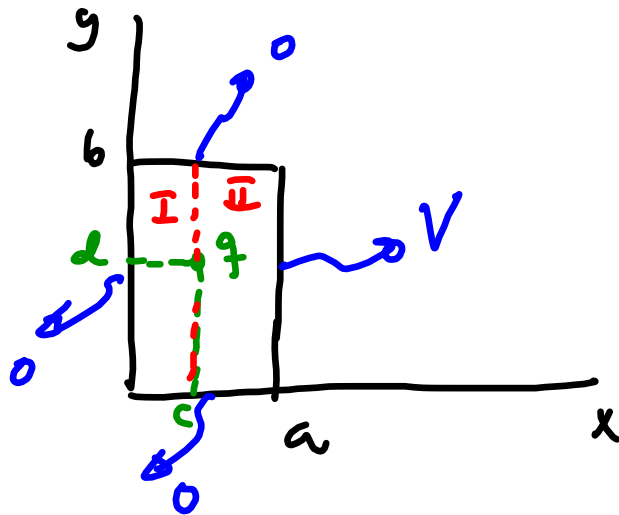


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# Principle of superposition.

Consider this problem:



Find  $\phi(x,y)$  inside  
the box.

- Hard way:
- Separate volume in 2 regions.
  - Propose two solutions.
  - Use b.c.'s to find the coefficients.

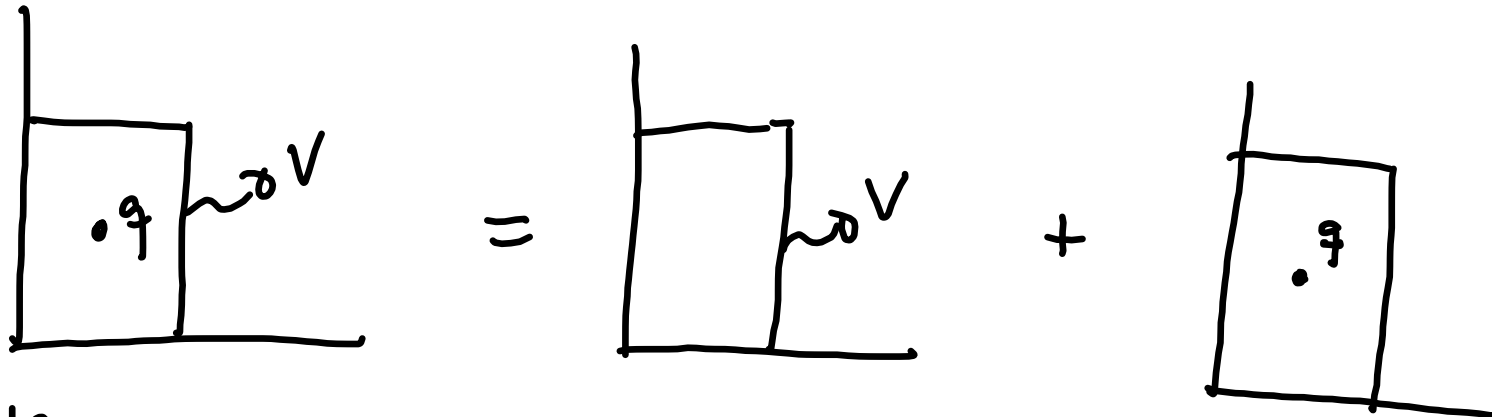
$$\phi^I(x, y) = \sum_{n=1}^{\infty} A_n \sin \frac{n\pi x}{b} \sin \frac{n\pi y}{b}$$

$$\phi^{II}(x, y) = \sum_{n=1}^{\infty} \sin \frac{n\pi y}{b} \left[ B_n e^{\frac{n\pi x}{b}} + C_n e^{-\frac{n\pi x}{b}} \right]$$

