

9/4

Homework:

Weisong Tu - Office Hours
Wed 2 to 3 PM
@ Tutoring Center
Room 201 Nielsen.

Vectors and Covectors.

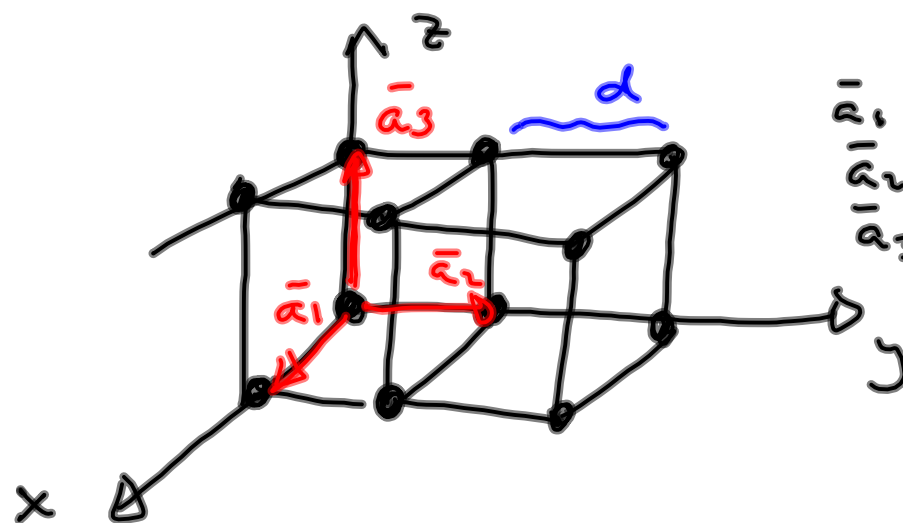
Solid State Physics:

A crystal is formed by an orderly array of ions that form a lattice in real space. The lattice is called Bravais lattice and there is a small number of this lattices such as cubic, face centered cubic, body center cubic, etc.

To define the position of a ion in the lattice we define a basis:

Ex: cubic lattice

d : lattice spacing



$$\begin{aligned}\bar{a}_1 &= d \hat{x} \\ \bar{a}_2 &= d \hat{y} \\ \bar{a}_3 &= d \hat{z}\end{aligned}$$

\bar{a}_i : are vectors of length d that connect the ion chosen as the origine with a nearest neighbor

$\{\bar{a}_i\}$ are linearly independent.

Now the position of any ion at site \bar{R} can be written as:

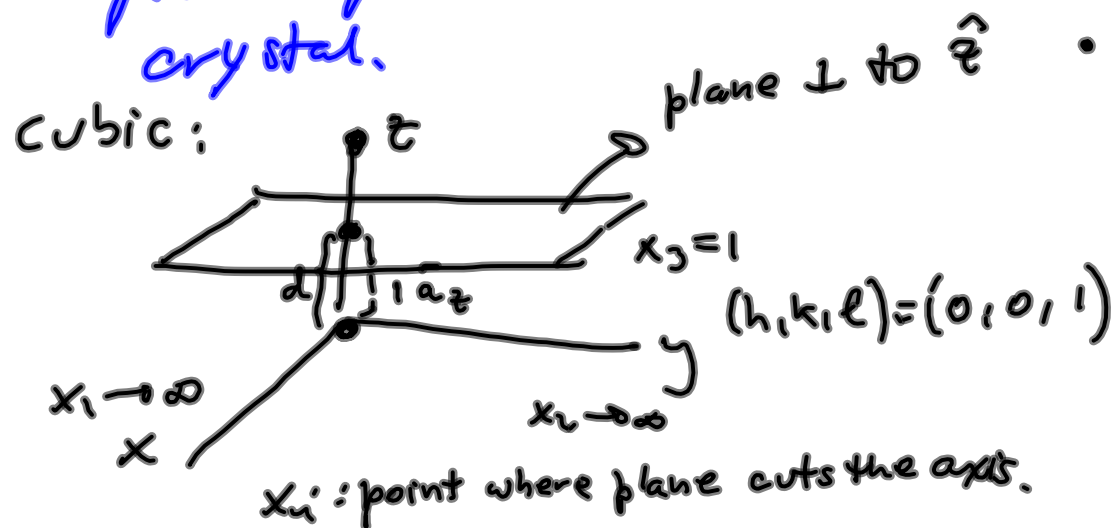
$$\bar{R} = m_1 \bar{a}_1 + m_2 \bar{a}_2 + m_3 \bar{a}_3 \quad \text{with } m_i \text{ integers}$$

We can use the vectors \bar{R} to define our lattice but we also can define the lattice using families of planes perpendicular to the vectors \bar{R} .

For example: in the cubic lattice we can use the families of planes perpendicular to \hat{x} , \hat{y} and \hat{z} .

Miller indices: are used by crystallographers to identify the families of planes in a system.

