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Last time:

$W_N(n_1)$ probability of taking n_1 steps to the right after N steps.

$$\langle n_1 \rangle = Np \quad \langle n_2 \rangle = Nq$$

$$\overline{\Delta n_1^2} = \langle n_1^2 \rangle - \langle n_1 \rangle^2 = Npq \quad \text{and} \quad \Delta^* n_1 = \sqrt{Npq}$$

Now let's find $\langle m \rangle$ and $\overline{\Delta m^2}$ for $P_N(m)$

$$m = n_1 - n_2 \quad N = n_1 + n_2 \therefore m = 2n_1 - N$$

$$\begin{aligned} \langle m \rangle &= 2\langle n_1 \rangle - N = 2Np - N = N(2p - 1) = \\ &= N(p + \frac{p-1}{q}) = N(p-q) \end{aligned}$$

\checkmark
probability
of being at
a distance
 m from the
start after N
steps

$$\Delta m = m - \langle m \rangle = 2n_1 - N - N(p-q) = 2n_1 - N - N(2p-1),$$

$$= 2(n_1 - Np) - N + N = 2(n_1 - \langle n_1 \rangle) = 2\Delta n_1^2$$

$$\overline{\Delta m^2} = 4 \overline{\Delta n_1^2} = 4Npq$$

$$\Delta^* m = 2\sqrt{Npq}$$

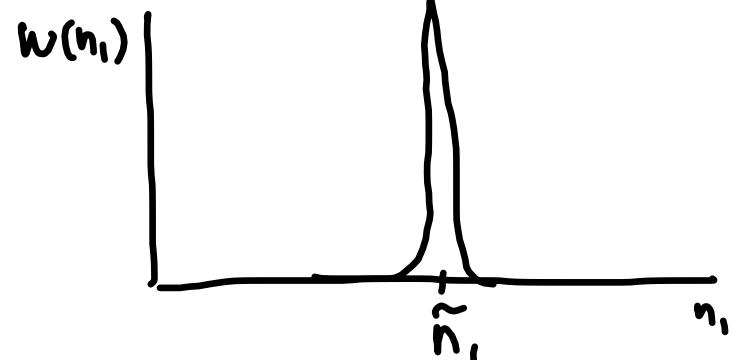
$$\text{If } p=q \text{ then } \langle m \rangle = 0 \text{ and } \Delta^* m = \sqrt{N}$$

Large N limit of the probability distribution.

If N is very large we'll see that $N_N(n_i)$ and $P_n(n_i)$ become gaussian except for the case in which $\rho \ll 1$ or $\eta \ll 1$. Then you'll get a Poisson distribution that does not have a finite dispersion (see Hw #1).

$$\text{If } N \rightarrow \infty \quad \frac{\Delta n_i^+}{\langle n_i \rangle} = \sqrt{\frac{q}{p}} \frac{1}{\sqrt{N}}$$

Now



Since $W(n_i)$ has a sharp maximum at $n_i = \tilde{n}_i$, we can make a Taylor expansion of $W(n_i)$ about \tilde{n}_i .

Notice that close to \tilde{n}_i , $|W(n_i + 1) - W(n_i)| \ll W(n_i)$
so n_i can be considered continuous.

We know that

$$\frac{\partial W}{\partial n_i} \Big|_{\tilde{n}_i} = 0 \quad \text{or} \quad \frac{d \ln W}{d n_i} \Big|_{\tilde{n}_i} = 0$$

A series for
 $\ln W$ converges faster
than for W .

$$\ln W(n_1) = \underbrace{\ln \tilde{W}}_{\tilde{W}} + \underbrace{\frac{d \ln W}{d n_1} \Big|_{\tilde{n}_1}}_0 \gamma + \frac{1}{2} \underbrace{\frac{d^2 \ln W}{d n_1^2} \Big|_{\tilde{n}_1}}_{-\beta_2} \gamma^2 + \dots$$

Then

$$W(n_1) \approx \tilde{W} e^{-\frac{1}{2} \beta_2 \gamma^2} \quad \textcircled{1}$$

Also we know that

$$W_N(n_1) = \frac{N!}{n_1! (N-n_1)!} \rho^{n_1} q^{N-n_1}$$

$$\begin{aligned} \ln W_N(n_1) &= \ln N! - \ln n_1! - \ln (N-n_1)! + n_1 \ln \rho + \\ &\quad (N-n_1) \ln q \quad \textcircled{2} \end{aligned}$$

