

3 images!

①

I need an arrangement of charges such that the potential cancels at the surfaces shaded.

If I use only one, say ^{image $-q$} at $(-a, a, 0)$ I can cancel the potential at the zy plane but not at the other plane. If I use both $(-a, a, 0)$ and $(a, -a, 0)$, the contribution of the three does not cancel at each plane.

Then, I need 3 images! Together with the real charge I have a total of 4 and by pairs they produce the desired result. The total potential at an arbitrary point (x, y, z) in the first quadrant is:

$$\Phi = \frac{q}{4\pi\epsilon_0} \left[\frac{1}{\sqrt{(x-a)^2 + (y-a)^2 + z^2}} + \frac{1}{\sqrt{(x+a)^2 + (y+a)^2 + z^2}} - \frac{1}{\sqrt{(x+a)^2 + (y-a)^2 + z^2}} - \frac{1}{\sqrt{(x-a)^2 + (y+a)^2 + z^2}} \right]$$

It is easy to confirm that this potential cancels at $x=0$ and at $y=0$.

(b) The force at the real charge is made of 3 terms, due to the images.

Diagram showing a charge q at $(a, a, 0)$ in a 2D coordinate system with x and y axes. The charge is positive (+). The x-axis is horizontal and the y-axis is vertical. The charge is located in the first quadrant. There are also signs for charges at $(0,0)$: '+' at the origin, '-' at $(a,0)$, and '-' at $(0,a)$.

$$\vec{F}_{\text{at } (a,a,0)} = \frac{q^2}{4\pi\epsilon_0} \left[-\frac{\hat{e}_x}{(2a)^2} - \frac{\hat{e}_y}{(2a)^2} + \frac{\hat{e}_x + \hat{e}_y}{\sqrt{2}(\sqrt{2}2a)^2} \right]$$

Annotations for the force equation:

- to normalize unit vector to 1 (pointing to the $\sqrt{2}$ in the denominator)
- distance from $(a, a, 0)$ to $(-a, a, 0)$ (pointing to the $2a$ in the denominator)

$$= \frac{q^2}{4\pi\epsilon_0} \left[-\frac{\hat{e}_x}{4a^2} - \frac{\hat{e}_y}{4a^2} + \frac{(\hat{e}_x + \hat{e}_y)}{8a^2\sqrt{2}} \right].$$

(c) To find the Dirichlet Green function $G_D(\vec{x}, \vec{x}')$ we use the discussion in the book and lectures involving the analogy of G_D with the potential Φ found in (a). Basically we have to locate the charge q at an arbitrary point (x', y', z') instead of $(a, a, 0)$, find the images at $(-x', y', z')$, $(x', -y', z')$ and $(-x', -y', z')$ and then get the potential of $q = 4\pi\epsilon_0$.

$$G_D(\vec{x}, \vec{x}') = \left[\frac{1}{\sqrt{(x-x')^2 + (y-y')^2 + (z-z')^2}} + \frac{1}{\sqrt{(x+x')^2 + (y+y')^2 + (z-z')^2}} \right] - \left[\frac{1}{\sqrt{(x+x')^2 + (y-y')^2 + (z-z')^2}} - \frac{1}{\sqrt{(x-x')^2 + (y+y')^2 + (z-z')^2}} \right]$$

At $y'=0$ it cancels. At $x'=0$ it cancels.

(d) Formally the surface integral that I need is

$$\text{" } -\frac{1}{4\pi} \oint_S \Phi(x') \frac{\partial G_D}{\partial n'} da' \text{" from the book.}$$

In our case the surface is the ^{top} half plane at $x'=0$

So \oint_S means $\int_0^{\infty} \int_{-\infty}^{+\infty} dz'$. " $\Phi(x')$ " is just a constant V .

" da' " is just $dy'dz'$, already written in the integral.

