

$$H = \hbar\omega \left( a_- a_+ - \frac{1}{2} \right)$$

Or, using the other order,  
we get (left as exercise):

$$H = \hbar\omega \left( a_+ a_- + \frac{1}{2} \right)$$

**Theorem:** if  $\psi$  satisfies  $H\psi = E\psi$ ,  
then  $H(a_+\psi) = (E + \hbar\omega)(a_+\psi)$

$$\begin{aligned} H(a_+\psi) &= \hbar\omega \left( a_+ a_- + \frac{1}{2} \right) (a_+\psi) = \hbar\omega \left( a_+ a_- a_+ + \frac{1}{2} a_+ \right) \psi \\ &= \hbar\omega a_+ \left( a_- a_+ + \frac{1}{2} \right) \psi = a_+ \left[ \hbar\omega \left( a_- a_+ + 1 - \frac{1}{2} \right) \psi \right] \\ &= a_+ (H + \hbar\omega) \psi = a_+ (E + \hbar\omega) \psi = (E + \hbar\omega) (a_+\psi). \end{aligned}$$

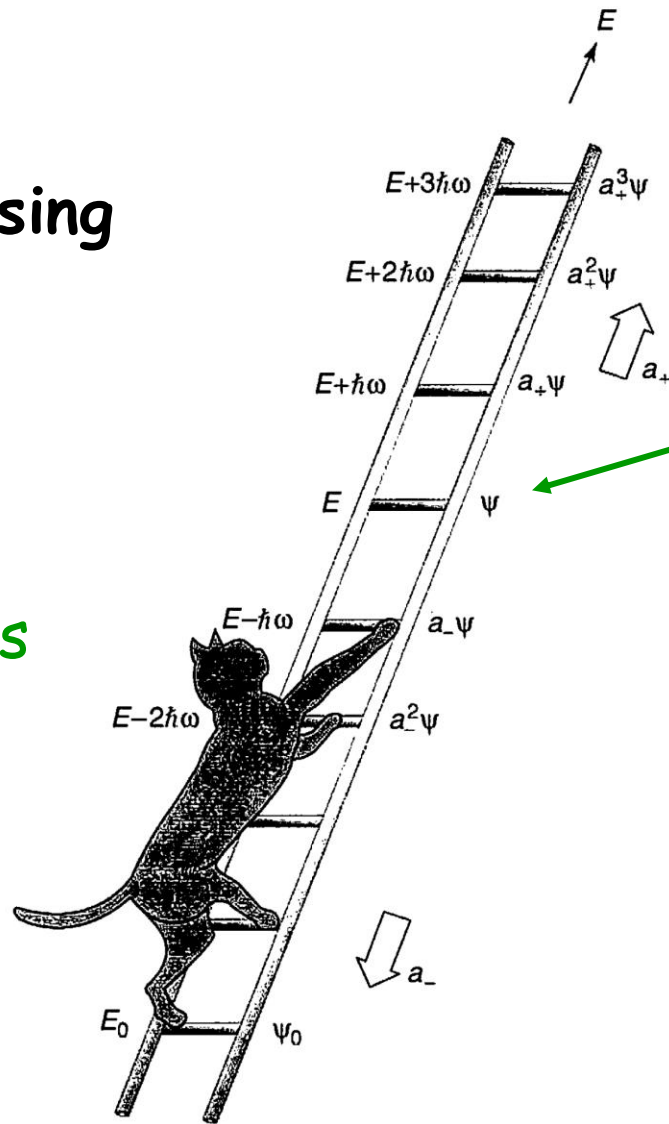
If I know one solution, I know another solution ...

Another theorem: if  $\psi$  satisfies  $H\psi = E\psi$ , then  $H(a_-\psi) = (E - \hbar\omega)(a_-\psi)$  (left as exercise)

$a_+$  is the raising operator

Equally spaced energies. This is why it works.

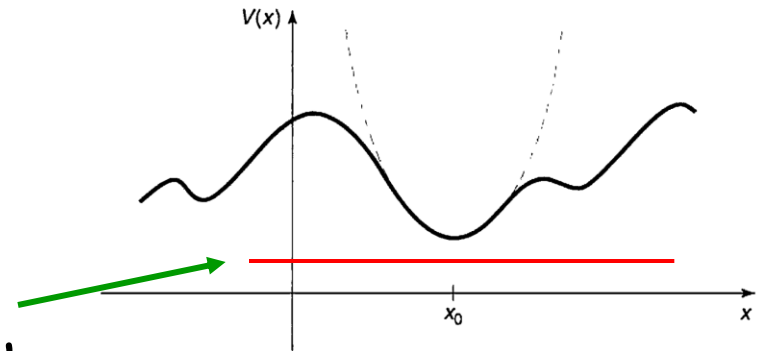
Not general: valid for oscillators.