For large n , $\ln n!$ can be considered continuous.

$$\frac{d \ln n!}{dn} \approx \frac{\ln(n+1)! - \ln(n)!}{1} = \ln\left[\frac{(n+1)!}{n!}\right] = \ln(n+1) \approx \\ \approx \ln n \quad \text{since } n \gg 1. \quad ③$$

Then

$$\frac{d \ln W}{dn_1} \approx -\ln n_1 + \ln(N-n_1) + \ln p - \ln q$$

We can find \tilde{n}_1 since $\frac{d \ln W}{dn_1}|_{\tilde{n}_1} = 0$

$$0 = \ln \underbrace{\frac{(N-\tilde{n}_1)}{\tilde{n}_1}}_1 \frac{p}{q}$$

Then $N - \tilde{n}_1 = \tilde{n}_1 \frac{q}{p}$

$$N = \tilde{n}_1 \left(\frac{q}{p} + 1 \right) = \tilde{n}_1 \frac{q+1}{p}$$

$\tilde{n}_1 = Np$

D180

$$\begin{aligned} B_2 &= \left. \frac{d^2 \ln W}{du_1^2} \right|_{\tilde{n}_1} = -\frac{1}{n_1} \left|_{\tilde{n}_1} \right. - \frac{1}{n-n_1} \left|_{\tilde{n}_1} \right. = -\frac{1}{Np} - \frac{1}{N-Np} \\ &= -\frac{1}{Np} - \frac{1}{N(\underbrace{1-p}_{q})} = -\frac{q+p}{Npq} = -\frac{1}{Npq} \end{aligned}$$

$$|B_2| = \frac{1}{Npq} \equiv \frac{1}{\Delta n^2}$$

To find \tilde{W} we request normalization of $W(u_1)$:

$$1 = \int_{-\infty}^{\infty} W(u_1) du_1 \approx \int_{-\infty}^{\infty} \tilde{W}(\tilde{n}_1 + \gamma) d\gamma = \int_{-\infty}^{\infty} \tilde{W} e^{-\frac{|B_2|\gamma^2}{2}} d\gamma \stackrel{\text{because } \tilde{W}(\gamma) \geq 0 \text{ for } \gamma \text{ away from } \tilde{n}_1}{=} 1$$

We know that

$$\int_{-\infty}^{\infty} e^{-\alpha x^2} dx = \sqrt{\frac{\pi}{\alpha}} \quad \text{then}$$

$$I = \tilde{W} \sqrt{\frac{2\pi}{|B_2|}} = 1 \Rightarrow \tilde{W} = \sqrt{\frac{|B_2|}{2\pi}} = \sqrt{\frac{1}{2\pi Npq}}$$

Then

$$W(n_1) = \sqrt{\frac{1}{2\pi Npq}} e^{-\frac{1}{2} \frac{(n_1 - np)^2}{2Npq}}$$

or to remove N :

$$W(n_1) = \sqrt{\frac{1}{2\pi (\sigma n_1)^2}}$$

$$e^{-\frac{1}{2} \frac{(n_1 - \langle n_1 \rangle)^2}{2(\sigma n_1)^2}}$$

(4)

Gaussian
shape.

Gaussian probability distribution

From ④ we can also obtain the large N limit of $P_N(m)$:

