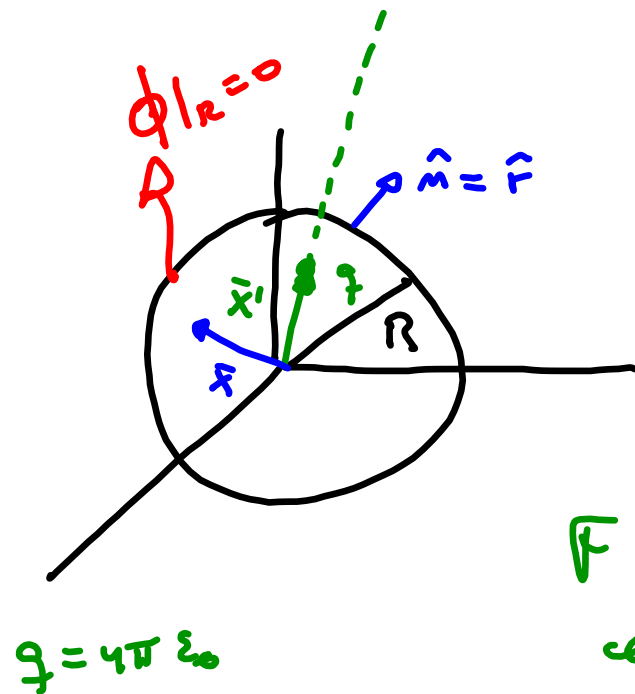


# Green Functions

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Spherical symmetry:



We want  $G(\bar{x}, \bar{x}')$   
for  $r \leq R$  ( $V$  is volume  
inside the sphere).

Dirichlet b.c.

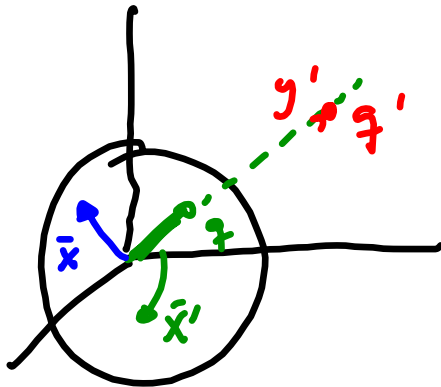
$$G(\bar{x}, \bar{x}') = \frac{1}{|\bar{x} - \bar{x}'|} + F(\bar{x}, \bar{x}')$$

$F(\bar{x}, \bar{x}')$  is the potential of  
external charges that satisfy

$$G|_{x=R} = 0.$$

$$G(\bar{x}, \bar{x}') = \frac{1}{|\bar{x} - \bar{x}'|} + \frac{q'}{4\pi\epsilon_0 |\bar{x} - \bar{y}'|} =$$

we need to find  $q'$  and  $\bar{y}'$   
so that  $G|_R = 0$ .



$$\hat{m} = \frac{|\bar{x}|}{R}$$

$$\hat{m}' = \frac{|\bar{x}'|}{R}$$

$$= \frac{1}{|x\hat{m} - x'\hat{m}'|} + \frac{q'}{4\pi\epsilon_0 |x\hat{m} - y'\hat{m}'|}$$

$$G(\bar{x}, \bar{x}')|_R = 0 = \frac{1}{|R\hat{m} - x'\hat{m}'|} + \frac{q'}{4\pi\epsilon_0 |R\hat{m} - y'\hat{m}'|}$$

$$0 = \frac{1}{R|\hat{m} - \frac{x'}{R}\hat{m}'|} + \frac{q'}{4\pi\epsilon_0 y'|\frac{R}{y'}\hat{m} - \hat{m}'|}$$

We request that

$$\frac{1}{R} + \frac{q'}{4\pi\epsilon_0 y'} = 0 \quad \therefore \boxed{\frac{1}{R} = -\frac{q'}{4\pi\epsilon_0 y'}} \quad (1)$$

and

$$\left| \hat{m} - \frac{x'}{R} \hat{m}' \right| = \left| \frac{R}{y'} \hat{m} - \hat{m}' \right|$$

$$\left( \hat{m} - \frac{x'}{R} \hat{m}' \right) \cdot \left( \hat{m} - \frac{x'}{R} \hat{m}' \right) = \left( \frac{R}{y'} \hat{m} - \hat{m}' \right) \cdot \left( \frac{R}{y'} \hat{m} - \hat{m}' \right)$$

$$\cancel{\hat{m} \cdot \hat{m}} - 2 \hat{m} \cdot \hat{m}' \frac{x'}{R} + \frac{x'^2}{R^2} \hat{m}' \cdot \hat{m}' = \frac{R^2}{y'^2} \hat{m} \cdot \hat{m} - 2 \hat{m}' \cdot \hat{m} \frac{R}{y'} + \cancel{\hat{m}' \cdot \hat{m}'}$$

$$\boxed{\frac{x'}{R} = \frac{R}{y'}} = \boxed{y' = \frac{R^2}{x'}} \quad (2)$$

Plugging ② in ① we find  $q'$ :

$$q' = -\frac{4\pi \epsilon_0 R}{x'} \quad \text{③}$$

Then replacing  $y'$  and  $q'$  in  $G$ :

$$G(\bar{x}, \bar{x}') = \frac{1}{|\bar{x} - \bar{x}'|} - \frac{R}{x' \left| \bar{x} - \frac{R^2}{x'} \hat{m}' \right|}$$

