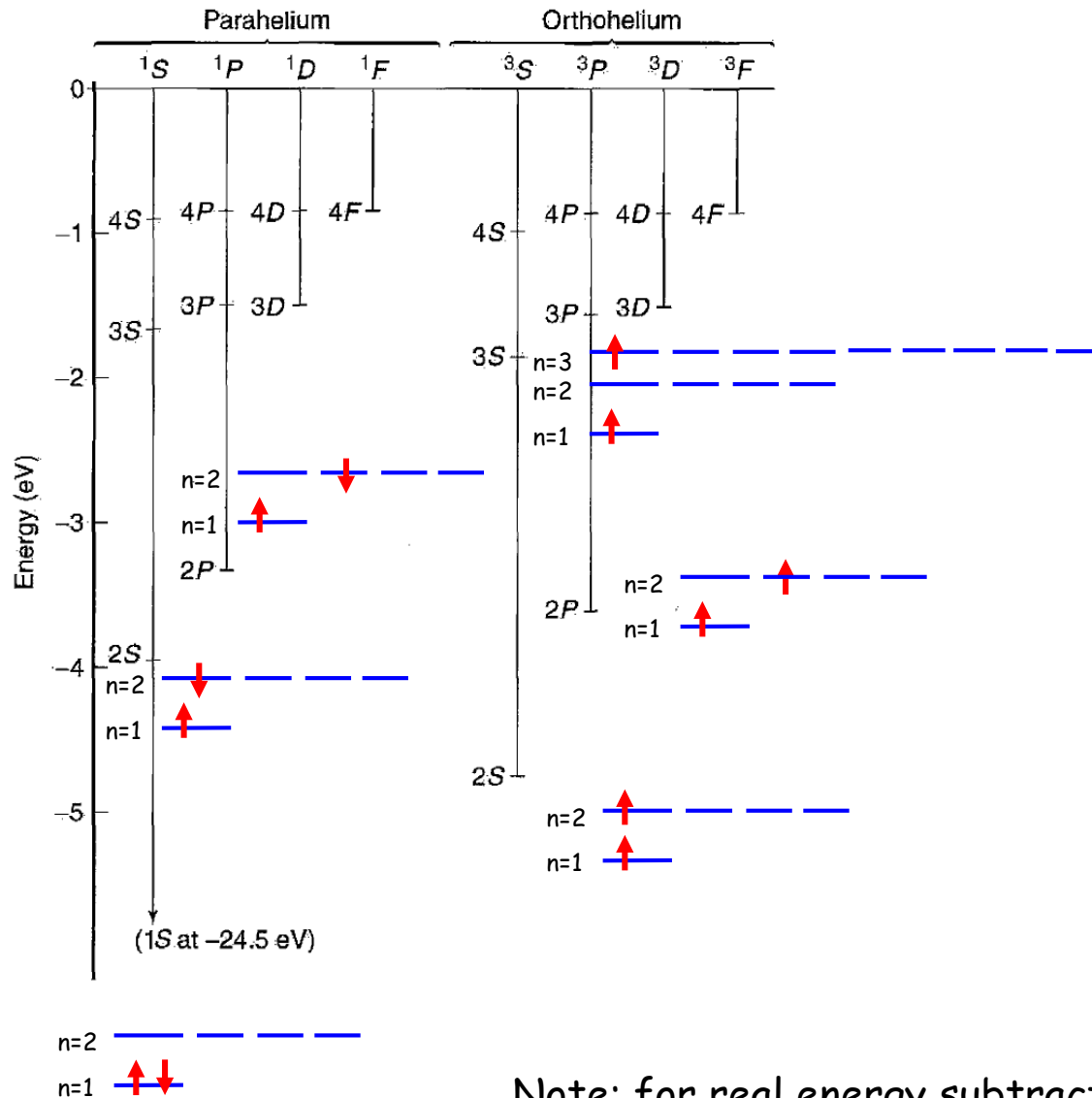
If $e-e$ is brought back, at least qualitatively, then the degeneracy is broken

Due to the effective "exchange forces", the AS space-like combination keeps the two electrons a bit further apart ... (for AS sector the exchange force is "repulsive"; for S sector is "attractive")

Then, the energy levels for two electrons is:



Real numbers from book

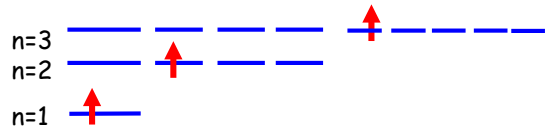


Note: for real energy subtract 54.4 eV
 (1 eV $\sim 1.6 \cdot 10^{-19}$ J)

Not in book

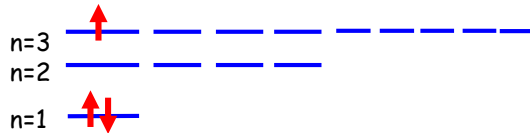
Z=3, Lithium

$S=3/2$
 $L=3$



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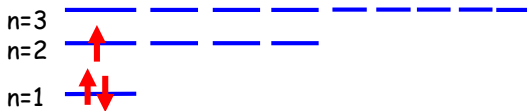
$S=1/2$
 $L=0$



$S=1/2$
 $L=1$

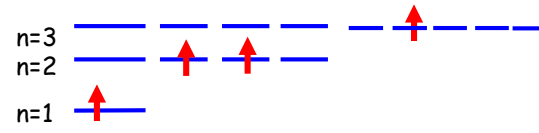


$S=1/2$
 $L=0$



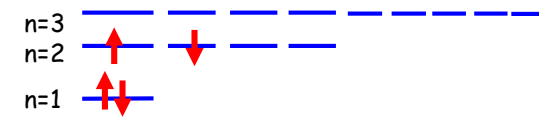
Z=4, Beryllium

$S=2$
 $L=4$

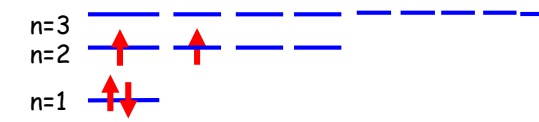


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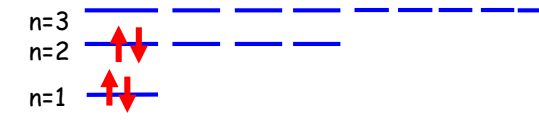
$S=0$
 $L=1$



$S=1$
 $L=1$



$S=0$
 $L=0$



The excited states become complicated fast!