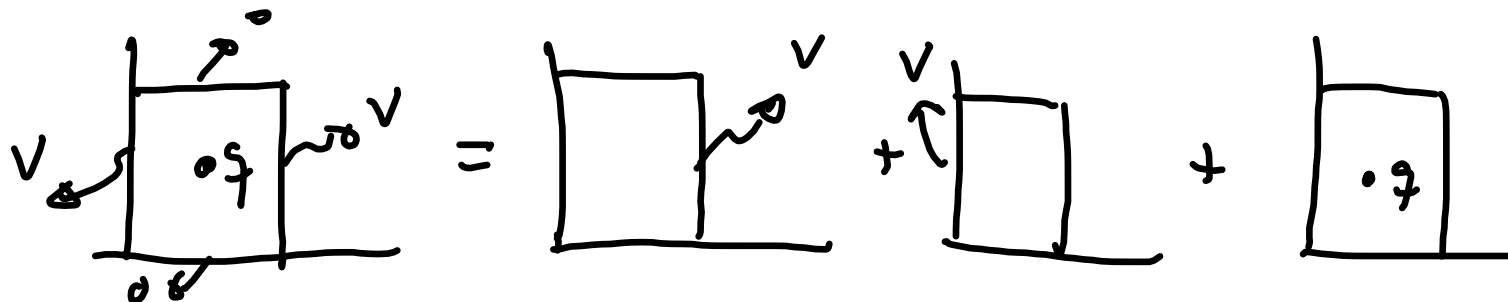
$A_n$ ,  $B_n$  and  $C_n$  can be determined by the  
2 b.c.'s at  $x=c$  and from  $\phi(x=a, y)=V$ .

But it is very tedious.

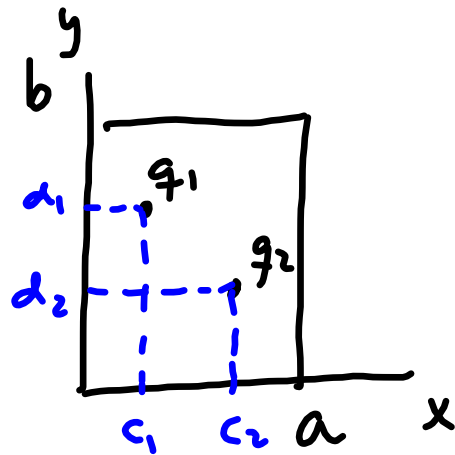
Then it is simpler if you use the principle of superposition:



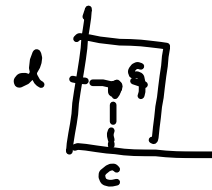
Also



Also:



The solution is the sum of the solution for  $d$  specializing

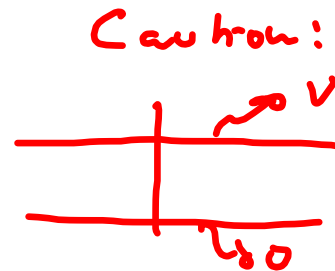
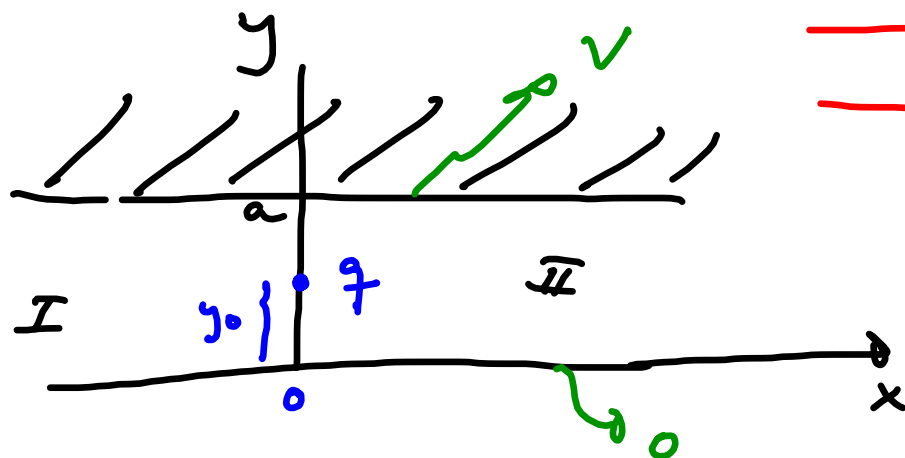


in  $q_1$  at  $(c_1, d_1) + q_2$  at  $(c_2, d_2)$

The sum of the 2 solutions will have 3 different forms:

$\phi^I$	for	$0 \leq x \leq c_1$
$\frac{\phi^I}{\phi^II}$	for	$c_1 \leq x \leq c_2$
$\frac{\phi^I}{\phi^III}$	for	$c_2 \leq x \leq a$ .

