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Fundamental tensor: metric tensor.

g_{ij} is a tensor that will allow us to lower indices of a tensor. Its inverse

g^{ij} ($g^{ij} \equiv (g_{ij})^{-1}$) will allow us to raise indices of a tensor.

g_{ij} can be any arbitrary symmetric rank 2 tensor that has an inverse ($\det g_{ij} \neq 0$).

But we are going to use the so called "metric tensor" as g_{ij} .

Metric Tensor!

- Let's consider a system S with coordinates $\{x^i\}$ and a system S' with coordinates $\{x'^j\}$.
- In S consider $d\bar{r} = dr^k$.
- In S' consider $d\bar{r}' = dr'^k$.

We know that $d\bar{r} = d\bar{r}'$ only their components are different.

- we can define vectors $\bar{\epsilon}_j$ in S and $\bar{\epsilon}'_j$ in S' so that they are the covariant basis vectors for each system.

- Then

$$d\bar{r} = \bar{\epsilon}_j dx^j \quad \textcircled{1}$$

In cartesian coordinates

$$\bar{\epsilon}_i \equiv \hat{e}_i \text{ (unit vector)}$$

and

$$d\bar{r}' = \bar{\epsilon}'_j dx'^j \quad \textcircled{2}$$

- Then:

$$\text{From } \textcircled{1}: \quad \bar{\epsilon}_j = \frac{d\bar{r}}{dx^j}$$

$$\text{From } \textcircled{2}: \quad \bar{\epsilon}'_j = \frac{d\bar{r}'}{dx'^j}$$

Notice that the "components" of $\bar{\epsilon}_j$ are:

$$(\bar{\epsilon}_j)^k$$

In tensor notation:

$$\bar{r} = r^k \quad \text{then} \quad \bar{\epsilon}_j = (\bar{\epsilon}_j)^k$$

- Now let's define the rank 2 tensor g_{ij} as the direct product of the basis vectors:

$$g_{ij} = \bar{\epsilon}_i \cdot \bar{\epsilon}_j \equiv (\epsilon_i)^k (\epsilon_j)_k$$

scalar
but component
of a tensor
of rank 2.

In cartesian coordinates:

$$\bar{r} = x^k \hat{e}_k = x^1 \hat{e}_1 + x^2 \hat{e}_2 + x^3 \hat{e}_3$$

tensor notation $\bar{r} = r^k \equiv x^k = (x^1, x^2, x^3)$

$$g_{ij} = \bar{\xi}_i \cdot \bar{\xi}_j = \begin{pmatrix} \hat{e}_1 \cdot \hat{e}_1 & \hat{e}_1 \cdot \hat{e}_2 & \hat{e}_1 \cdot \hat{e}_3 \\ \hat{e}_2 \cdot \hat{e}_1 & \hat{e}_2 \cdot \hat{e}_2 & \hat{e}_2 \cdot \hat{e}_3 \\ \hat{e}_3 \cdot \hat{e}_1 & \hat{e}_3 \cdot \hat{e}_2 & \hat{e}_3 \cdot \hat{e}_3 \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} \\ = \mathbb{I}$$

In cartesian coordinates g_{ij} is the identity -

Let's see that $g_{ij} = \bar{\mathbf{E}}_i \cdot \bar{\mathbf{E}}_j$ is a tensor:

In S :

$$ds^2 = d\bar{\mathbf{r}} \cdot d\bar{\mathbf{r}} = \frac{\partial \bar{\mathbf{r}}}{\partial x^i} \cdot \frac{\partial \bar{\mathbf{r}}}{\partial x^j} dx^i dx^j \stackrel{\textcircled{1}}{=} \bar{\mathbf{E}}_i \cdot \bar{\mathbf{E}}_j dx^i dx^j$$

↙
scalar
rank 0

$$= g_{ij} \underbrace{dx^i dx^j}_{\text{vectors: tensors of rank 1}}$$

vectors: tensors of
rank 1

We know that

$$d\bar{\mathbf{r}} = \frac{\partial \bar{\mathbf{r}}}{\partial x^i} dx^i$$

also

$$d\bar{\mathbf{r}} = \bar{\mathbf{E}}_j dx^j$$

$$\text{then } \bar{\mathbf{E}}_j = \frac{\partial \bar{\mathbf{r}}}{\partial x^j} \quad \textcircled{1}$$

g_{ij} has to be a rank 2
tensor using the
quotient rule.