If I am given one solution, I get infinite more.

$a_-$  is the lowering operator

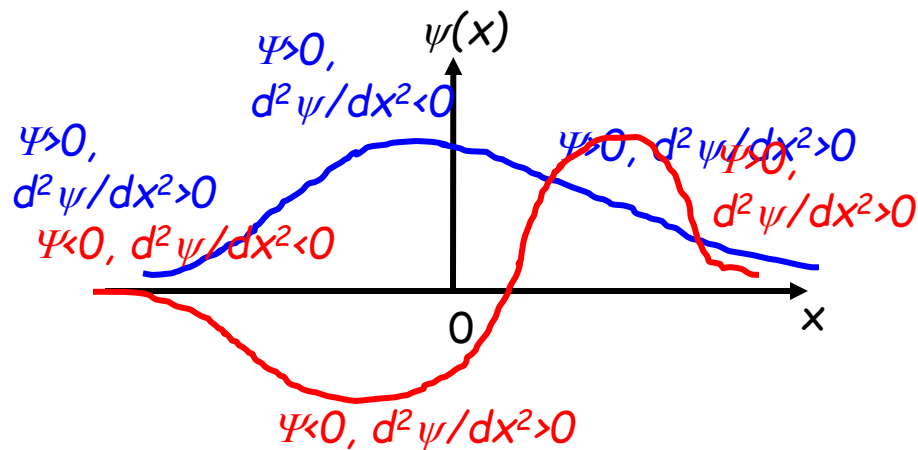
However, there is a problem:  
the energy cannot continue  
going down!



Theorem:  $E$  less than  $V(x)$  cannot happen

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} + V \psi = E \psi \quad \longrightarrow \quad \frac{d^2 \psi}{dx^2} = \frac{2m}{\hbar^2} \underbrace{[V(x) - E]}_{>0} \psi$$

If this is  $>0$  for all  $x$ , then  $\psi$  and  $d^2 \psi / dx^2$  must have same sign.



Because the energy cannot  
continue going down forever,  
the chain down must stop...

Eventually once true ground state  $\psi_0$  is reached, then  $a_-\psi_0$  must be 0. We can use this condition to find  $\psi_0$ .

$$\underbrace{\frac{1}{\sqrt{2\hbar m\omega}} \left( \hbar \frac{d}{dx} + m\omega x \right)}_{a_-} \psi_0 = 0 \quad \longrightarrow \quad \frac{d\psi_0}{dx} = -\frac{m\omega}{\hbar} x \psi_0$$

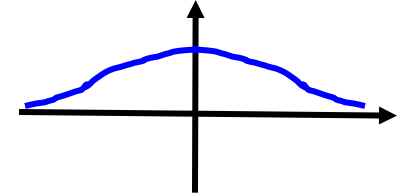
$$\int \frac{d\psi_0}{\psi_0} = -\frac{m\omega}{\hbar} \int x dx \quad \Rightarrow \quad \ln \psi_0 = -\frac{m\omega}{2\hbar} x^2 + \text{constant},$$

$$\psi_0(x) = Ae^{-\frac{m\omega}{2\hbar} x^2}.$$

After normalization:  
(left as exercise)

$$\psi_0(x) = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-\frac{m\omega}{2\hbar}x^2}$$

What is the energy  $E_0$ ?



$$\underbrace{\hbar\omega(a_+a_- + 1/2)}_{H \text{ (Hamiltonian)}} \psi_0 = E_0 \psi_0$$

