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Dirac delta function in 3D:

$$\delta(\bar{x} - \bar{X}) = \delta(x_1 - X_1) \delta(x_2 - X_2) \delta(x_3 - X_3)$$

$$\int_V \delta(\bar{x} - \bar{X}) d^3x = \begin{cases} 1 & \text{if } \bar{X} \text{ is in } V. \\ 0 & \text{if } \bar{X} \text{ is not in } V. \end{cases}$$

Examples: $\rho(\bar{x}) = q \delta(\bar{r} - \bar{x})$ for a point charge.
we see that

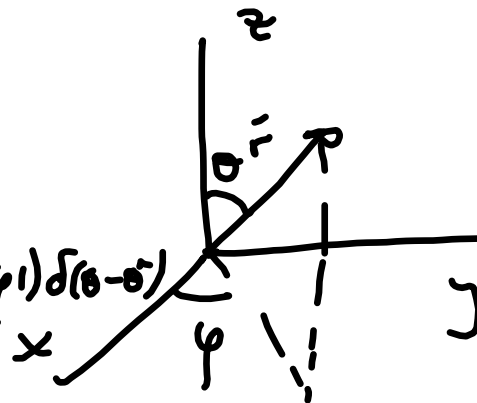
$$\int_V \rho(\bar{x}) d^3x = q \int_V \delta(\bar{r} - \bar{x}) d^3x = q \quad \text{if } V = \text{all space.}$$

Since many times boundary conditions are given on spheres or cylinders rather than planes we need to express $\delta(\vec{x} - \vec{x}')$ in different systems of coordinates:

Spherical coordinates: (r, θ, φ)

We know that $\int_{\text{all space}} \delta(\vec{r}) dV = 1$

$$1 = \int_0^\infty \int_0^\pi \int_0^{2\pi} A \underbrace{r^2 \sin\theta \, dr \, d\theta \, d\varphi}_{dV} \delta(r-r') \delta(\varphi-\varphi') \delta(\theta-\theta')$$



$$= \underbrace{\int_0^{\infty} A \delta(r-r') r^2 dr}_{A r'^2} \underbrace{\int_0^{\pi} \delta(\theta-\theta') \sin \theta d\theta}_{\sin \theta'} \underbrace{\int_0^{2\pi} \delta(\varphi-\varphi') d\varphi}_1$$

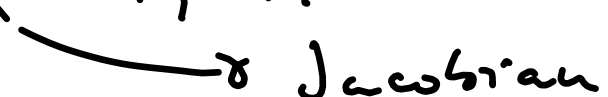
$$\text{Then } A = \frac{1}{r'^2 \sin \theta'}$$

Then

$$\delta(\vec{r}-\vec{r}') = \frac{\delta(r-r') \delta(\theta-\theta') \delta(\varphi-\varphi')}{r'^2 \sin \theta'} \equiv \frac{\delta(r-r') \delta(\theta-\theta') \delta\varphi}{r^2 \sin \theta}$$

\mathcal{I} n general:

$$\delta(\bar{x} - \bar{x}') = \frac{\delta(\xi_1 - \xi_1') \delta(\xi_2 - \xi_2') \delta(\xi_3 - \xi_3')}{|J(x_i, \xi_j)|}$$


 Jacobian

$$J(x_i, \xi_j) = \begin{pmatrix} \frac{\partial x_1}{\partial \xi_1} & \frac{\partial x_1}{\partial \xi_2} & \frac{\partial x_1}{\partial \xi_3} \\ \frac{\partial x_2}{\partial \xi_1} & \frac{\partial x_2}{\partial \xi_2} & \frac{\partial x_2}{\partial \xi_3} \\ \frac{\partial x_3}{\partial \xi_1} & \frac{\partial x_3}{\partial \xi_2} & \frac{\partial x_3}{\partial \xi_3} \end{pmatrix} \quad x_i = x_i(\xi_j)$$

$$\mathbb{R}^3 \{ \hat{i} \} = (r, \theta, \varphi)$$

$$x = r \sin \theta \cos \varphi$$

$$y = r \sin \theta \sin \varphi$$

$$z = r \cos \theta$$

$$J = \begin{vmatrix} \sin \theta \cos \varphi & r \cos \theta \cos \varphi & -r \sin \theta \sin \varphi \\ \sin \theta \sin \varphi & r \cos \theta \sin \varphi & +r \sin \theta \cos \varphi \\ \cos \theta & -r \sin \theta & 0 \end{vmatrix} = r^2 \sin \theta$$

$$\text{If } \{r_i\} = (r, \omega\theta, \varphi)$$

$$\delta(\vec{r} - \vec{r}') = \frac{\delta(r - r') \delta(\omega\theta - \omega\theta') \delta(\varphi - \varphi')}{r^2}$$

Expansion of $\delta(x - x')$ in terms of orthogonal functions. [Review: Ch. 5.1.6]

$\{\varphi_n(x)\}$ are orthogonal in the interval (a, b) :

$$\int_a^b \varphi_m^*(x) \varphi_n(x) dx = \delta_{m,n} C$$

C : constant
if $C = 1$

$\{\varphi_n(x)\}$: orthonormal!