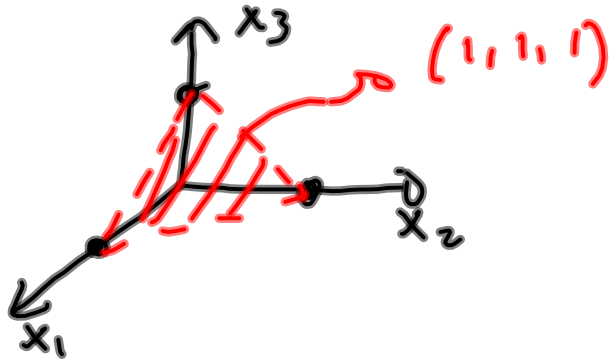
Planes are defined by 3 points in a Bravais lattice and the collection of all planes parallel to each other expands the crystal.



- The family is composed by planes parallel to $x-y$ and separated a distance d from each other.
- Find the plane in the family closer to the origin.
- define $(h, k, l) = \left(\frac{1}{x_1}, \frac{1}{x_2}, \frac{1}{x_3}\right)$

$(1, 0, 0)$ defines the family of planes \perp to \bar{a}_1
 $(0, 1, 0)$ " " " " " " " \bar{a}_2
 $(0, 0, 1)$ " " " " " " " \bar{a}_3

Notice that in some lattices we also can define families of planes that are NOT perpendicular to any basis vector.

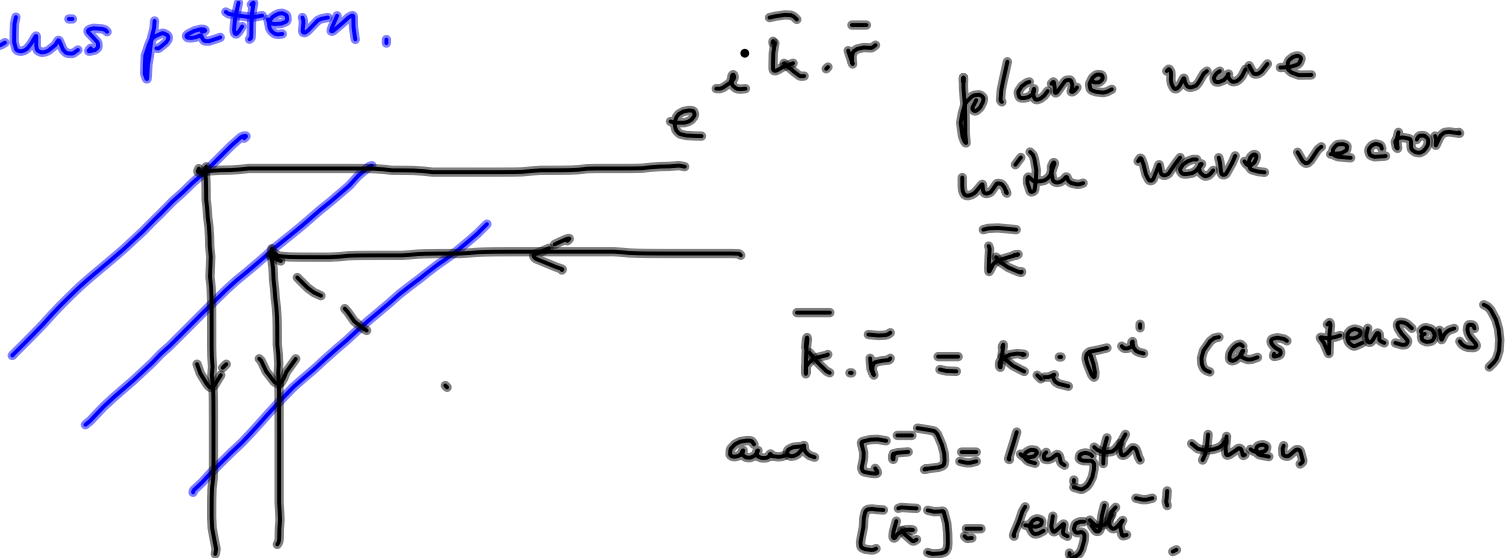


orthogonal to

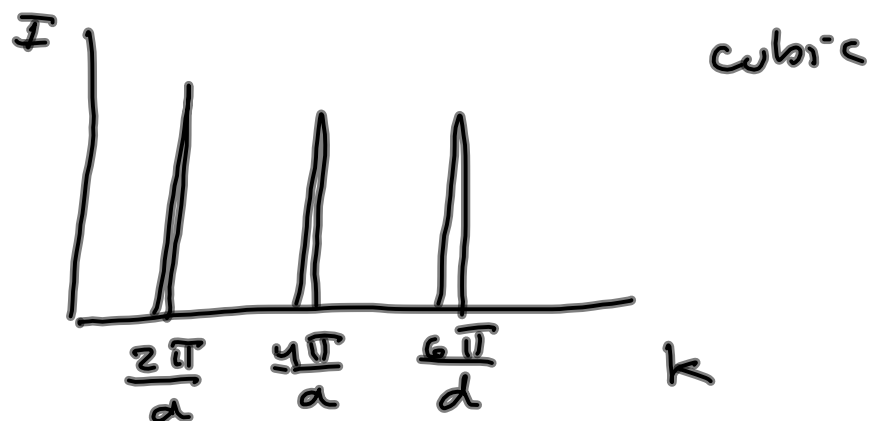
$$\bar{R} = \bar{a}_1 + \bar{a}_2 + \bar{a}_3$$

Why are planes important?