=0

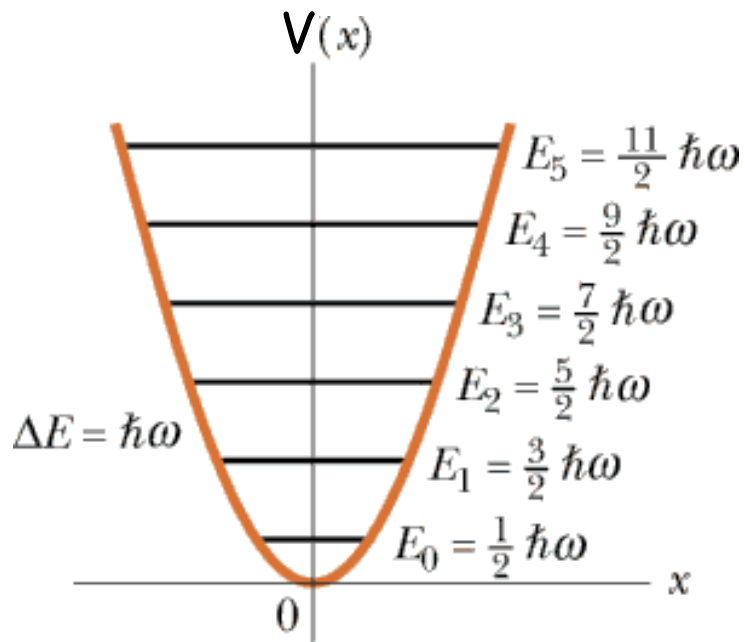
Then:

$$E_0 = \frac{1}{2}\hbar\omega$$

>0 as expected.

“Zero point energy”.

The harmonic oscillators  
are never still!



