

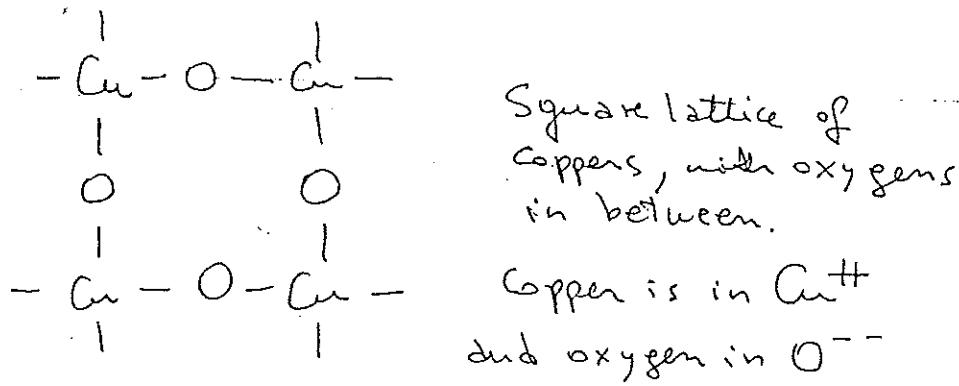
## High-T<sub>c</sub> Superconductors

Discovered in 1986.

Current record 133K (154K under pressure).

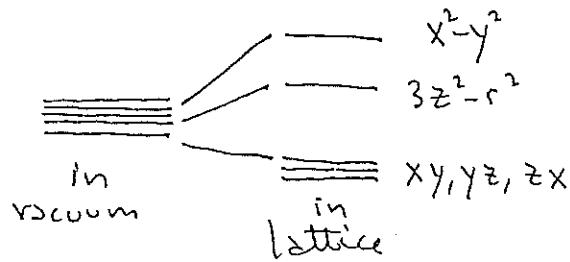
Widely believed that all the action is in the common factor i.e. in the CuO<sub>2</sub> planes.

The structure is two dimensional.

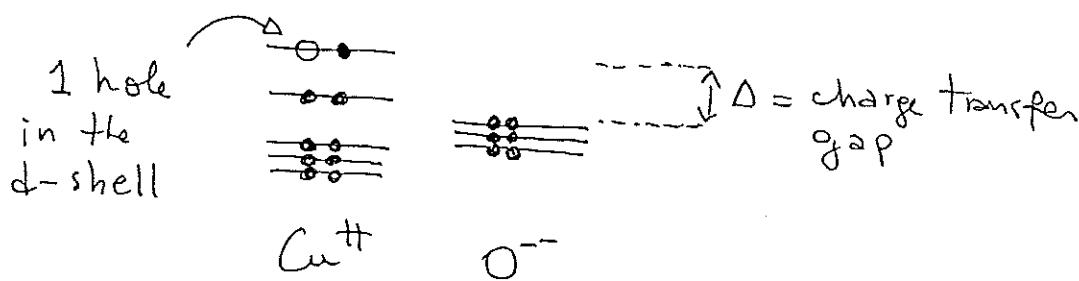


O<sup>-</sup> means that the p-orbitals are full.

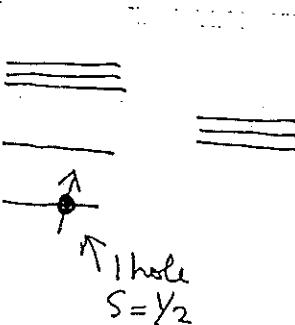
Cu<sup>+</sup> is a bit more complicated. The shell that matters is the "d". If we would have rotational invariance this orbital would have a degeneracy 5. However, in the standard "environment" of the transition metal Cu<sup>+</sup>, which is a cage of oxygens, rotational invariance is broken:



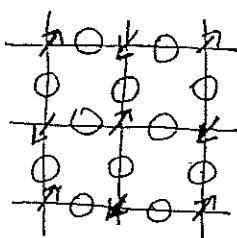
2. The number of electrons corresponds to the following situation:



this usually is inverted and the hole language is used.

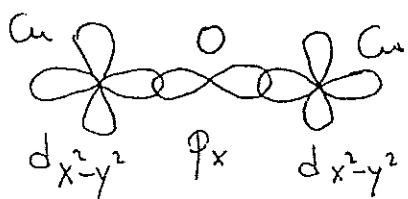


Then, the situation is:



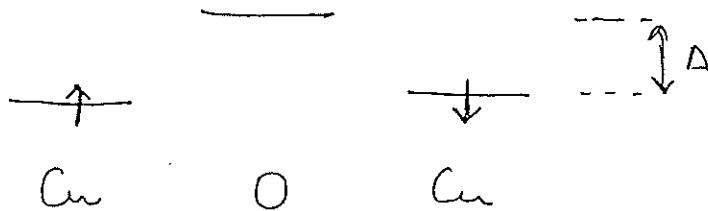
(we are assuming that there is one hole per Cu<sup>+</sup>)

Now, how do these spins interact?