Any "good behaved" function $f(x)$ can be expanded in terms of $\{\varphi_n(x)\}$:

$$f(x) = \sum_{n=0}^{\infty} b_n \varphi_n(x)$$

with

$$b_n = \int_a^b f(x) \varphi_n^*(x) dx$$

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$$\int_a^b \varphi_m^*(x) f(x) dx = \sum_{n=0}^{\infty} b_n \underbrace{\int_a^b \varphi_m^*(x) \varphi_n(x) dx}_{\delta_{m,n}}$$

Now let's do the same for $\delta(x-t)$:

$$\delta(x-t) = \sum_{n=0}^{\infty} a_n(t) \varphi_n(x) \quad (1)$$

Let's find $a_n(t)$:

$$a_n(t) = \int_a^b \varphi_n^*(x) \delta(x-t) dx = \quad (1)$$

$\varphi_n^*(t)$

$$= \int_a^b \varphi_n^*(x) \sum_{m=0}^{\infty} a_m(t) \varphi_m(x) dx =$$

$$= \sum_{m=0}^{\infty} a_m(t) \int_a^b \underbrace{\varphi_m^*(x) \varphi_m(x)}_{\delta_{n,m}} dx =$$
$$= a_m(t)$$

Then $a_m(t) = \varphi_m^*(t)$

$$\therefore \delta(x-t) = \sum_{m=0}^{\infty} \varphi_m^*(t) \varphi_m(x)$$

Example: $\{\varphi_n(x)\} = \left\{ \sqrt{\frac{2}{a}} \cos \frac{2n\pi x}{a} \right\}$ in $(-\frac{a}{2}, \frac{a}{2})$
 orthonormal since

$$\frac{2}{a} \int_{-a/2}^{a/2} \cos \frac{2n\pi x}{a} \cos \frac{2m\pi x}{a} dx = \delta_{m,n}$$

Then:

$$\delta(x-t) = \sum_{n=0}^{\infty} a_n(t) \cos \frac{2n\pi x}{a}$$

Multiply both sides by $\sqrt{\frac{2}{a}} \cos \frac{2m\pi x}{a}$ and integrate over x on $(-\frac{a}{2}, \frac{a}{2})$ to find $a_n(t)$.

$$\sqrt{\frac{2}{a}} \int_{-a/2}^{a/2} \delta(x-t) \cos \frac{2\pi n x}{a} dx = \sum_{n=0}^{\infty} a_n(t) \int_{-a/2}^{a/2} \cos \frac{2\pi n x}{a} dx$$

$\underbrace{\hspace{15em}}_{\cos \frac{2\pi n t}{a}} \qquad \underbrace{\hspace{15em}}_{\frac{2}{a} \cdot \cos \frac{2\pi n x}{a} dx}$

$$a_n(t) = \sqrt{\frac{2}{a}} \cos \frac{2\pi n t}{a}$$

$$\delta(x, t) = \frac{2}{a} \sum_{n=0}^{\infty} \cos \frac{2\pi n t}{a} \cos \frac{2\pi n x}{a}$$

Tensor analysis (Ch. 4)

A tensor of rank k in dimension N has N^k elements.

$k=0 \Rightarrow$ scalar 1 component

$k=1 \Rightarrow$ vector N components

\vdots

So far we know that:

- Tensors of rank 0 are invariant under transformations.

- Tensors of rank 1 transform according to:

$$A^{i'} = \frac{\partial x^{i'}}{\partial x^j} A^j$$

for contravariant components.

$$B_{j'} = \frac{\partial x^i}{\partial x^{i'}} B_i$$

for covariant components.

$x^{i'}$: components in S'

x^i : components in S

We use Einstein's notation summing over repeated indices.

Now let's study tensors with $k > 1$:

Direct product or outer product of tensors:

$$A^{\mu} B^{\nu} = C^{\mu\nu}$$

$C^{\mu\nu}$ has rank 2
and it results from
the direct product
of A^{μ} with B^{ν} .