Homework:



$$\phi(x,y) \equiv \phi(y)$$

$$\nabla^2 \phi = 0 \Rightarrow$$

$$\frac{d^2 y}{dy^2} = 0$$

$$y(y) = Ay$$

$$y(0) = 0 \text{ then } A = \frac{V}{a}$$

1) Solve

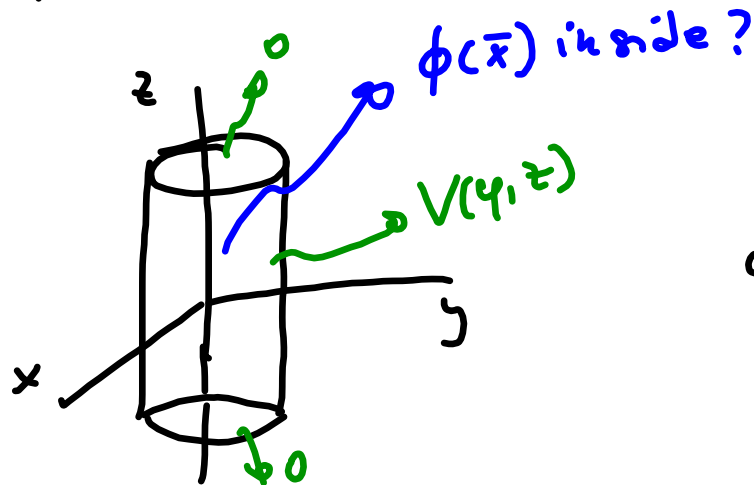
$$\phi^I(x,y) = \sum_{n=1}^{\infty} A_n \sin \frac{n\pi y}{a} e^{\frac{n\pi x}{a}}$$

$$\phi^{II}(x,y) = \sum_{n=1}^{\infty} B_n \sin \frac{n\pi y}{a} e^{-\frac{n\pi x}{a}}$$

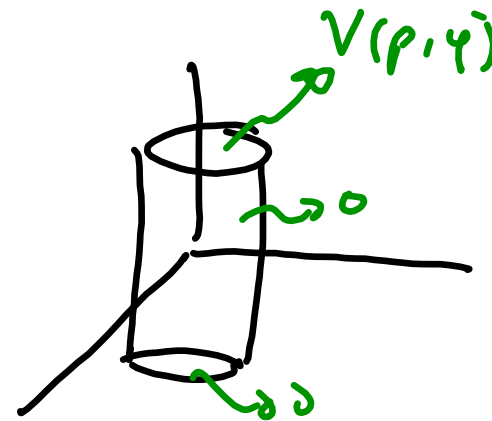
Find  $A_n$  and  $B_n$  from b.c.'s at  $x=0$ .

## Cylindrical Coordinates.

If an electrostatic problem has b.c.'s given on a cylindrical surface then you should solve  $\nabla^2 \phi = 0$  in cylindrical coordinates.



or



$$\nabla^2 \phi(\rho, \varphi, z) = 0$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left( \rho \frac{\partial \phi}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \phi}{\partial \varphi^2} + \frac{\partial^2 \phi}{\partial z^2} = 0 \quad (1)$$

Separation of variables:

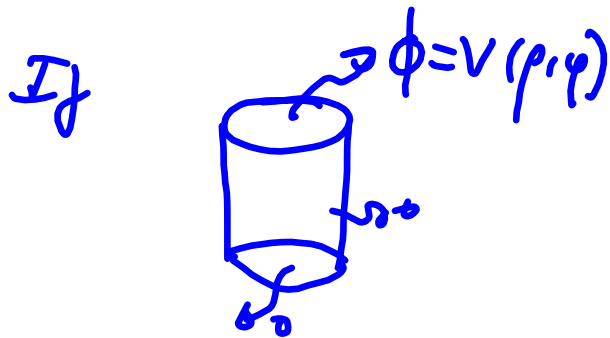
$$\phi(\rho, \varphi, z) = P(\rho) Q(\varphi) Z(z) \quad (2)$$

Plug (2) in (1) and obtain 3 ODE  
and 2 separation constants:

$$Qz \frac{1}{\rho} \frac{\partial}{\partial \rho} \left( \rho \frac{dP}{d\rho} \right) + \frac{Pz}{\rho^2} \frac{d^2 Q}{d\rho^2} + Pz \frac{d^2 z}{dz^2} = 0$$

Divide by  $PQz$ :

$$\underbrace{\frac{1}{P} \frac{1}{\rho} \frac{\partial}{\partial \rho} \left( \rho \frac{\partial P}{\partial \rho} \right) + \frac{1}{\rho^2 Q} \frac{d^2 Q}{d\rho^2}}_{-k^2} + \underbrace{\frac{1}{z} \frac{d^2 z}{dz^2}}_{k^2} = 0$$



we want now periodic solutions for  $z$ .