" $\frac{\partial G_D}{\partial n'}$ " means $\nabla G_D \cdot \vec{n}'$ and \vec{n}' is $(-\hat{e}_x)$. So $\nabla G_D \cdot \vec{n}'$ means $-\frac{d}{dx'} G_D(\vec{x}, \vec{x}')$ at $x'=0$.
↑
from the volume considered outwards.

$$\begin{aligned} \frac{d}{dx'} G_D(\vec{x}, \vec{x}') &= \frac{d}{dx'} \left[(x-x')^2 + (y-y')^2 + (z-z')^2 \right]^{-1/2} \\ &+ \frac{d}{dx'} \left[(x+x')^2 + (y+y')^2 + (z-z')^2 \right]^{-1/2} \\ &+ \frac{d}{dx'} \left[- \left[(x+x')^2 + (y-y')^2 + (z-z')^2 \right]^{-1/2} \right] \\ &+ \frac{d}{dx'} \left[- \left[(x-x')^2 + (y+y')^2 + (z-z')^2 \right]^{-1/2} \right] = \end{aligned}$$

$$\begin{aligned} &= -\frac{1}{2} \left[(x-x')^2 + (y-y')^2 + (z-z')^2 \right]^{-3/2} 2(x-x')(-1) \\ &- \frac{1}{2} \left[(x+x')^2 + (y+y')^2 + (z-z')^2 \right]^{-3/2} 2(x+x') \\ &+ \frac{1}{2} \left[(x+x')^2 + (y-y')^2 + (z-z')^2 \right]^{-3/2} 2(x+x') \\ &+ \frac{1}{2} \left[(x-x')^2 + (y+y')^2 + (z-z')^2 \right]^{-3/2} 2(x-x')(-1) \end{aligned}$$

Now I have to take $x'=0$ to be at the surface.

$$\begin{aligned} \left. \frac{d}{dx'} G_D(\vec{x}, \vec{x}') \right|_{\substack{\text{surface} \\ x'=0}} &= x \left[x^2 + (y-y')^2 + (z-z')^2 \right]^{-3/2} - x \left[x^2 + (y+y')^2 + (z-z')^2 \right]^{-3/2} \\ &+ x \left[x^2 + (y-y')^2 + (z-z')^2 \right]^{-3/2} - x \left[x^2 + (y+y')^2 + (z-z')^2 \right]^{-3/2} \\ &= \frac{2x}{\left[x^2 + (y-y')^2 + (z-z')^2 \right]^{3/2}} - \frac{2x}{\left[x^2 + (y+y')^2 + (z-z')^2 \right]^{3/2}} \end{aligned}$$

The final expression then is:

$$\Phi(\vec{x}) = + \frac{V}{4\pi} \int_0^{\infty} dy' \int_{-\infty}^{+\infty} dz' \quad 2x \left[\frac{1}{[x^2 + (y-y')^2 + (z-z')^2]^{3/2}} - \frac{1}{[x^2 + (y+y')^2 + (z-z')^2]^{3/2}} \right]$$

remember
 $\vec{M}' = -\hat{e}_x$

$\vec{x} = (x, y, z)$ is a point in the first quadrant
 (y', z') are coordinates of a point at the surface
at constant potential V .

(e)

We use
$$\int \frac{du}{(a^2+u^2)^{3/2}} = \frac{u}{a^2(a^2+u^2)^{1/2}}$$

(i)
$$\int_{-\infty}^{+\infty} dz' \frac{1}{\sqrt{x^2+(y-y')^2+(z'-z)^2}} \stackrel{z'-z=u}{=} \int_{-\infty}^{+\infty} \frac{du}{\sqrt{x^2+(y-y')^2+u^2}} = \int_{-\infty}^{+\infty} \frac{du}{\sqrt{a^2+u^2}} =$$

$$= \frac{u}{a^2(a^2+u^2)^{1/2}} \Big|_{-\infty}^{+\infty} = \lim_{u \rightarrow \infty} \frac{u}{a^2(a^2+u^2)^{1/2}} - \lim_{u \rightarrow -\infty} \frac{u}{a^2(a^2+u^2)^{1/2}} =$$

