

$$\beta_1 = \frac{(\mu'-1)}{b^3(1+\frac{\mu'}{2})} \cdot \frac{3m}{4\pi} \frac{(1+\frac{\mu'}{2})b^3}{[(1+2\mu')(1+\frac{\mu'}{2})b^3 - (\mu'-1)^2 a^3]}$$

$$\beta_1 = \frac{3m}{4\pi} \frac{(\mu'-1)}{[(1+2\mu')(1+\frac{\mu'}{2})b^3 - (\mu'-1)^2 a^3]}$$

Then we return to α_1 :

$$\alpha_1 = \beta_1 b^3 + \delta_1 = \frac{3m}{4\pi} \frac{(\mu'-1)b^3 + (1+\frac{\mu'}{2})b^3}{[(1+2\mu')(1+\frac{\mu'}{2})b^3 - (\mu'-1)^2 a^3]}$$

$$= \frac{3m}{4\pi} \frac{\frac{3}{2}b^3}{[(1+2\mu')(1+\frac{\mu'}{2})b^3 - (\mu'-1)^2 a^3]} = \alpha_1$$

Finally: $\delta_1 = \frac{1}{a^3} (\beta_1 a^3 + \delta_1 - \frac{m}{4\pi})$

$$= \frac{m}{4\pi a^3} \left[\frac{3[(\mu'-1)a^3 + (1+\frac{\mu'}{2})b^3]}{[(1+2\mu')(1+\frac{\mu'}{2})b^3 - (\mu'-1)^2 a^3]} - 1 \right] = \delta_1$$

(c) Consider now the limit of large μ
 i.e. we keep the dominant term
 in this limit.

$$\gamma_1 \approx \frac{3m}{4\pi\mu'} \frac{(b^3/2)}{(b^3 - a^3)}$$

$$\beta_1 \approx \frac{3m}{4\pi\mu'} \frac{1}{(b^3 - a^3)}$$

$$\alpha_1 \approx \frac{3m}{4\pi\mu'} \frac{(\frac{3}{2}b^3)}{(b^3 - a^3)}$$

$$\delta_1 \approx \frac{m}{4\pi a^3} \left[\frac{3}{\mu'} \frac{(a^3 + b^3/2)}{(b^3 - a^3)} - 1 \right]$$

$$\gamma_1, \beta_1, \alpha_1 \rightarrow 0 \quad \text{as } \mu' \rightarrow \infty$$

$$\delta_1 \rightarrow -\frac{m}{4\pi a^3} \quad \text{as } \mu' \rightarrow \infty$$

But we have to be careful! the field that
 matters is \vec{B} , but Φ_M generates \vec{H} .

To answer (c), let us focus on α_1 , the coefficient that matters outside the shell.

$$\Phi_M \stackrel{r > b}{=} \frac{\alpha_1}{r^2} P_1(\cos\theta) \approx \frac{3m}{4\pi(b^3 - a^3)\frac{\mu}{\mu_0}} \frac{3}{2} b^3 \frac{\cos\theta}{r^2} =$$

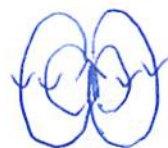
$$= \frac{\frac{9}{2} b^3}{(b^3 - a^3)\frac{\mu}{\mu_0}} \cdot \underbrace{\frac{m \cos\theta}{4\pi r^2}}_{\text{standard potential of a dipole}}$$

dimensionless factor in front that brings down the intensity of this dipole at large μ .

$$\vec{B} \stackrel{r > b}{=} -\mu_0 \nabla \Phi_M = -\mu_0 \left[\frac{\frac{9}{2} b^3}{(b^3 - a^3)\frac{\mu}{\mu_0}} \right] \nabla \left(\frac{m \cos\theta}{4\pi r^2} \right)$$

usual gradient that leads to the conical shape

Coefficient that reduces the magnitude of the dipole for $r > b$.