Then

$$\frac{d^2 z}{dz^2} = k^2 z \Rightarrow$$

$$z(z) \propto e^{\pm kz}$$

Then we are left with:

$$\frac{1}{\rho P} \frac{d}{d\rho} \left( \rho \frac{dP}{d\rho} \right) + \frac{1}{\rho^2 Q} \frac{d^2 Q}{d\rho^2} = -k^2 \quad (3)$$

Multiply (3) by  $\rho^2$ :

$$\frac{\rho}{P} \frac{d}{d\rho} \left( \rho \frac{dP}{d\rho} \right) + \frac{1}{Q} \frac{d^2 Q}{d\rho^2} = -k^2 \rho^2$$

Then

$$\underbrace{\frac{\rho}{P} \frac{d}{d\rho} \left( \rho \frac{dP}{d\rho} \right) + k^2 \rho^2}_{v^2} = - \underbrace{\frac{1}{Q} \frac{\partial^2 Q}{\partial \varphi^2}}_{+v^2}$$

Then

$Q(\varphi) \propto e^{\pm i v \varphi}$  periodic in  $\varphi$ .

For  $P(\rho)$  we have:

$$\rho \frac{d}{d\rho} \left( \rho \frac{dP}{d\rho} \right) + k^2 \rho^2 P - v^2 P = 0$$

$$\rho^2 \frac{d^2 P}{d\rho^2} + \rho \frac{dP}{d\rho} + (k^2 \rho^2 - \nu^2) P = 0$$

Divide by  $\rho^2$ :

$$\frac{d^2 P}{d\rho^2} + \frac{1}{\rho} \frac{dP}{d\rho} + \left( k^2 - \frac{\nu^2}{\rho^2} \right) P = 0$$

Define  $x = k\rho \Rightarrow \rho = \frac{x}{k}$   
 $dx = k d\rho \Rightarrow d\rho = \frac{dx}{k}$

$$k^2 \frac{d^2 P}{dx^2} + \frac{k^2}{x} \frac{dP}{dx} + \left( k^2 - \frac{\nu^2 k^2}{x^2} \right) P = 0$$

Divide by  $k^2$ :

$$\frac{d^2 P}{dx^2} + \frac{1}{x} \frac{dP}{dx} + \left(1 - \frac{\nu^2}{x^2}\right) P = 0$$

Ch. 14

Bessel's eq.

Solution is  $P_\nu(x)$ , Bessel function of order  $\nu$ .

Once we find  $P_\nu(x)$  we can solve  $\nabla^2 \phi = 0$ :

$$\phi(\rho, \varphi, z) = \sum_{\nu, k} a_{\nu, k} P_{\nu, k}(\rho) \phi_\nu(\varphi) z_k(z)$$

Let's solve Bessel's equation:

Frobenius method (see 7.5 in book).

- Solves linear, second order, homogeneous ordinary diff. eqs. of the form:

$$y'' + P(x)y' + Q(x)y = 0$$

where  $y'' = \frac{d^2y}{dx^2}$        $y' = \frac{dy}{dx}$

Most general solution is given by:

$$y(x) = c_1 y_1(x) + c_2 y_2(x)$$

Propose a solution of the form:

$$y(x) = \sum_{\lambda=0}^{\infty} a_{\lambda} x^{k+\lambda} \quad (1)$$

and set equations to find  $a_{\lambda}$  and  $k$ .