$$= \frac{1}{a^2} + \frac{1}{a^2} = \frac{2}{a^2} = \frac{2}{x^2+(y-y')^2}$$

(ii) Same but with " a^2 " = $x^2+(y+y')^2$

Then
$$\int_{-\infty}^{+\infty} dz' \left[\frac{1}{(\quad)^{3/2}} - \frac{1}{(\quad)^{3/2}} \right] = \frac{2}{x^2+(y-y')^2} \leftrightarrow \frac{2}{x^2+(y+y')^2}$$

What remains is:

$$\frac{+\sqrt{2}x}{4\pi} \int_0^{\infty} dy' 2 \left[\frac{1}{x^2+(y-y')^2} \leftrightarrow \frac{1}{x^2+(y+y')^2} \right] =$$

$$= + \frac{Vx}{\pi} \left[\int_0^{\infty} dy' \frac{1}{x^2 + \underbrace{(y-y')^2}_{\text{or } (y'-y)^2}} - \int_0^{\infty} dy' \frac{1}{x^2 + (y+y')^2} \right]$$

$$= + \frac{Vx}{\pi} \left[\int_{-y}^{\infty} du \frac{1}{x^2 + u^2} - \int_{+y}^{\infty} du \frac{1}{x^2 + u^2} \right]$$

\nearrow $y' - y = u$ \nearrow $y' + y = u$

$$= + \frac{Vx}{\pi} \left[\frac{1}{x} \tan^{-1}\left(\frac{u}{x}\right) \Big|_{-y}^{\infty} - \frac{1}{x} \tan^{-1}\left(\frac{u}{x}\right) \Big|_{+y}^{\infty} \right]$$

$$= + \frac{V}{\pi} \left[\cancel{\tan^{-1}(\infty)} - \tan^{-1}\left(-\frac{y}{x}\right) - \left(\cancel{\tan^{-1}(\infty)} - \tan^{-1}\left(\frac{y}{x}\right) \right) \right]$$

$$= \frac{V}{\pi} \left[\tan^{-1}\left(\frac{y}{x}\right) - \underbrace{\tan^{-1}\left(-\frac{y}{x}\right)}_{\text{odd function}} \right] = \frac{2V}{\pi} \tan^{-1}\left(\frac{y}{x}\right)$$

$$\Phi(\vec{x}) = + \frac{2V}{\pi} \tan^{-1}\left(\frac{y}{x}\right)$$

indep. of z
which is correct
by symmetry

If $y=0$, $\tan^{-1}(0) = 0$
which is the correct boundary
condition!

$$\text{If } x=0, \text{ we get } \Phi(\vec{x}) = + \frac{2V}{\pi} \cdot \frac{\pi}{2}$$

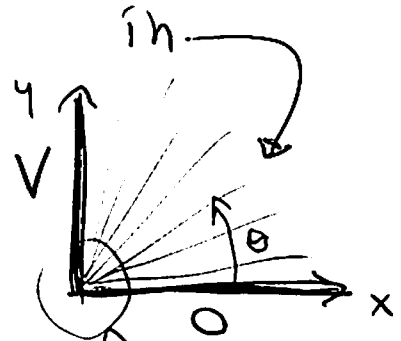
$= +V$ which is

~~which is correct~~ correct!

$$\boxed{\Phi(\vec{x}) = \frac{2V}{\pi} \tan^{-1}\left(\frac{y}{x}\right)}$$

(indep. of z)

for any point



$$x = r \cos \theta$$

$$y = r \sin \theta$$

$$\tan \theta = \frac{y}{x}$$

$$\theta = \tan^{-1}\left(\frac{y}{x}\right)$$

$$\boxed{\Phi(\vec{x}) = \frac{2V}{\pi} \theta}$$

Huge fields here
because $\nabla \phi$ (edge)
is huge.