(d) Inside the shell we have to be careful, since Φ_M may be reduced as $\frac{1}{\mu}$ at large μ but to get \vec{B} we need to multiply by μ .

$$\vec{B} = -\mu \nabla \Phi_M$$

↑ important!

$$\begin{aligned} \Phi_M & \text{ at large } \mu = (\beta_1 r + \frac{\gamma_1}{r^2}) \cos \theta \stackrel{\text{large } \mu}{=} \frac{3m}{4\pi \mu'} \left[\frac{r}{(b^3 - a^3)} + \frac{b^3/2}{(b^3 - a^3)r^2} \right] \cos \theta \\ & = \frac{3m}{4\pi \mu' (b^3 - a^3)} \left(r + \frac{b^3}{2r^2} \right) \cos \theta \end{aligned}$$

$$\vec{B} = -\mu \nabla \Phi_M = \frac{-\mu 3m}{4\pi \mu' (b^3 - a^3)} \nabla \left[\left(r + \frac{b^3}{2r^2} \right) \cos \theta \right]$$

μ_0

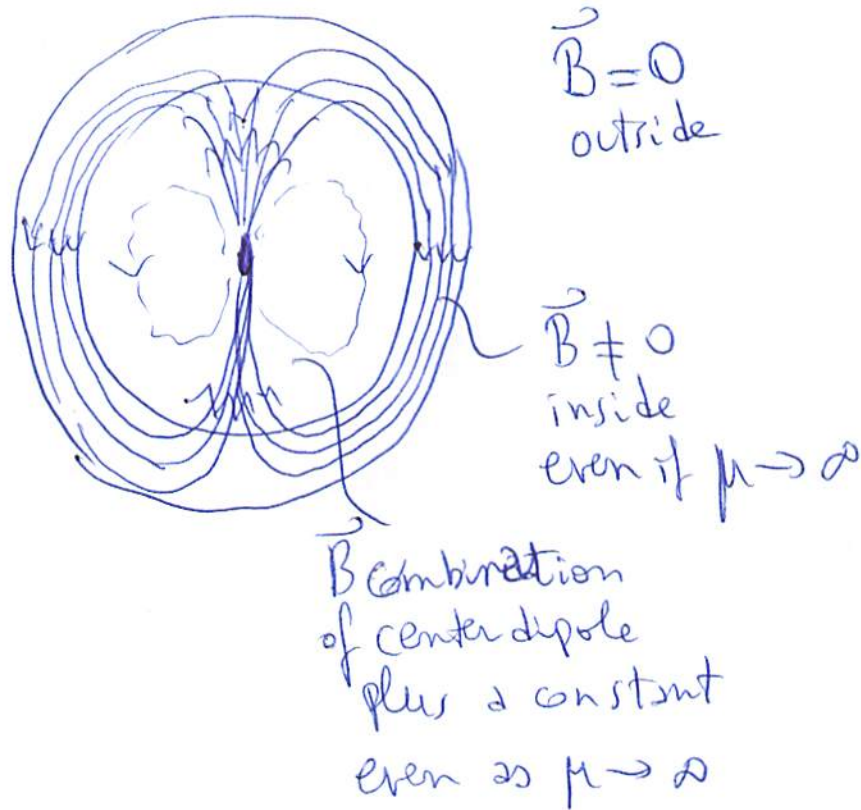
$$\vec{B} = \frac{-3m\mu_0}{4\pi (b^3 - a^3)} \nabla \left[\left(r + \frac{b^3}{2r^2} \right) \cos \theta \right]$$

\vec{B} is independent of μ inside the shell.

$$\vec{B} \neq 0 \text{ even if } \mu \rightarrow \infty$$

Combination of a constant \vec{B} (because $r \cos \theta = z$) and a dipole

For completeness, in the center hole the field \vec{B} is also non zero. Drawing lines basically you get at $\mu \rightarrow \infty$ the following lines of \vec{B} :



This is only a crude sketch to show that as $\mu \rightarrow \infty$, the outside has no lines, the "shell" has dense lines and the inside also has non zero lines of \vec{B} .