Lattice structure is studied by looking at the interference pattern obtained when x-rays are reflected on the lattice. The properties of the lattice planes determined this pattern.



Experimentalists can vary \bar{R} by controlling the energy (wavelength) of the electromagnetic wave -
 They find that the reflected intensity is a function of \bar{k} .



The reflected light has the form:

$$e^{i \bar{k} (\bar{r} + \bar{R})}$$

where \bar{R} is the position of the scatterer (ion) and \bar{r} is a general position in space.

The maximum intensity (the peaks) occur when \bar{k} is such that

$$e^{i\bar{k} \cdot \bar{R}} = 1 \quad \text{then } \bar{k} \cdot \bar{R} = 2\pi n \quad \downarrow \text{integer}$$

The values of \bar{k} such that $\bar{k} \cdot \bar{R} = 2\pi n$ will be called \bar{K} so that $\bar{K} \cdot \bar{R} = 2\pi n$.

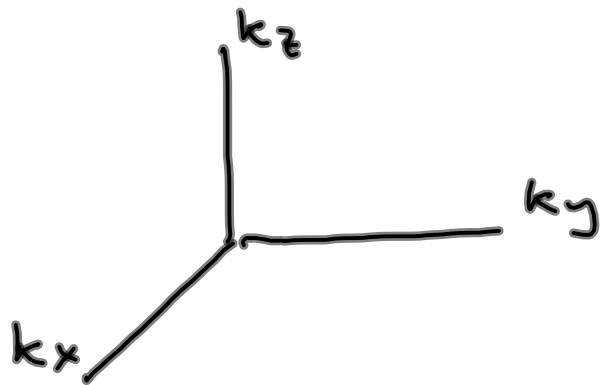
It can be shown that the vectors \bar{K} also form a periodic lattice. This lattice is in k -space or reciprocal space. The reciprocal space is called dual space by mathematicians.

Reciprocal basis:

In dual space (or reciprocal space) we can define a basis $\{\bar{b}_i\}$ that will generate all the vectors \bar{K} that form the reciprocal lattice (in tensor notation we will see that

$$\bar{b}_i = \hat{b}^i)$$

Example: Find the dual basis for the cubic lattice.



$$\bar{K} = m_1 \bar{b}_1 + m_2 \bar{b}_2 + m_3 \bar{b}_3$$

↓
 m_i are integers
 point in reciprocal lattice.

$$\int \bar{K} \cdot \bar{R} = K_i R^i = 2\pi n \quad \bar{K} \text{ belongs to reciprocal lattice.}$$

Then

$$\begin{aligned} \bar{K} \cdot \bar{R} = & m_1 n_1 \bar{a}_1 \cdot \bar{b}_1 + m_2 n_2 \bar{a}_2 \cdot \bar{b}_2 + m_3 n_3 \bar{a}_3 \cdot \bar{b}_3 + \\ & + n_1 m_2 \bar{a}_1 \cdot \bar{b}_2 + n_1 m_3 \bar{a}_1 \cdot \bar{b}_3 + n_2 m_3 \bar{a}_2 \cdot \bar{b}_3 + \dots = 2\pi n \end{aligned}$$

This is true if:

$$\bar{a}_i \cdot \bar{b}_i = 2\pi \quad \text{and} \quad \bar{a}_i \cdot \bar{b}_j = 0 \quad \text{if } i \neq j.$$

In general to find the reciprocal basis
we require:

$$a_i \cdot b^j = \delta_i^j 2\pi \quad \text{for physicists}$$

Mathematicians use:

$$a_i \cdot b^j = \delta_i^j$$

This works in any dimension.

For 3D you can use:

$$\bar{b}_i = 2\pi \frac{\bar{a}_j \times \bar{a}_k}{\underbrace{\bar{a}_i \cdot (\bar{a}_j \times \bar{a}_k)}_{\text{volume of unit cell in real space}}} \quad i, j, k \text{ in cyclic order.}$$

Clearly

$$\bar{a}_i \cdot \bar{b}_i = 2\pi$$

$$\bar{a}_j \cdot \bar{b}_i = 0 \quad \text{if } i \neq j.$$

Ex: cubic lattice:

$$\bar{b}_1 = 2\pi \frac{\overbrace{\bar{a}_2 \times \bar{a}_3}^{d^2 \hat{x}}}{\underbrace{\bar{a}_1 \cdot (\bar{a}_2 \times \bar{a}_3)}_{d^3}} = \frac{2\pi d^2 \hat{x}}{d^3} = \frac{2\pi}{d} \hat{x}$$

$$\bar{b}_i = \frac{2\pi}{a} \bar{x}_i \quad \text{cubic lattice of size } \frac{2\pi}{d} \quad \left(\text{volume } \frac{(2\pi)^3}{d^3} \right).$$