If our equation is Bessel's eq:

$$x^2 y'' + x y' + (x^2 - n^2) y = 0 \quad (2)$$

$$\text{From (1): } y'(x) = \sum_{\lambda=0}^{\infty} a_{\lambda} (k+\lambda) x^{k+\lambda-1} \quad (3)$$

$$\text{and } y''(x) = \sum_{\lambda=0}^{\infty} a_{\lambda} (k+\lambda)(k+\lambda-1) x^{k+\lambda-2} \quad (4)$$

Plug ①, ③ and ④ in ②:

$$\sum_{\lambda=0}^{\infty} a_{\lambda} (k+\lambda) (k+\lambda-1) x^{k+\lambda} + \sum_{\lambda=0}^{\infty} a_{\lambda} (k+\lambda) x^{k+\lambda} + \sum_{\lambda=0}^{\infty} a_{\lambda} x^{k+\lambda+2} - \sum_{\lambda=0}^{\infty} a_{\lambda} m^2 x^{k+\lambda} = 0$$

Consider  $\lambda=0$ . This gives us the lowest power of  $x$  in the solution which is  $x^k$ .

We need to ask  $k$  such that the coefficient of  $x^k$  vanishes.

$$a_0 k(k-1) + k a_0 - a_0 m^2 = 0$$

$$(k^2 - m^2) a_0 = 0 \Rightarrow \text{if } a_0 \neq 0 \text{ then}$$

$$k^2 = m^2 \text{ or } \boxed{k = \pm m}$$

indicial  
equation.

$k = m$  and  $k = -m$  will give us the two independent solutions  $y_1(x)$  and  $y_2(x)$ .

Now consider the terms for  $\lambda = 1$ . The lowest power of  $x$  that appears is  $x^{k+1}$  whose coefficient has to vanish to solve the eq.



$$a_1 [(k+1)k + k + 1 - n^2] = 0$$

$$a_1 [k^2 + 2k + 1 - n^2] = 0$$

$$a_1 [(k+1)^2 - n^2] = 0$$

$$a_1 (k+1-n)(k+1+n) = 0$$

You could use this as  
 indicial eq. with  $a_1 \neq 0$   
 but then you'll have to  
 ask  $a_0 = 0$ . In BOTH cases  
 you will obtain the same  
 $y_1$  and  $y_2$  solutions.

Notice that  
 if  $k = n$  or  $-n$  this  
 does not vanish then  
 we need that  $a_1 = 0$ .

If  $n = -\frac{1}{2}$  this vanishes  
 so we will have to do  
 something different in  
 that case.

Now to find the other  $a_\lambda$ 's let's consider  $\lambda = j$  we will obtain the coefficients of  $x^{k+j}$  with  $k=n$  we obtain:

$$a_j [(n+j)(n+j-1) + (n+j) - n^2] + a_{j-2} = 0$$

Let's redefine  $i = j - 2$   $i + 2 = j$

$$a_{i+2} [(n+i+2)(n+i+1) + (n+i+2) - n^2] + a_i = 0$$

Now call  $i = j$ :

$$a_{j+2} = -a_j \frac{1}{(j+2)(2n+j+2)}$$

Then you see that all the  $a_j$ 's will have even  $j$ .

$$a_2 = -\frac{a_0}{2(2n+2)} = -\frac{a_0}{4(n+1)} = -\frac{a_0 n!}{2^2 1! (n+1)!}$$

$$a_4 = -\frac{a_2}{4(2n+4)} = \frac{-a_2}{8(n+2)} = \frac{a_0 n!}{2^4 2! (n+2)!}$$

in general:

$$a_{2p} = \frac{(-1)^p a_0 n!}{2^{2p} p! (n+p)!}$$

So all the  $a_\lambda$  coefficients with even  $\lambda$  are obtained in terms of  $a_0$  while all the coefficients with odd  $\lambda$  are 0.

Then

$$y_1(x) = a_0 x^n \left[ 1 - \frac{n! x^2}{2^2 1! (n+1)!} + \frac{n! x^4}{2^4 2! (n+2)!} + \dots \right]$$

$$= a_0 \sum_{j=0}^{\infty} \frac{(-1)^j n! x^{n+2j}}{2^{2j} j! (n+j)!} = a_0 2^n n! \underbrace{\sum_{j=0}^{\infty} \frac{(-1)^j}{j! (n+j)!} \left(\frac{x}{2}\right)^{n+2j}}_{J_n(x)}$$

$J_n(x)$ : Bessel function of order  $n$   
(Ch. 14) for all  $n \neq -\frac{1}{2}$ .

If  $k = -n$  you will obtain  $y_2(x)$ .