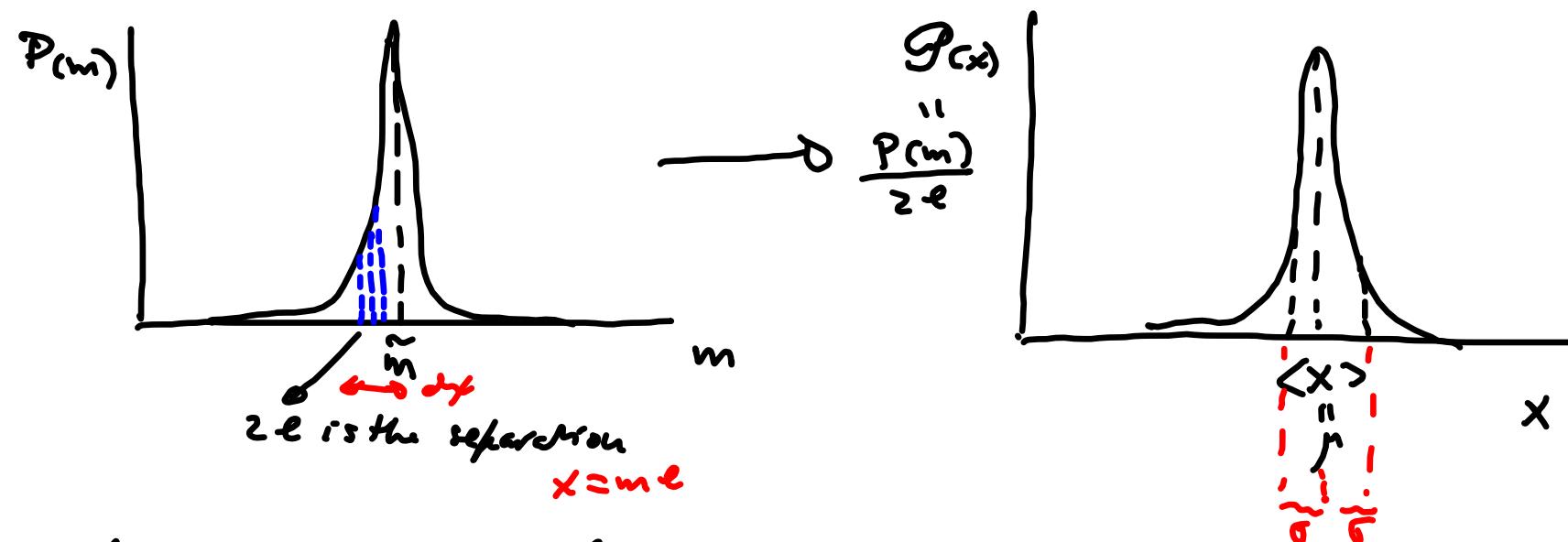
$$m = 2n_1 - N \Rightarrow n_1 = \frac{m+N}{2}$$

$$P(m) = W\left(\frac{N+m}{2}\right) = [2\pi Npq]^{-1/2} e^{-\frac{[m-N(p-q)]^2}{8Npq}}$$

$$= [2\pi Npq]^{-1/2} e^{-\frac{m-\langle m \rangle}{2\langle \Delta m^2 \rangle}}$$

We can express $P(m)$ in terms of the actual displacement $x=ml$ l : size of the step.

If $\ell \ll L$ L : length of interest such as the length of a container.
 Then x can be considered ℓ : of the order of the mean free path.



in each interval dx
there are $\frac{dx}{2e}$ values of m .

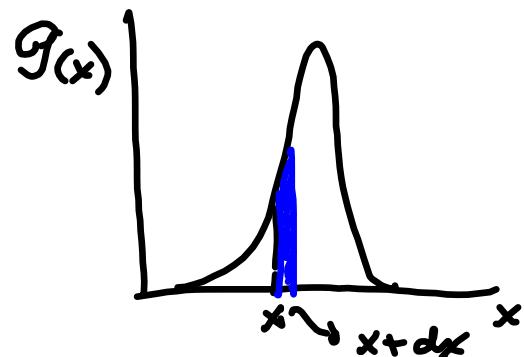
$$\mu \in (p-\gamma)N^p$$

$$\sigma \equiv 2\sqrt{Np\gamma} +$$

$P(x)$ is a probability density.

$$P(x) dx = \frac{P(m)}{2e} dx = \frac{1}{2\pi\sigma} e^{-\frac{(x-\mu)^2}{2\sigma^2}} dx$$
Gaussian distribution

$P(x) dx$ is the probability of finding the walker in the interval $(x, x+dx)$ and it is given by the blue area in the figure.



Notice that if $p \neq q$ the gaussian approximation is still valid when $N \rightarrow \infty$ as long as p is not too small. If $p \ll 1$ then $\langle \Delta n_i^2 \rangle = \sqrt{\frac{q}{p}} \sqrt{\frac{1}{n}}$ and it does not go to zero. You'll get Poisson's distribution.