- $\bar{x} = (r, \theta, \varphi)$  best option for coordinates
- $\frac{\partial G}{\partial m'} \Big|_R \equiv \frac{\partial G}{\partial r'} \Big|_R$
- Expand in terms of  $Y_e^m(\theta, \varphi)$ .

We found that

$$\frac{1}{|\bar{x} - \bar{x}'|} = 4\pi \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} \frac{r_c^{\ell}}{r_>^{\ell+1}} \frac{1}{2\ell+1} Y_{\ell m}^*(\theta', \varphi') Y_{\ell m}(\theta, \varphi)$$

$r_c$  ( $r_>$ ) smaller (larger) between  $x$  and  $x'$ ,

$$\frac{R}{x' |\bar{x} - \frac{R^2}{x'} \hat{m}'|} = \frac{R}{x'} 4\pi \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} \frac{r_c^{\ell}}{r_>^{\ell+1}} \frac{1}{2\ell+1} Y_{\ell m}^*(\theta, \varphi) Y_{\ell m}(\theta', \varphi')$$

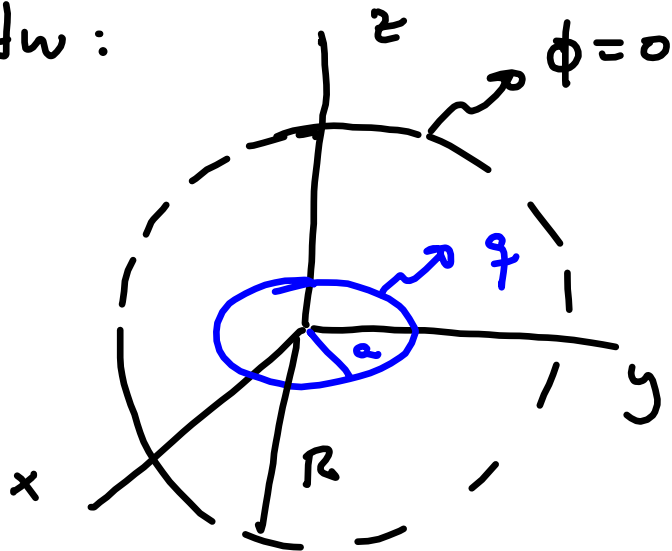
$r_c$  ( $r_>$ ) smaller (larger) between ( $x$  and  $\frac{R^2}{x'}$ ).

Try at home:

$$G(\bar{r}, \bar{r}') = 4\pi \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} \frac{1}{2\ell+1} \left( \frac{r_c^{\ell}}{r_>^{\ell+1}} - \frac{r^{\ell} r'^{\ell}}{R^{2\ell+1}} \right) Y_{\ell m}(\theta, \varphi) Y_{\ell m}^*(\theta', \varphi')$$

Now with  $G(\vec{r}, \vec{r}')$  you can solve several problems:

HW:



$$\rho(\vec{x}) = \frac{q}{2\pi a^2} \delta(r-a) \delta(\cos\theta)$$

Plug  $G$  and  $\rho$   
in expression for  
 $\phi(\vec{r})$  obtained  
last time.

$$\phi(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int_V \rho(\vec{r}') G(\vec{r}, \vec{r}') d^3r'$$

$$\phi(\vec{r}) = \frac{1}{4\pi\epsilon_0} 4\pi \frac{q}{2\pi a^2} \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{1}{2l+1} Y_{lm}(\theta, \varphi)$$

$$\int_0^R r'^2 dr' \delta(r'-a) \left( \frac{r_{<}^l}{r_{>}^{l+1}} - \frac{r^l r'^l}{R^{2l+1}} \right)$$

$\underbrace{\delta_{m,0} 2\pi \sqrt{\frac{2l+1}{4\pi}} P_l(\cos\theta)}_{2\bar{u}}$

$$\int_{-1}^1 d(\cos\theta) \delta(\cos\theta) \int_0^{2\pi} d\varphi' Y_{lm}^*(\theta', \varphi')$$