Equally spaced levels

$$E_n = \left( n + \frac{1}{2} \right) \hbar \omega$$

Solutions are

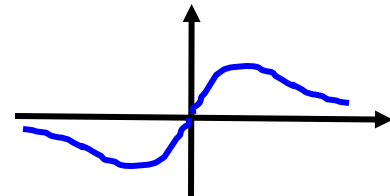
$$\psi_n(x) = A_n (a_+)^n \psi_0(x)$$

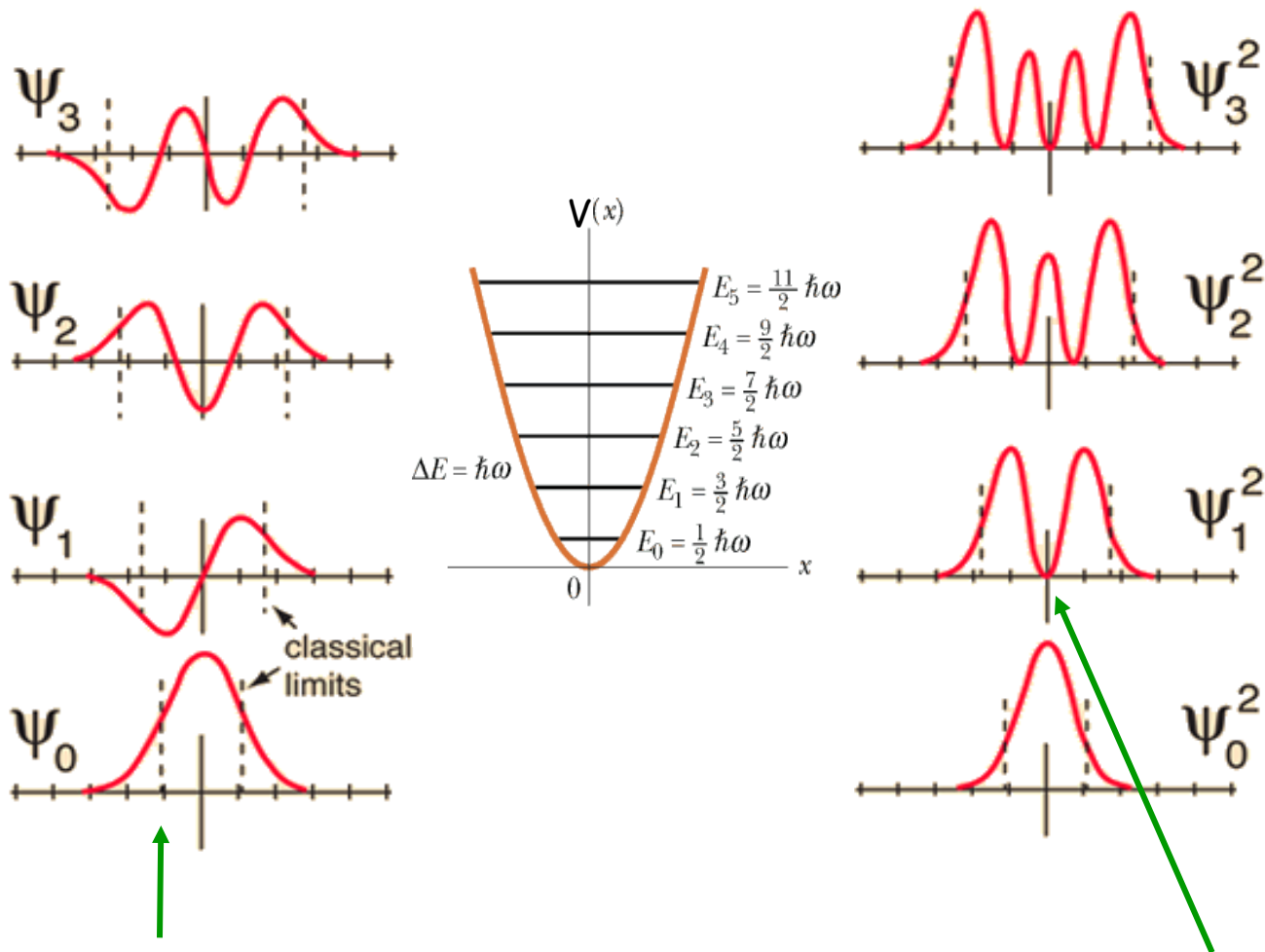
Example: construct state 1

$$\psi_1(x) = A_1 a_+ \psi_0 = \frac{A_1}{\sqrt{2\hbar m\omega}} \left( -\hbar \frac{d}{dx} + m\omega x \right) \left( \frac{m\omega}{\pi \hbar} \right)^{1/4} e^{-\frac{m\omega}{2\hbar} x^2}$$

$$= A_1 \left( \frac{m\omega}{\pi \hbar} \right)^{1/4} \sqrt{\frac{2m\omega}{\hbar}} x e^{-\frac{m\omega}{2\hbar} x^2}$$

Brings down an  $x$





**Note:** particle can be found outside the classical region

Zero chance at the nodes.

It can be shown, following the textbook, that:

$$\begin{aligned} a_+ \psi_n &= \sqrt{n+1} \psi_{n+1} & \psi_n &= \frac{1}{\underbrace{\sqrt{n!}}_{A_n}} (a_+)^n \psi_0 \\ a_- \psi_n &= \sqrt{n} \psi_{n-1} \end{aligned}$$

$$\int_{-\infty}^{\infty} \psi_m^* \psi_n dx = \delta_{mn}$$

Orthonormal like we found before for square well.


**Example 2.5:** find expectation value of  $V$  in the  $n$ th state.

$$\langle V \rangle = \left\langle \frac{1}{2}m\omega^2 x^2 \right\rangle = \frac{1}{2}m\omega^2 \int_{-\infty}^{\infty} \psi_n^* x^2 \psi_n dx.$$

It may be tempting to write  $\psi_n$  as a Gaussian with some polynomial in front. However, there is a **simpler** path.

From  $a_{\pm} \equiv \frac{1}{\sqrt{2\hbar m\omega}} (\mp ip + m\omega x)$  used before, deduce:

$$x = \sqrt{\frac{\hbar}{2m\omega}} (a_+ + a_-); \quad p = i\sqrt{\frac{\hbar m\omega}{2}} (a_+ - a_-).$$


$$x^2 = \frac{\hbar}{2m\omega} \left[ (a_+)^2 + (a_+ a_-) + (a_- a_+) + (a_-)^2 \right]$$

Creates  $\psi_{n-2}$  orthog to  $\psi_n$

$$\langle V \rangle = \frac{\hbar\omega}{4} \int \psi_n^* \left[ (a_+)^2 + (a_+a_-) + (a_-a_+) + (a_-)^2 \right] \psi_n dx.$$

Creates  $\psi_{n+2}$  orthog to  $\psi_n$

$$a_+ \psi_n = \sqrt{n+1} \psi_{n+1}, \quad a_- \psi_n = \sqrt{n} \psi_{n-1}$$

$$a_- a_+ \psi_n = a_- \sqrt{n+1} \psi_{n+1} = \sqrt{n+1} \underbrace{a_- \psi_{n+1}}_{\sqrt{n+1} \psi_n} = (n+1) \psi_n$$

$a_+ a_-$        $a_- a_+$

$$\langle V \rangle = \frac{\hbar\omega}{4} (n + n + 1) = \frac{1}{2} \hbar\omega \left( n + \frac{1}{2} \right) = \frac{1}{2} E_n$$