- Inverse of the metric tensor:

Let's define:

$$g^{ij} \text{ so that } g^{ij} g_{jk} = \delta^i_k$$

Then we see that

$$g^{ij} = (g_{jk})^{-1}$$

Now define that

$$g^{ik} A_k = A^i \quad (\text{definition}).$$

then:

$$g_{ik} A^k = g_{ik} [g^{kl} A_l] = \underbrace{g_{ik} g^{kl}}_{\delta_i^l} A_l = A_i$$

We see that g^{ik} raises indices and g_{ik} lowers indices.

What happens in S' ?

$$\text{In } S \quad g^{ij} g_{jk} = \delta^i_k$$

In S' :

$$\begin{aligned} g'^{ij} g'_{jk} &= \frac{\partial x'^i}{\partial x^\alpha} \frac{\partial x'^j}{\partial x^\beta} g^{\alpha\beta} \frac{\partial x^\gamma}{\partial x'^j} \frac{\partial x^\delta}{\partial x'^k} g_{\gamma\delta} = \\ &= \frac{\partial x'^i}{\partial x^\alpha} \frac{\partial x^\gamma}{\partial x'^j} \frac{\partial x'^j}{\partial x^\beta} \frac{\partial x^\delta}{\partial x'^k} g^{\alpha\beta} g_{\gamma\delta} = \\ &= \frac{\partial x'^i}{\partial x^\alpha} \frac{\partial x^\gamma}{\partial x'^k} \underbrace{\frac{\partial x^\beta}{\partial x'^j}}_{\delta^{\beta j}} \underbrace{g^{\alpha\beta} g_{\gamma\delta}}_{\delta^{\alpha\gamma}} = \frac{\partial x'^i}{\partial x^\alpha} \frac{\partial x^\alpha}{\partial x'^k} = \frac{\partial x'^i}{\partial x'^k} = \delta^i_k \end{aligned}$$

Then we saw that in S' :

$$g'^{ij} g'_{jk} = \delta'^i_k \quad \text{because } g'^{ij} \text{ and } g'_{ij} \text{ are tensors.}$$

• Properties of g_{ij} :

• $g_{ij} = g_{ji} \Rightarrow g_{ij}$ is symmetric.

D/ $g_{ij} = \bar{\mathbf{e}}_i \cdot \bar{\mathbf{e}}_j = \bar{\mathbf{e}}_j \cdot \bar{\mathbf{e}}_i = g_{ji}$

then $g_{ij} A^j = g_{ji} A^j = A_i$

Examples:

1) Metric tensor in polar coordinates.

• Consider system S with coordinates

$$(x^1, x^2) \equiv (\rho, \varphi).$$

• and S' with coordinates $(x^{1'}, x^{2'}) \equiv (x, y)$

• Transformations:

$$x^i(x^{i'}):$$

$$\left\{ \begin{array}{l} \rho = (x^2 + y^2)^{1/2} \\ \varphi = \tan^{-1}(y/x) \end{array} \right.$$

$$x^{i'}(x^i):$$

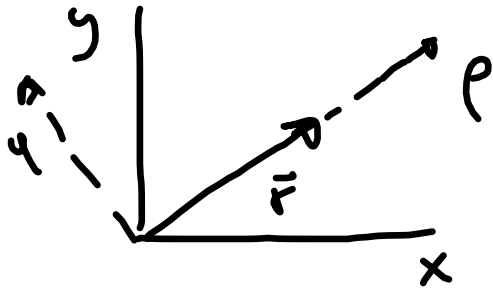
$$\left\{ \begin{array}{l} x = \rho \cos \varphi \\ y = \rho \sin \varphi \end{array} \right.$$