$\underbrace{\hspace{10em}}_{2\bar{u} \delta_{m,0} \text{ since } Y_{lm} \propto e^{-im\varphi'}}$

$$Y_{lm}^*(\theta, \varphi) = \sqrt{\frac{2l+1}{4\pi} \frac{(l-m)!}{(l+m)!}} P_l^m(\cos\theta) e^{-im\varphi}$$

Then the integral over  $\theta'$  and  $\varphi'$  gives us:

$$P_\ell(0) \sqrt{\frac{2\ell+1}{4\pi}} 2\pi \delta_{m,0} \underbrace{\sqrt{\frac{2\ell+1}{4\pi}} P_\ell(\cos\theta)}$$

Then:

$$\phi(r) = \frac{2\pi}{4\pi \epsilon_0} \frac{4\pi}{2\pi a^2} \sum_{\ell=0}^{\infty} \frac{1}{2\ell+1} \underbrace{Y_{\ell 0}(\theta, \varphi)}_{0 \text{ for } \ell \text{ odd}}$$

$$\int_0^R r'^2 dr' \delta(r'-a) \underbrace{\left( \frac{r'^\ell}{r_0^{2\ell+1}} - \frac{r'^\ell r'^\ell}{R^{2\ell+1}} \right)}_{\neq} P_\ell(0) \sqrt{\frac{2\ell+1}{4\pi}}$$

$P_{2j}(0) = \frac{(-1)^j (2j)!}{2^{2j} (j!)^2}$



Let's consider  $\mathbb{I}$ : (replacing  $l$  by  $z_j$ )

For  $r > a$   $r_c = r'$  and  $r_s = r$ .

$$\mathbb{I} = a^{z_j+2} \left( \frac{1}{r^{z_j+1}} - \frac{r^{z_j}}{R^{z_j+1}} \right)$$

For  $r < a$   $r_c = r$  and  $r_s = r'$ :

$$\mathbb{I} = a^2 \left( \frac{r^{z_j}}{a^{z_j+1}} - \frac{r^{z_j} a^{z_j}}{R^{z_j+1}} \right)$$

Then:

Then we obtain:

$$\phi^{r > a}(r, \theta) = \frac{q}{4\pi\epsilon_0} \sum_{j=0}^{\infty} P_{2j}(\cos\theta) \frac{(-1)^j (2j)!}{z^{2j} (j!)^2} a^{2j}$$

$$\left( \frac{1}{r^{2j+1}} - \frac{r^{2j}}{R^{4j+1}} \right)$$

$$\phi^{r < a}(r, \theta) = \frac{q}{4\pi\epsilon_0} \sum_{j=0}^{\infty} P_{2j}(\cos\theta) \frac{(-1)^j (2j)!}{z^{2j} (j!)^2} r^{2j}$$

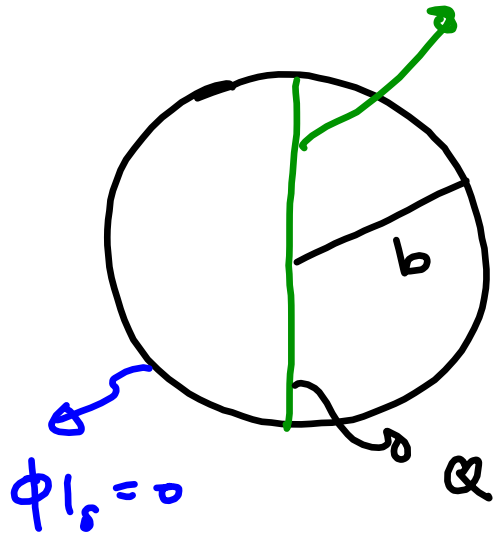
or

$$\left( \frac{1}{a^{2j+1}} - \frac{a^{2j}}{R^{4j+1}} \right)$$

$r_c(r)$  smaller (larger)  
between  $r$  and  $a$ .

$$\phi(r, \theta) = \frac{q}{4\pi\epsilon_0} \sum_{j=0}^{\infty} P_{2j}(\cos\theta) \frac{(-1)^j (2j)!}{z^{2j} (j!)^2} r_c^{2j} \left( \frac{1}{r_c^{2j+1}} - \frac{r_c^{2j}}{R^{4j+1}} \right)$$

Next time:



$$\rho(\vec{r}') = \frac{Q}{2b} \frac{1}{2\pi r'^2} [\delta(\cos\theta' - 1) + \delta(\cos\theta' + 1)]$$

$$\int_V \rho(\vec{r}') d^3r' = Q \text{ (check!)}$$

Find  $\phi(\vec{r})$  inside the grounded sphere.

$$\phi(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int_V G(\vec{r}, \vec{r}') \rho(\vec{r}') d^3r'$$