Probability Distributions for more than 1 random variables: discrete case.

$$P(u_i, v_j)$$

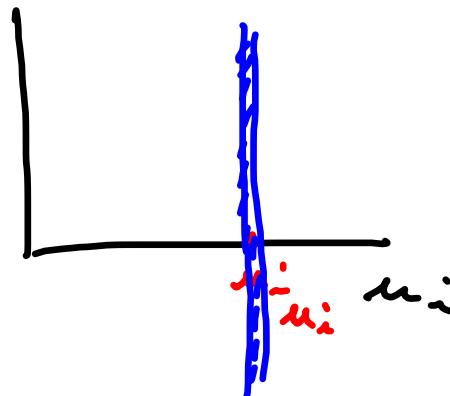
prob that a random variable u takes the value u_i of M possible values and a random variable v takes the value v_j of N possible values,
 $i=1, \dots, M$ and $j=1, \dots, N$.

. $P(u_i, v_j)$ is normalized

$$\sum_{i=1}^M \sum_{j=1}^N P(u_i, v_j) = 1$$

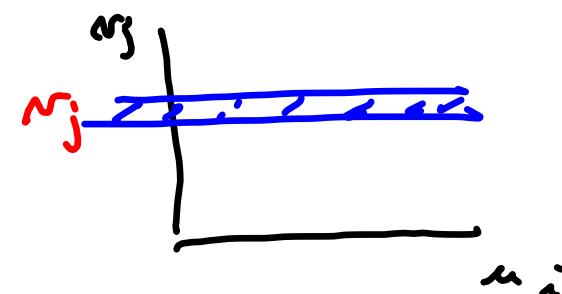
- $P_u(u_i) = \sum_{j=1}^N P(u_i, v_j) \textcircled{S} v_j$

prob of selecting v_i regardless of the value of v .



- $P_v(v_j) = \sum_{i=1}^M P(u_i, v_j)$

prob of selecting v_j regardless of the value of u .



- Then

$$\sum_{i=1}^N P_u(u_i) \textcircled{S} \sum_{i=1}^N \sum_{j=1}^M P(u_i, v_j) = 1$$

- If u and v are independent variables, or uncorrelated variables then

$$P(u_i, v_j) = P_u(u_i) P_v(v_j)$$

- $\bar{F}_{(u,v)} = \sum_{i,j} F(u_i, v_j) P(u_i, v_j)$

- $\bar{f}(u) = \sum_{i,j} f(u_i) P(u_i, v_j) = \sum_{i=1}^M f(u_i) \underbrace{\sum_{j=1}^N P_{u_i, v_j}}_{P_u(u_i)}$

⑤ $\sum_{i=1}^M f(u_i) P_u(u_i)$

- $\overline{F + G} = \bar{F} + \bar{G}$

- $\overline{f(u) g(v)} = \bar{f}(u) \bar{g}(v)$

Probability Density (continuous case)

If the length of interest is much larger than the interval between the values of the discrete random variables we can define a continuous probability density.

$$\frac{P(u)}{\delta u} du = P(u) du$$

probability density.

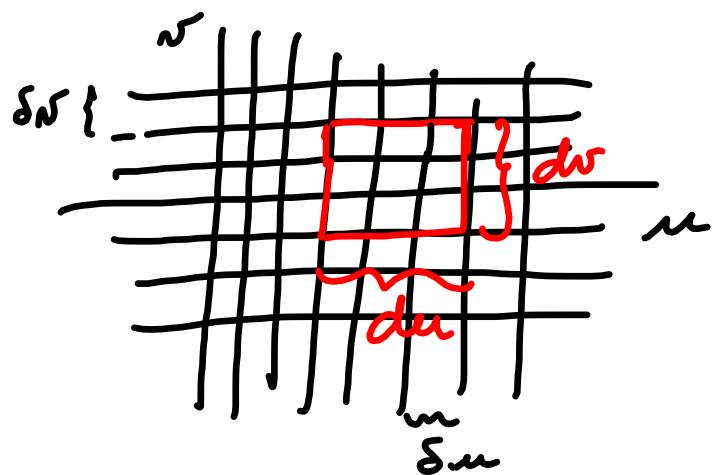
$\delta u \ll du \ll U$

Now sums become integrals:

$$\sum_{i=1}^N P(x_i) = 1 \rightarrow \int_{a_1}^{a_2} P(x) dx = 1$$

$a_1 - a_2$
range of x

Beyond 1 variable:



$$\underbrace{P(x, v) dx dv}_{\text{probability that } x \text{ and } v \text{ fall in } dA = dx dv} = \frac{P(x, v) dx dv}{\delta x \delta v}$$

Now

$$\int_{a_1}^{a_2} \int_{b_1}^{b_2} du dv P(u, v) = 1 \quad \text{etc.}$$