In 2D:

$$A^{\mu} = (A^1, A^2) \quad B^{\nu} = (B^1, B^2)$$

$$C^{\mu\nu} = \begin{pmatrix} A^1 B^1 & A^1 B^2 \\ A^2 B^1 & A^2 B^2 \end{pmatrix} = \begin{pmatrix} C^{11} & C^{12} \\ C^{21} & C^{22} \end{pmatrix}$$

How does $C^{\mu\nu}$ transform as we go from S to S' ?

$$\begin{aligned}
 C'^{\mu\nu} &= A'^{\mu} B'^{\nu} = \frac{\partial x'^{\mu}}{\partial x^{\alpha}} A^{\alpha} \frac{\partial x'^{\nu}}{\partial x^{\beta}} B^{\beta} = \\
 &= \frac{\partial x'^{\mu}}{\partial x^{\alpha}} \frac{\partial x'^{\nu}}{\partial x^{\beta}} \underbrace{A^{\alpha} B^{\beta}}_{C^{\alpha\beta}} = \\
 &= \frac{\partial x'^{\mu}}{\partial x^{\alpha}} \frac{\partial x'^{\nu}}{\partial x^{\beta}} C^{\alpha\beta}
 \end{aligned}$$

$C^{\alpha\beta}$ is our prototype contravariant tensor of rank 2.

Notice that \mathbb{I} can form the outer product of covariant and contravariant vectors then I can define:

$$\left. \begin{aligned}
 C^{\mu}_{\nu} &= A^{\mu} B_{\nu} \\
 C_{\mu}^{\nu} &= A_{\mu} B^{\nu}
 \end{aligned} \right\} \text{mixed rank tensors}$$

$$\left. \begin{aligned}
 C_{\mu\nu} &= A_{\mu} B_{\nu}
 \end{aligned} \right\} \text{covariant tensor of rank 2.}$$

You can show that

$$C'^{\mu}_{\nu} = \frac{\partial x'^{\mu}}{\partial x^{\alpha}} \frac{\partial x^{\beta}}{\partial x'^{\nu}} C^{\alpha}_{\beta}$$

$$C'^{\mu\nu} = \frac{\partial x^{\alpha}}{\partial x'^{\mu}} \frac{\partial x'^{\nu}}{\partial x^{\beta}} C_{\alpha\beta}$$

$$C'_{\mu\nu} = \frac{\partial x^{\alpha}}{\partial x'^{\mu}} \frac{\partial x^{\beta}}{\partial x'^{\nu}} C_{\alpha\beta}$$

Higher rank tensors can be constructed by performing the outer product of several vectors or of tensors of lower rank:

$$C^{\mu\nu} \beta^\sigma = T^{\mu\nu\sigma}$$

Contravariant tensor
of rank 3.

The number of indices defines the rank of the tensor.

Ex: in 2D

$$T^{\mu\nu\sigma} = \begin{array}{cc|cc} T^{111} & T^{112} & T^{211} & T^{212} \\ T^{121} & T^{122} & T^{221} & T^{222} \end{array} \quad \left(\begin{array}{c} 2 \times 2 \times 2 \\ \text{cube} \end{array} \right)$$

8 components.

How does $T^{\mu\nu\sigma}$ transform?

$$T'^{\mu\nu\sigma} = C'^{\mu\nu} B'^{\sigma} = \frac{\partial x'^{\mu}}{\partial x^{\alpha}} \frac{\partial x'^{\nu}}{\partial x^{\beta}} C^{\alpha\beta} \frac{\partial x'^{\sigma}}{\partial x^{\gamma}} B^{\gamma}$$

$$= \frac{\partial x'^{\mu}}{\partial x^{\alpha}} \frac{\partial x'^{\nu}}{\partial x^{\beta}} \frac{\partial x'^{\sigma}}{\partial x^{\gamma}} T^{\alpha\beta\gamma}$$

sums over
 α, β and γ

Other examples:

$$B^{\nu}_{\mu} \alpha T_{\delta\epsilon}^{\rho} = M^{\nu}_{\mu} \alpha \delta_{\epsilon}^{\rho} \quad \text{mixed tensor.}$$

$$M^{\nu}_{\mu} \alpha \delta_{\epsilon}^{\rho} = \frac{\partial x'^{\nu}}{\partial x^{\alpha}} \frac{\partial x^{\beta}}{\partial x'^{\mu}} \frac{\partial x^{\gamma}}{\partial x^{\delta}} \frac{\partial x^{\epsilon}}{\partial x'^{\rho}} M^{\alpha}_{\beta\gamma\delta}$$