Notice that

$$\begin{aligned}\vec{r}' = (x, y) &= (\rho \cos \varphi, \rho \sin \varphi) = \rho \cos \varphi \hat{e}_1 + \rho \sin \varphi \hat{e}_2 = \\ &= \rho \bar{\vec{e}}_\rho + \varphi \bar{\vec{e}}_\varphi = \vec{r}\end{aligned}$$

$$\text{Then } \cdot \bar{\vec{e}}_\rho = \frac{\partial \vec{r}}{\partial \rho} = (\cos \varphi, \sin \varphi) \quad \text{Cartesian components of } \bar{\vec{e}}_\rho$$

$$|\bar{\vec{e}}_\rho| = 1$$



$$\begin{aligned}\cdot \bar{\vec{e}}_\varphi &= \frac{\partial \vec{r}}{\partial \varphi} = (-\rho \sin \varphi, \rho \cos \varphi) = \\ &= \rho (-\sin \varphi, \cos \varphi) \\ |\bar{\vec{e}}_\varphi| &= \rho \neq 1!\end{aligned}$$

Now

$$g_{ij} = \bar{\xi}_i \cdot \bar{\xi}_j = \begin{pmatrix} \bar{\xi}_\rho \cdot \bar{\xi}_\rho & \bar{\xi}_\rho \cdot \bar{\xi}_\psi \\ \bar{\xi}_\psi \cdot \bar{\xi}_\rho & \bar{\xi}_\psi \cdot \bar{\xi}_\psi \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & \rho^2 \end{pmatrix}$$

g_{ij} is diagonal because S is an orthogonal system of coordinates but $g_{ij} \neq \mathbb{I}$.

then

$$g^{ij} = (g_{ij})^{-1} = \begin{pmatrix} 1 & 0 \\ 0 & 1/\rho^2 \end{pmatrix} \neq g_{ij}$$

2) Oblique system:

Now S : cartesian system $(x^1, x^2) = (x, y)$

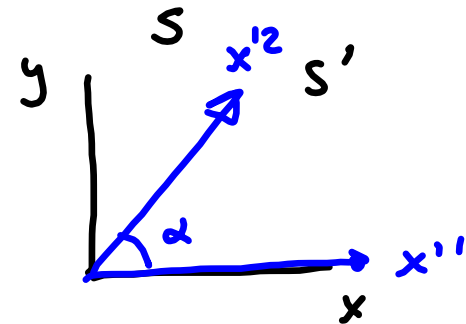
S' : oblique system (x'^1, x'^2)

In S : $g_{ij} = g^{ij} = \mathbb{I}$

In S' : $g'_{ij} = \bar{\xi}'_i \cdot \bar{\xi}'_j$

In homework you found that

$$\bar{\xi}'_1 = \hat{e}_1 = (1, 0) \quad \bar{\xi}'_2 = (\cos \alpha, \sin \alpha)$$



then

$$g'^{ij} = \begin{pmatrix} \bar{\xi}'_1 \cdot \bar{\xi}'_1 & \bar{\xi}'_1 \cdot \bar{\xi}'_2 \\ \bar{\xi}'_2 \cdot \bar{\xi}'_1 & \bar{\xi}'_2 \cdot \bar{\xi}'_2 \end{pmatrix} = \begin{pmatrix} 1 & \cos \alpha \\ \cos \alpha & 1 \end{pmatrix} \text{ symmetric}$$

non diagonal because \wedge it is not orthogonal
 if $\alpha = \frac{\pi}{2} \Rightarrow S'$ orthogonal then g'^{ij} diag.

then

$$g'^{ij} = (g'_{ij})^{-1} = \frac{1}{\sin^2 \alpha} \begin{pmatrix} 1 & \cos \alpha \\ \cos \alpha & 1 \end{pmatrix}$$

Notice that

$$x'^i = g'^{ij} x'^j !$$

$$\begin{aligned} x'_1 &= g'^{11} x'^1 + g'^{12} x'^2 = \\ &= x'^1 + \cos \alpha x'^2 = x'_1 \\ \text{same for } x'_2 \end{aligned}$$

Then we see that for scalar product
we can write:

$$a^\mu b_\mu = g^{\mu\nu} a_\nu b_\mu = g_{\mu\nu} a^\mu b^\nu$$

If $g^{\mu\nu} = \mathbb{I}$ it is irrelevant where you put
the indices but if $g^{\mu\nu} \neq \mathbb{I}$ it really
matters!!