

6.2 Degenerate Perturbation Theory:

What happens if the state you are correcting by perturbation theory is degenerate?

Consider first **degeneracy two**. We will drop the index "n" because degenerate states are not generic but special cases. This means:

$$H^0 \psi_a^0 = E^0 \psi_a^0 \quad H^0 \psi_b^0 = E^0 \psi_b^0 \quad \langle \psi_a^0 | \psi_b^0 \rangle = 0$$

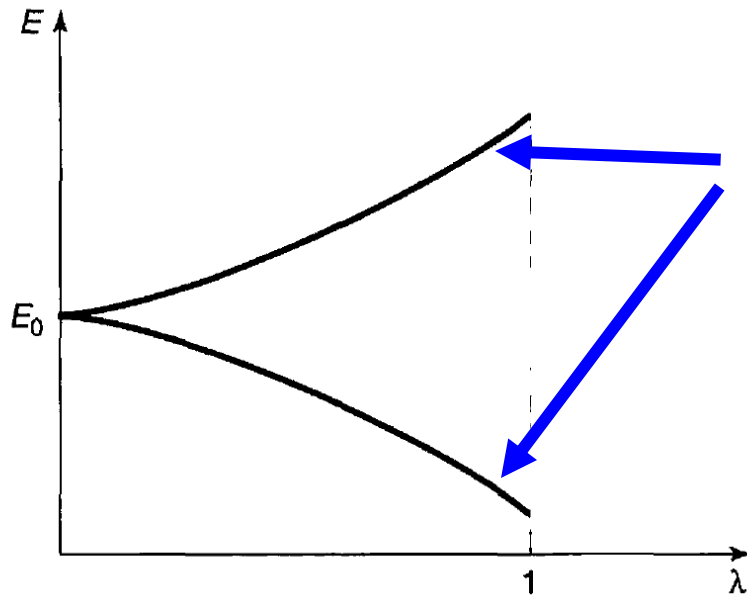
Same energy, like 2s and 2p levels in H atom.

Any linear combination has the same energy.

$$\psi^0 = \alpha \psi_a^0 + \beta \psi_b^0$$

$$H^0 \psi^0 = E^0 \psi^0$$

Usually, the perturbation H' will break the degeneracy.



Each split state is a particular linear combination of the two original degenerate states.

We will repeat the same steps as before:

$$H = H^0 + \lambda H'$$

$$E = E^0 + \lambda E^1 + \lambda^2 E^2 + \dots$$

$$\psi = \psi^0 + \lambda \psi^1 + \lambda^2 \psi^2 + \dots$$

To lowest order we arrive to the same equation as before, just without the index n.

$$H^0 \psi^1 + H' \psi^0 = E^0 \psi^1 + E^1 \psi^0$$

Again, follow the **same procedure** as before, namely take the **inner product** but using first ψ_a^0 and then ψ_b^0 :

$$\cancel{\langle \psi_a^0 | H^0 \psi^1 \rangle} + \langle \psi_a^0 | H' \psi^0 \rangle = E^0 \cancel{\langle \psi_a^0 | \psi^1 \rangle} + E^1 \underbrace{\langle \psi_a^0 | \psi^0 \rangle}$$

Via hermiticity of H^0 , again these two cancel.

Gives only α

$$\alpha \langle \psi_a^0 | H' | \psi_a^0 \rangle + \beta \langle \psi_a^0 | H' | \psi_b^0 \rangle = \alpha E^1$$

$$\alpha W_{aa} + \beta W_{ab} = \alpha E^1$$

When the "b" degenerate state is used, we find a similar formula:

$$\alpha W_{ba} + \beta W_{bb} = \beta E^1$$

$$\alpha W_{aa} + \beta W_{ab} = \alpha E^1$$

It is like a 2x2 matrix problem. α and β are the components of the eigenfunction and E^1 the eigenvalue.

$$E^1_{\pm} = \frac{1}{2} \left[W_{aa} + W_{bb} \pm \sqrt{(W_{aa} - W_{bb})^2 + 4|W_{ab}|^2} \right]$$