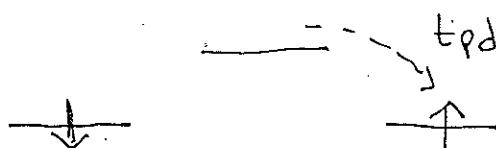
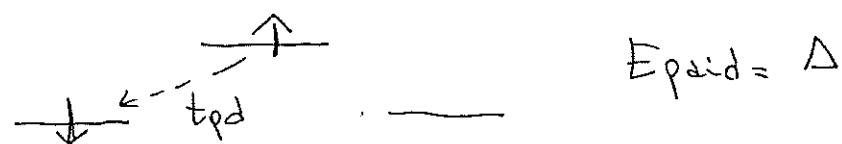
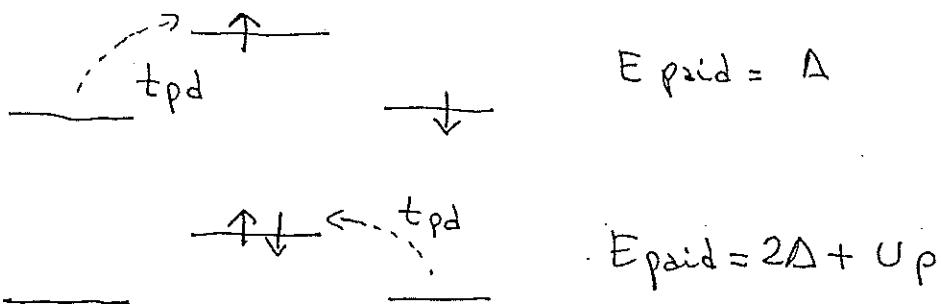


There is an overlap between orbitals that allows for the movement of holes from Cu to O and viceversa.

3.



These spins can be interchanged in four steps



$$\text{Effective coupling } J \propto \frac{t_{pd}^4}{\Delta^2(\Delta + U_p)} \quad (\text{superexchange})$$

$$J \sim 0.1 \text{ eV}, \quad t_{pd} \sim 1 \text{ eV}$$

$$\Delta \sim 2-3 \text{ eV}$$

$$U_p \sim 4 \text{ eV}$$

$$U_d \sim 10 \text{ eV}$$

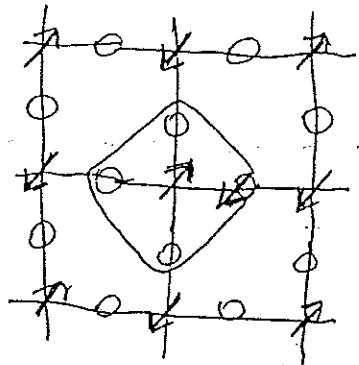
4.

$$H = J \sum \vec{S}_i \cdot \vec{S}_j \quad (\text{Heisenberg})$$

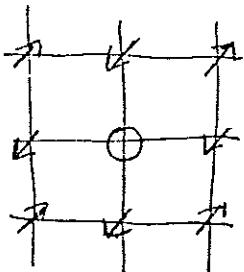
$\langle i,j \rangle$

Cu-Cu links

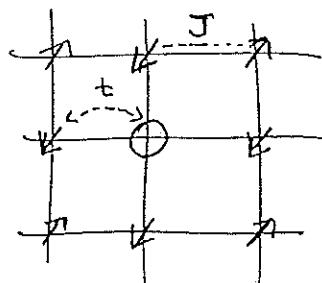
What happens with more holes?



The new hole goes to an oxygen (on Cu a large U is paid) and it tends to form a singlet state with a copper. This is called Zhang-Rice singlet. The effect is a state  $S=0$  centered at the copper.



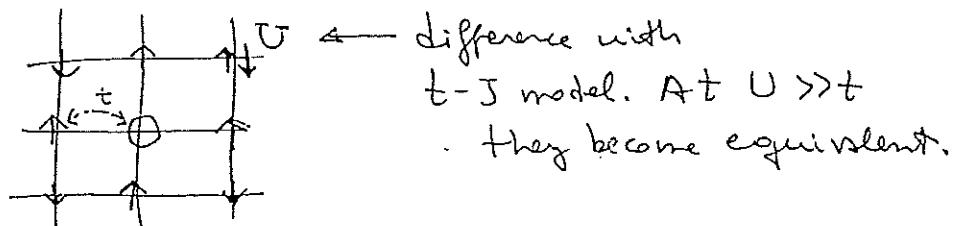
Oxygens are gone. Per site there are 3 possibilities: ↑, ↓ or ○. These holes can move into what we call the "t-J model"



An equivalent model is the one band Hubbard model:

$$H = -t \sum_{\langle ij \rangle} (c_{i\sigma}^{\dagger} c_{j\sigma} + h.c.) + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

$\sigma = \uparrow, \downarrow$



At very large  $U/t$  and half-filling the situation is that only one particle per site exists. Through a 2 steps process, the effective model at large  $U/t$  is the Heisenberg model:

$$\begin{array}{c} \uparrow \quad \downarrow \\ \nearrow t \quad \searrow t \\ - \quad \nabla \end{array} \quad E = U$$

$$\begin{array}{c} \leftarrow t \quad \uparrow \\ \uparrow \quad \uparrow \end{array}$$

$$J \propto \frac{t^2}{U}.$$

Then, it is worth studying the one band Hubbard model at  $\langle n \rangle = 1$  to see the appearance of AF.